

# 1 Labrador Sea sub-surface density as a precursor of 2 multi-decadal variability in the North Atlantic: a multi-model 3 study.

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14 **Abstract.** The Subpolar North Atlantic (SPNA) is a region with prominent decadal variability that has experienced  
15 remarkable warming and cooling trends in the last few decades. These observed trends have been preceded by  
16 slow-paced increases and decreases in the Labrador Sea density (LSD), which are thought to be a precursor of large  
17 scale ocean circulation changes. This article analyses the inter-relationships between the LSD and the wider North  
18 Atlantic across an ensemble of coupled climate model simulations. In particular, it analyses the link between  
19 subsurface density and the deep boundary density, the Atlantic Meridional Overturning Circulation (AMOC), the  
20 Subpolar Gyre (SPG) circulation, and the upper ocean temperature in the eastern SPNA.

21

22 All simulations exhibit considerable multidecadal variability in the LSD and the ocean circulation indices, which are  
23 found to be interrelated. LSD is strongly linked with the strength of subpolar AMOC and gyre circulation, and is  
24 also linked with the subtropical AMOC, although the strength of this relationship is model dependent and affected  
25 by the inclusion of the Ekman component. The connectivity of LSD with the subtropics is found to be sensitive to  
26 different model features, including: the mean density stratification in the Labrador Sea; the strength and depth of the  
27 AMOC; and the depth at which the LSD propagates southward along the western boundary. Several of these  
28 quantities can also be computed from observations, and comparison with these observation-based quantities suggests  
29 that models representing a weaker link with the subtropical AMOC might be more realistic. This would imply  
30 that RAPID AMOC measurements might not be adequate to represent decadal to multidecadal changes in the  
31 subpolar overturning circulation.

32 :

## 33 1. Introduction

34

35 The North Atlantic Ocean is a key component in Earth's climate through, for example, its role in redistributing heat  
36 and in taking up excess heat and carbon from the atmosphere. It is also a region that has varied significantly in the  
37 past. This is particularly true for the North Atlantic subpolar gyre, that has varied significantly on multi-decadal  
38 timescales across a range of different variables (H kkinen and Rhines, 2004; Holliday et al., 2020; Reverdin, 2010;  
39 Robson et al., 2018b). Basin-mean sea surface temperature (SST) over the North Atlantic has also been observed to  
40 vary on multi-decadal timescales (Schlesinger and Ramankutty, 1994), and has been linked to a range of important

41 climate impacts, including hurricane numbers and rainfall in monsoon regions (Knight et al., 2006; Monerie et al.,  
42 2019; Zhang and Delworth, 2006). The North Atlantic is also expected to change significantly in the future due to  
43 the effects of climate change, and consequently produce substantial climate impacts on the surrounding regions  
44 (Sutton and Hodson, 2005; Woollings et al., 2012). On decadal timescales, it is the interaction between natural  
45 variability and externally forced changes that will shape how the Atlantic regions climate will evolve. Therefore, in  
46 order to improve predictions of the North Atlantic, it is imperative that we improve our understanding of the  
47 processes that control decadal timescale changes in this region.

