We thank the referee for the thorough evaluation of the manuscript and the interesting points raised. These will help improve the clarity of the paper. A point to point response follows, where the referee comments are in italic and our respective answers given below.

The authors identify the importance of sea-ice for the occurrence of abrupt climate changes and couple its dynamics into a Stommel box model. The authors may want to put their model with its inherent dynamic mechanisms in context to other similar attempts highlighting the different implied dynamical mechanisms; for example the conceptual models considered by Boers et al, Proc Natl Acad Sci 115(47):E11005–E11014 and by Gottwald, Clim Dyn (2021) 56:227-243.

We agree that it is appropriate to add some comments regarding similar efforts with conceptual models. This will be added in the introduction. The work by Gottwald considers sea ice as an intermittent thermal insulator to the polar ocean, represented by the Stommel model, using a chaotic Lorenz ‘84 model to act as effectively stochastic forcing. The sea ice component then becomes a deterministic approximation of a correlated additive and multiplicative noise process, extremes of which can trigger temporary excursions of the Stommel model. This is similar in spirit to our stochastic forcing of the sea ice component. However, we include a deterministic underlying parameter shift (e.g. changes in ice sheet extent and height) as the main cause of the abrupt transitions, which is in our opinion more in line with existing evidence of the mechanism of DO events. Boers et al also consider a coupling of sea ice/ice shelf to an ocean box model. Here, the sea ice evolves due to a prescribed piecewise-linear feedback, which leads to self-sustained oscillations. Our model differs in that the sea ice dynamics also involve a tipping point, and the ocean component features rate-induced tipping, thus giving the possibility of the cascades of tipping points discussed in the manuscript.

As a result, our study differs from previous studies both in terms of the dynamical mechanism, as well as the interpretation of the underlying driver of the abrupt transitions.

The authors tune the parameter h in (5) to allow for what they coin smooth bifurcations. Are there observations suggesting that transitions are smooth or non-smooth? Also, it might be helpful to show the two cases of non-smooth and smooth in Figure 3 (and also for the Stommel model). This would clarify what the authors mean by smooth and non-smooth bifurcations; depending on the background of the reader these terms may invoke different associations. Might also be worthwhile defining this in the manuscript.
We acknowledge that the slightly loose use of the word “smooth” causes some confusion, which we will resolve with more careful wording and some additional explanations in the revised manuscript.

First, we are actually not attempting to coin the term “smooth bifurcation”, we just say that in the sea ice model a larger value for $h$ gives a smoother transition of the albedo. This results in a “smoother” appearance of the bifurcation diagram, in the sense of a more “rounded” S-curve of the fold-fold bifurcations. The extreme case for $h \rightarrow 0$ would be a Z-curve instead of an S-curve. What cannot be seen from the bifurcation diagram is that in the less smooth case the curvature of the underlying potential only starts to change significantly as one get very close to the bifurcation, limiting the detectability of critical slowing down. We will consider using terminology other than “smooth” here to distinguish from the following.

On the other hand, the Stommel model does have a truly “non-smooth” bifurcation, which comes from a discontinuity in the flow (it is a non-smooth dynamical system).

We will more carefully define our usage of the word non-smooth in this case. The main argument to do so is already given in the paper: The stable and unstable fixed points meet in a cusp, as opposed to a fold as is the case for a “smooth” bifurcation (see Fig. 4).

As a result, the attractor and saddle can come very close to each other already significantly far away from the bifurcation point. In the paper, we consider the shortest distance of the basin boundary to the attractor as a function of the control parameter.

Figure 7b can then be understood to define the smooth versus non-smooth bifurcation. To make this even clearer, we will include another figure as a subfigure (or in the supplemental material), which shows how the fixed points move as $\eta_1$ is changed, and how they merge tangentially in the bifurcation, together with the real part of the first eigenvalue of the Jacobian to show the non-smoothness of the flow. This will help illustrate how the non-smoothness of the flow is related to the collision of the fixed points in a cusp.

Regarding the choice of $h$ and real-world observations, it is difficult to say how smooth the albedo transition should be considered to be. This depends on what is actually modeled by the sea ice variable.

Since we are modeling a large ocean basin, we considered it more appropriate to use a more gradual albedo transition, corresponding to a wider range of partial sea ice cover.

The choice of $h$ does not change any of our results however, besides the fact that for lower values of $h$ (and thus a “less smooth” bifurcation diagram), it would be more difficult to detect a critical slowing down in the sea ice variable, which is however not a main focus of the paper. We will add bifurcation diagrams for smaller versus larger values of $h$ in the Supp material.
It was not clear to me how their model allows for the succession of abrupt climate changes such as the DO events mentioned in the introduction. The model seems to capture only single transitions. Can the authors comment on this?

Indeed the model in its present form only captures individual warming transitions. What we tried to point out in the manuscript is that the parameter shift itself can be modeled by a further slow variable, which could be a simple negative feedback reflecting, e.g., the influence of the AMOC on the ice sheets. The ice sheets could then drive the sea ice by their influence on atmospheric circulation. This allows for oscillations in between the ‘on’ and ‘off’ states, and thus successive abrupt climate changes. Many other extensions are possible, which could capture a variety of different dynamical mechanisms, such as excitability. We will try to be more explicit about this in the revised manuscript.

Regarding the new proposed warning indicator J. Am I correct in thinking that the reason why looking at the individual elements of the Jacobian rather than at the eigenvalues of the Jacobian is that the estimation of each element is done via finite-differencing (which is a bad estimator for noisy data) and calculating the eigenvalues exacerbates this via multiplication?

Indeed this is one of the reasons, as also mentioned on page 17 in the manuscript. The other reason being a bias in the estimation of the elements, which affects the eigenvalues systematically (Fig. S2). This could be due to the method of finite-differencing. We will investigate whether there are better methods in future studies, which will be more method-heavy and thus not suitable to include here. This does however not change the general idea of the early-warning method.

Figure 1. What are the values of $\eta_3$ in (b) and of $\eta_1$ in (c)

Will add the values to the caption.

Line 266: andBoers —> and Boers

Ok.