48

49 It has generally been thought that changes in the ocean circulation, and particularly the Atlantic Meridional  
50 Overturning Circulation (AMOC), have played a significant role in shaping the Atlantic Mmulti-decadal  
51 Vvariability across the North Atlantic (AMV; Knight et al. 2005). In particular, changes in the strength of the  
52 AMOC, and its related ocean heat transports have been shown to control multi-decadal internal variability in a range  
53 of coupled climate models (Danabasoglu, 2008; Dong and Sutton, 2005; Jungclaus et al., 2005; Ortega et al., 2011,  
54 2015). The proposed mechanisms to explain the multi-decadal variability involve interplays between the North  
55 Atlantic Oscillation (NAO), ~~The mechanisms proposed~~ One of the mechanisms consistently identified involves  
56 ~~changes in~~ North Atlantic Deep Water (NADW) formation, the boundary currents, the Gulf Stream and gyre  
57 circulations, and the horizontal density gradients (e.g. Joyce and Zhang, 2010; Polyakov et al., 2010; Ba et al., 2013;  
58 Nigam et al., 2018; Zhang et al., 2019). Changes in AMOC and the wider ocean circulation have been indeed used  
59 to explain the ~~observed~~ changes observed in the subpolar North Atlantic (SPNA); on decadal and longer timescales  
60 (Moat et al., 2019). In particular, the SPNA underwent a rapid warming and salinification in the mid 1990s before a  
61 decadal timescale cooling and freshening started in 2005, which is consistent with decadal-to-multidecadal  
62 variability of the AMOC (Robson et al., 2012, 2013, 2016). The recent cooling has been linked to climate impacts  
63 over the continents, including heat waves (Duchez et al., 2016), through an effect on the position on the jet stream  
64 (Josey et al., 2018). A long term relative cooling of the SPNA since ~1850 has also been attributed to a centennial  
65 weakening of the AMOC (Caesar et al., 2018; Rahmstorf et al., 2015), an AMOC reduction that most CMIP6 model  
66 projections predict to continue in the future (Weijer et al., 2020). However, a lack of direct observations of the  
67 strength of the AMOC or the ocean circulation more generally have hindered our ability to make a direct attribution  
68 of recent changes.

69

70 In order to understand the aforementioned changes in the SPNA on multi-decadal timescales many authors have  
71 turned to indirect measurements of the AMOC. One particular proxy of AMOC strength that has received some  
72 focus recently are density anomalies at depth in the western SPNA or Labrador Sea region. In climate models,  
73 density anomalies in the western SPNA are a key predictor of density anomalies further south on the western  
74 boundary, and hence of the AMOC strength via thermal wind balance (Hodson and Sutton, 2012; Ortega et al.,  
75 2017; Robson et al., 2014, 2016). Observations show considerable decadal variability in subsurface density  
76 anomalies; density anomalies in the western SPNA or Labrador Sea between ~1000-2500 m increased significantly  
77 and peaked in ~1995 and subsequently declined (Robson et al., 2016; Yashayaev and Loder, 2016). Therefore, these  
78 density anomalies have been interpreted as indicating that the AMOC peaked circa mid-to-late 1990s, and then  
79 declined, consistent with the warming and then cooling of the eastern SPNA (Hermanson et al., 2014; Ortega et al.,  
80 2017; Robson et al., 2016). Time series of subsurface density anomalies in the western SPNA are also consistent  
81 with other proxies of AMOC strength, including sea level based proxies (McCarthy et al., 2015; Sutton et al.,  
82 2018), sediment based proxies (Thornalley et al., 2018), and upper ocean heat content fingerprints (Caesar et al.,  
83 2018; Zhang, 2008). Furthermore, the decline in AMOC suggested by the above proxies is also consistent with the  
84 observed AMOC decline at 26°N since 2004 (Smeed et al., 2018), and also with the changes in AMOC seen in  
85 ocean data assimilation systems (Jackson et al., 2016, 2019). Therefore, there is confidence that large scale changes  
86 in North Atlantic Ocean circulation have occurred over the past few decades, and that they have had a significant  
87 impact on upper ocean heat content.

88

89 Although there is consistency across proxies of AMOC changes in the North Atlantic, there are considerable gaps in  
90 our understanding and major uncertainties to overcome. For example, the development of the subsurface density  
91 proxies has been investigated so far in just a few models (Ortega et al., 2017; Robson et al., 2014). However, there is  
92 considerable spread across climate models in the simulations of AMOC mean state and variability (Reintges et al.,  
93 2017; Zhang and Wang, 2013), and also in the latitudinal coherence of AMOC anomalies (Li et al., 2019; Roberts et  
94 al., 2020; Hirschi et al., 2020), which might reflect different roles of deep density anomalies in the western SPNA on  
95 the AMOC, as well as different interplays between the subpolar and subtropical gyre contributions (Zou et al.,  
96 2020). Models also do not resolve realistically many key features of AMOC, most notably the overflows, and this  
97 affects the subsurface stratification downstream and on the western boundary (Zhang et al., 2011). There also  
98 remains significant uncertainty for other important processes. For example, it is not yet clear whether the recent  
99 changes in the SPNA are an ocean response to buoyancy forcing, or whether mechanical wind forcing has shaped  
100 the recent observed changes (Robson et al, 2016; Piecuch et al. 2017). Local surface fluxes are also likely to explain  
101 a significant proportion of the recent cooling (Josey et al, 2018). Subsurface density anomalies are not just a proxy  
102 for the AMOC, but more generally for buoyancy forced (or thermohaline) circulation changes, including gyre  
103 changes (Ortega et al., 2017; Yeager, 2015). Finally, the AMOC variability is also thought to respond to local wind  
104 forcing on a range of timescales, especially at lower latitudes (Polo et al., 2014; Zhao and Johns, 2014), which could  
105 disrupt or “mask” the influence of subsurface density anomalies as they propagate further south.

106

107 There is also considerable uncertainty in how and where subsurface density anomalies are formed in the SPNA, and  
108 how they are related to the AMOC. In observations and models, most water transformation associated with the  
109 AMOC occurs within the SPNA, and particularly in the eastern SPNA (Desbruyères et al., 2019; Grist et al., 2014;  
110 Langehaug et al., 2012). However, subsurface density anomalies in the western SPNA on decadal timescales have  
111 often been linked with buoyancy forcing and changes in deep convection in the Labrador Sea or with changes in the  
112 volume of Labrador Sea Water production (Yashayaev and Loder, 2016; Yeager and Danabasoglu, 2014). Many  
113 studies have also reported that the basin-wide AMOC in ocean-only and coupled models is sensitive to heat flux or  
114 buoyancy forcing in the Labrador Sea (Kim et al., 2020~~19~~; Ortega et al., 2011, 2017; Xu et al., 2019; Yeager and  
115 Danabasoglu, 2014). **Indeed, idealised experiments have shown that persisting positive NAO phases can strengthen  
116 the AMOC by fostering deep water formation via increased surface cooling in the Labrador Sea, thus inducing  
117 changes in the zonal density gradient (Delworth and Zeng, 2016; Kim et al., 2020), and thermal wind responses.**  
118 However, the **real** link between deep convection, deep water formation, and density anomalies at depth in the  
119 Labrador Sea is complex, and not fully understood (Katsman et al., 2018). Observations suggest that very little  
120 water transformation and deep water formation actually occurs in the Labrador Sea (Pickart and Spall 2007; Lozier  
121 et al., 2019). Indeed, recently it has been shown that the Labrador Sea (i.e. OSNAP-west) played very little role in  
122 the interannual variability so far observed across the whole OSNAP line (Lozier et al., 2019), **with the Irminger Sea  
123 playing a more dominant role. The Irminger Sea is a region that in some models controls the AMOC and SPNA  
124 variability, and that is especially sensitive to advective processes (Ba et al., 2013) and Arctic overflows (Fröb et al.,  
125 2016).** Moreover, ocean-only models appear to significantly overestimate the amount of deep water formed within  
126 the Labrador Sea, with likely implications for coupled models (Li et al., 2019). These inconsistencies raise the  
127 question of whether models are simulating the right relationships.

128

129 In this study we will address some of the above uncertainties by performing a multi-model analysis of the North  
130 Atlantic in coupled climate models. We focus on the question of how robust is the relationship between subsurface  
131 Labrador Sea density anomalies and the basin-wide Atlantic Ocean circulation on decadal timescales. We also assess  
132 the question of whether Labrador Sea density can robustly induce density changes over the western continental slope  
133 and generate a geostrophic response in the meridional circulation (Bingham and Hughes, 2009; Roussenov et al.,  
134 2008). Shedding new light on these links is important for, among other reasons, determining to what extent the  
135 RAPID measurements represent the variability of the basin-wide AMOC cell, as well as to identify the models that  
136 can produce more reliable predictions and projections of the SPNA. For this, we will assess specifically the

137 connection between subsurface density and AMOC at high and low latitudes via the western boundary. Furthermore,  
138 we will determine whether models consistently support an impact of AMOC changes on the SPNA upper ocean  
139 temperatures, and if not, investigate why. Our primary aim is to provide, for the first time in a multi-model context,  
140 a broad characterization of these relationships using consistent analysis frameworks and tools, and to document the  
141 uncertainty. The reasons for the uncertainty in the relationships will also be explored, establishing links with key  
142 model climatological properties that could eventually be exploited as emergent constraints. We intentionally do not  
143 explore in detail how subsurface density anomalies are formed in these models, and leave this for further study.

144

145 The paper is organised as follows. Section 2 describes the experiments and methods. Labrador Sea density, and its  
146 link with the ocean circulation and the wider North Atlantic are explored across the multi-model ensemble in  
147 Section 3. The characteristics of the intermodel spread in the previous relationships are explored in Section 4, and  
148 Section 5 presents the main conclusions of this study and discusses its implications.

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150

## 151 2. Experiments and methods

152

153 Here we provide an overview and brief description of the models used in this study and provide some statistical  
154 considerations for the intermodel comparison.

155

### 156 2.1. Experiment selection

157

158 For the multimodel analysis, we use the preindustrial control simulations (picontrol) from the fifth phase of the  
159 Coupled Model Intercomparison Project (CMIP5; Taylor et al. 2012), in which forcing values of GHGs, aerosols,  
160 ozone and solar irradiance are fixed to 1850 levels. We chose to use control over historical simulations to focus  
161 exclusively on internal variability and benefit from the more robust statistics that the long preindustrial experiments  
162 provide. Furthermore, we avoid the forced trends present in the historical experiments, which can lead to  
163 correlations that are difficult to interpret objectively (Tandon and Kushner, 2015). From the CMIP5 ensemble in  
164 particular, we only use those models in which 3D fields of ocean temperature and salinity, as well as the  
165 streamfunctions of meridional overturning circulation and/or the barotropic circulation, were available. Twenty  
166 different models meet this condition, and their main characteristics and number of simulation years have been  
167 summarized in Table 1. Most of the models have a nominal horizontal resolution in the ocean close to  $1^\circ$ , and,  
168 therefore, cannot resolve the effects of eddies. Menary et al. (2015) has shown for these same model simulations that  
169 the effective horizontal resolution can be higher over the Labrador Sea, due to the non-regular grids. Effective  
170 resolutions over the Labrador Sea area range from  $0.2147^\circ$  in the GC2MPI model to  $1.16^\circ$  in  
171 GISS-E2-R/GISS-E2-R-CC/CanESM2IPSE, with these differences determining to a large extent the mean state  
172 model biases and the dominant drivers (i.e. salinity or temperature) of the Labrador Sea density changes.

173

174 Complementing these simulations, we also consider two control experiments with eddy-permitting resolutions.  
175 Specifically, we use a present day control simulation (i.e. with fixed radiative forcing levels from year 1990) of the  
176 HiGEM model, with nominal horizontal resolution in the ocean of  $1/3^\circ$ , and of  $0.83^\circ$  latitude  $\times$   $1.25^\circ$  longitude in  
177 the atmosphere (Shaffrey et al., 2009), and a pre-industrial control of HadGEM3-GC2 (hereafter, GC2; Ortega et al.  
178 2017) with a nominal resolution in the ocean of  $1/4^\circ$  (ORCA025) and N216 in the atmosphere (i.e. approximately 60  
179 km in the mid-latitudes). The GC2 simulation is the same one employed for the previous analyses of Labrador Sea  
180 variability in Robson et al. (2016) and Ortega et al. (2017). Note that we will assume that the present day control in  
181 HiGEM can be compared with the other preindustrial simulations due to the large uncertainty these later show in  
182 their climatological biases, and so, for the sake of simplicity, we will only refer to preindustrial control experiments

183 from now on. [Figure 1](#) demonstrates that this assumption is reasonable, since the mean Labrador Sea stratification in  
 184 HiGEM is very similar to that in the other models.

185

186 As an observationally-constrained reference, this study also includes the assimilation run from DePreSys3, a decadal  
 187 prediction system from the MetOffice based on GC2 (Dunstone et al., 2016). In the ocean, the assimilation is  
 188 performed through a strong nudging (ten-day relaxation timescale) towards the full fields of a three-dimensional  
 189 objective temperature and salinity analysis (Smith and Murphy, 2007). Since it covers a comparatively shorter  
 190 period (1960-2013), and therefore different timescales than the control experiments, its comparison with the other  
 191 simulations will be done with caution, in particular regarding the indices of the large-scale Atlantic circulation, for  
 192 which other assimilation products show important discrepancies (Karspeck et al., 2015), thus highlighting  
 193 significant uncertainty. For evaluation purposes, we also use EN4.2.1 (Good et al., 2013), an objective analysis of  
 194 monthly temperature and salinity 3D observations developed at the MetOffice.

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**Table 1:** List of the models used for this study, their characteristics and those of their picontrol simulations. For further details on the CMIP5 model configurations and components please refer to Table 9.A.1 in Flato et al (2013) and references therein.

Model ID	Lon x Lat ocean resolution (number of vertical levels)	Length	Key variables available
HadGEM3-GC2	1/4° x 1/4° (75 levels)	311 years	AMOC, SPGSI, LSD, NOHT
HiGEM	1/3° x 1/3° (40 levels)	341 years	AMOC, SPGSI, LSD, NOHT
ACCESS1-0	1° x 1° enhanced near Equator and high latitudes (50 levels)	500 years	SPGSI, LSD, NOHT
ACCESS1-3	1° x 1° enhanced near Equator and high latitudes (50 levels)	500 years	SPGSI, LSD, NOHT
CCSM4	1.125° x 0.27–0.64° (60 levels)	1051 years	AMOC, SPGSI, LSD
CESM1-BGC	1.125° x 0.27–0.64° (60 levels)	500 years	AMOC, LSD
CESM1-CAM5	1.125° x 0.27–0.64° (60 levels)	319 years	AMOC, LSD
CESM1-FASTCHEM	1.125° x 0.27–0.64° (60 levels)	222 years	AMOC, LSD
CESM1-WACCM	1.125° x 0.27–0.64° (60 levels)	200 years	AMOC, LSD
CNRM-CM5	0.7° x 0.7° (42 levels)	850 years	AMOC, SPGSI, LSD
CanESM2	1.4° x 0.93° (40 levels)	996 years	AMOC, SPGSI, LSD
FGOALS-g2	1° x 1° with 0.5° meridional in the tropical region (30 levels)	700 years	AMOC, LSD
FGOALS-s2	1° x 1° with 0.5° meridional in the tropical region (30 levels)	501 years	SPGSI, LSD, NOHT
GFDL-ESM2G	1° x 0.85° (63 levels)	500 years	SPGSI, LSD



GISS-E2-R	1.25° x 1° (32 levels)	550 years	AMOC, LSD
GISS-E2-R-CC	1.25° x 1° (32 levels)	251 years	AMOC, LSD
MPI-ESM-LR	1.5° x 1.5° (40 levels)	1000 years	AMOC, SPGSI, LSD
MPI-ESM-MR	0.4° x 0.4° (40 levels)	1000 years	AMOC, SPGSI, LSD
MPI-ESM-P	1.5° x 1.5° (40 levels)	1156 years	AMOC, SPGSI, LSD
MRI-CGCM3	1° x 0.5° (51 levels)	500 years	AMOC, LSD, NOHT
NorESM1-M	1.125° x 1.125° (53 levels)	501 years	AMOC, SPGSI, LSD, NOHT
NorESM1-ME	1.125° x 1.125° (53 levels)	252 years	AMOC, SPGSI, LSD, NOHT

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200

## 2.2. Methodological considerations

201 Density values are computed from 3D salinity and potential temperature fields, using the International Equation of  
202 State of seawater (EOS-80), and are referred to the level of 2000 dbar ( $\sigma_2$ ), to give a stronger emphasis to the deep  
203 water properties.

204 Statistical significance of correlation coefficients is assessed following a two-tailed Student's t-test that takes into  
205 account the series' autocorrelation to correct the sample size, reducing the degrees of freedom of a series to its  
206 effective value (Bretherton et al., 1999).

207 Because our goal is to provide further insight into the suggested relationships established from observed trends in  
208 the North Atlantic (e.g., Robson et al., 2016), all statistical analyses in this study exploring the relationships between  
209 variables and associated lags are based on 10-year running trends. This is analogous to the calculation of a typical  
210 10-year running mean, but computing over each 10 year period a linear trend instead and keeping the slope value.  
211 Note also that our main results remain similar if decadal running means are applied instead (not shown), as both are  
212 alternative approaches to concentrate on the low-frequency variability. Running trends have also the particular  
213 advantage of not being sensitive to long-term drifts, which are still present (and can be important for some  
214 simulations and variables) when running means are computed. To illustrate how decadal running trends represent  
215 low-frequency variability, and how they compare with the decadal running means, both these have been included in  
216 [Figure 2b](#) (solid thick lines vs dashed thin lines) for an index of Labrador Sea density.

## 217 3. Labrador Sea density as an index of multi-decadal North Atlantic variability

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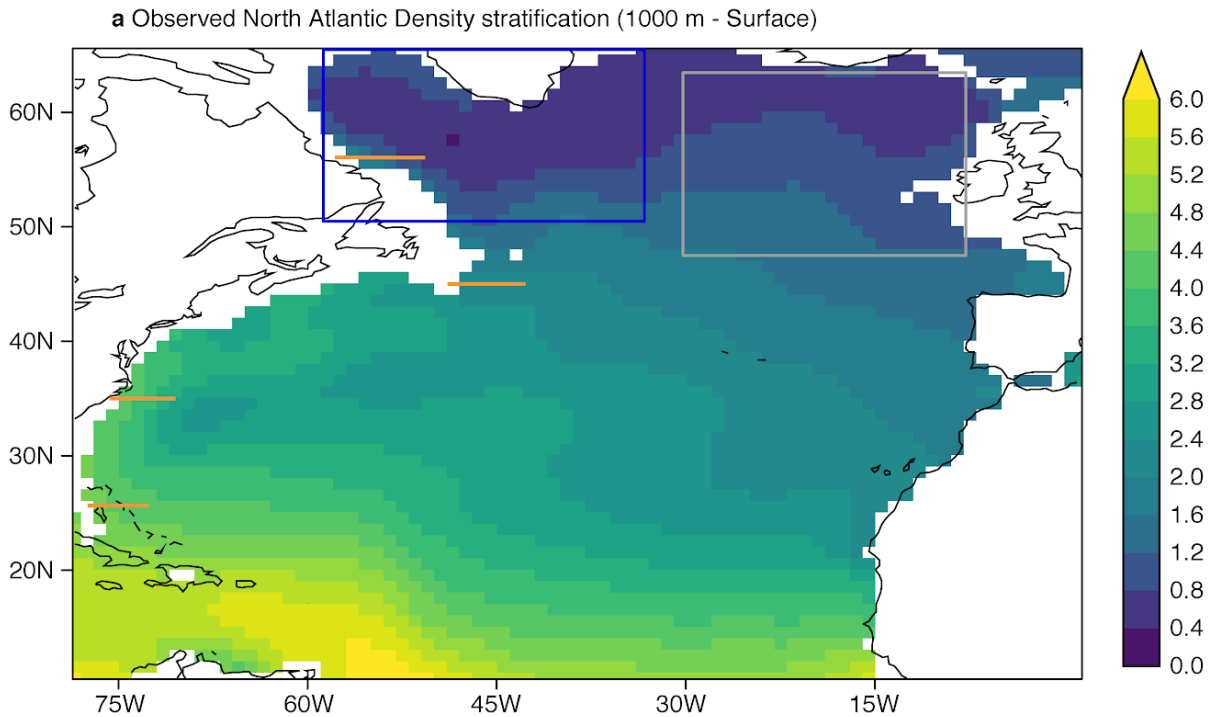
219 This section explores the potential of Labrador Sea density as a proxy of the ocean circulation changes in the North  
220 Atlantic. As in our previous studies (Ortega et al., 2017; Robson et al., 2016), the indices that we will herein define  
221 represent waters within the Labrador Sea and not those that are necessarily formed in the region (e.g. Labrador Sea  
222 Water). Since Labrador Sea variability is affected by different processes (e.g. vertical mixing, Arctic-Atlantic  
223 overflows, sea ice interactions) that can be represented differently in the models, both in time and space, we  
224 characterize its variability over a relatively broad box (60°W-35°W; 50°N-65°N, blue box [Figure 1a](#)) that also  
225 includes part of the Irminger Sea region. Note that over this large area EN4.2.1 (Good et al., 2013), an objective  
226 ~~analysis of monthly temperature and salinity 3D observations~~, shows the weakest density stratification in the North  
227 Atlantic (characterised in [Figure 1a](#) as the density difference between 1000m and the surface).

229 A first indicator of potential model discrepancies is Labrador Sea stratification, which can lead to differences in the  
230 representation of deep ocean convection (i.e. weaker density stratifications will facilitate the mixing, fostering  
231 convection activity, and vice versa for stronger density stratifications). [Figure 1b-d](#) illustrates the inter-model  
232 differences with the vertical profile of the spatially averaged Labrador Sea temperature, salinity and density. The  
233 largest discrepancies are seen for temperature. Most models present their warmest waters at the surface, and  
234 temperatures decrease sharply to minimum values around 100 m and increase again at deeper levels, reaching  
235 uniform conditions after approx 300 m. However, the location and magnitude of this temperature minimum and the  
236 two maxima are highly variable. It is important to note that the profile for one of the models, MRI-CGCM3, is  
237 noticeably different to the others, with a subsurface minimum more than 2 degrees colder than for any of the other  
238 models. In terms of salinity, the general profile is more coherent across models, with minimum salinity at the surface  
239 that progressively increases with depth and attains uniform values after 500 m. Density stratification seems to be  
240 determined by salinity, as their two vertical profiles show similar features. This similarity includes exceptionally  
241 strong density and salinity stratification in MRI-CGCM3 as compared with the other models. This stratification is so  
242 strong that it precludes the occurrence of deep convection (not shown). Because of this, MRI-CGCM3 is an outlier  
243 for many of the metrics used in the paper, and has been excluded for the subsequent analyses to facilitate the  
244 interpretation of our results. We also note that the profiles for the two eddy-permitting models (green and orange  
245 lines in [Figure 1b,d](#)) lie within the spread of the CMIP5 models, indicating that resolution (at least to  
246 eddy-permitting spatial scales) does not drastically change stratification in the region. The DePreSys3 assimilation  
247 run closely matches the stratification in EN4.2.1, which supports DePreSys3 assimilation run as a reasonable  
248 observation-constrained reference for the models. The comparison of both observation-based datasets with the rest  
249 of simulations suggests that, in the subsurface, all models are too warm and most of them are too salty, two biases  
250 that have a competing effect on the mean subsurface density. Because of these canceling effects, several models  
251 show a comparatively better representation of the subsurface densities when compared to EN4.2.1 and  
252 DePreSys3. This compensation of model shortcomings for temperature and salinity is clearly illustrated in HiGEM,  
253 which shows a remarkable agreement with EN4.2.1 below 500 m.

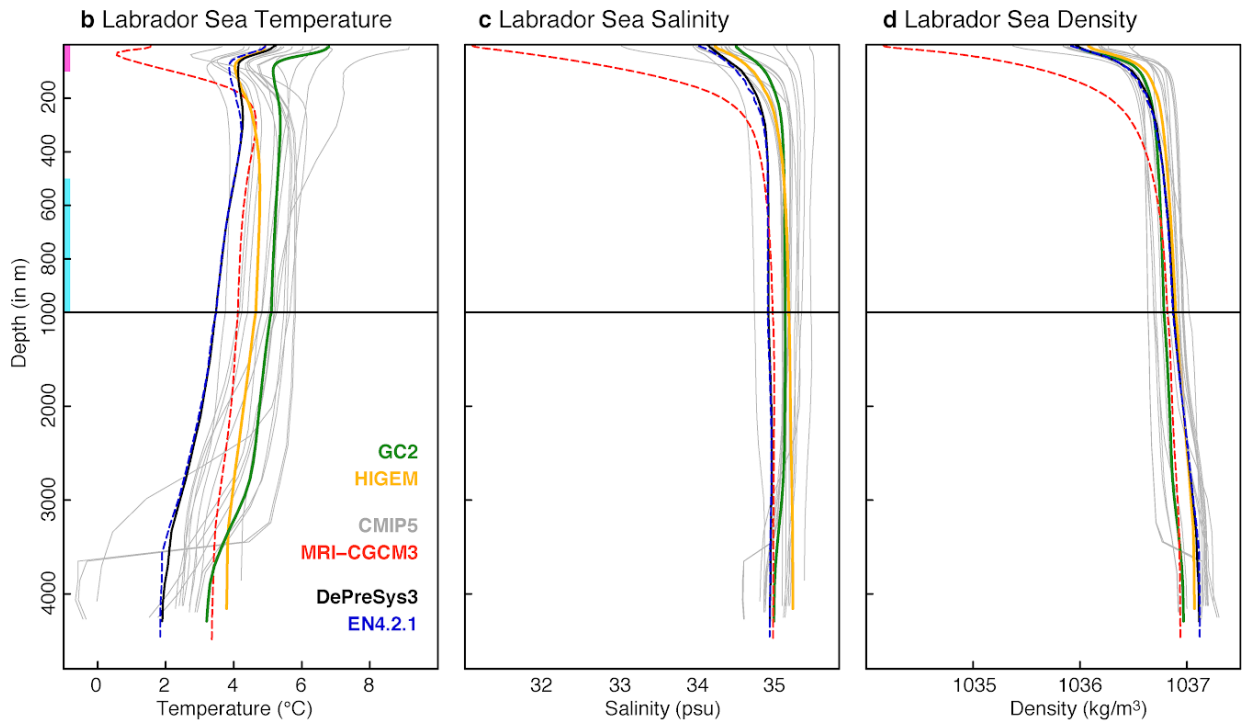
254 To represent the characteristic interannual variability of Labrador Sea densities (hereafter referred to as LSD for  
255 consistency with previous work), we perform an Empirical Orthogonal Function (EOF; [Storch and Zwiers, 1999](#))  
256 analysis and extract the leading mode for the spatially averaged annual means of LSD ([Figure 2a](#)), as in [Ortega et al.](#)  
257 (2017). For all simulations the first EOF of LSD exhibits a vertical structure where density values are largest at or  
258 near the surface and gradually decrease with depth. Thus, this first EOF typically reflects situations in which the  
259 density stratification, as described by the climatological vertical profile in [Figure 1c](#), is weakened or strengthened,  
260 which happens when depending on the sign of the corresponding principal component takes positive and negative  
261 values, respectively. Some inter-model discrepancies are evident, in particular regarding the depths where the  
262 maximum density values are found, which can happen between the surface and 500 m. Despite these differences, the  
263 dominant timescales of LSD variability seem to coincide between models. For example, [Figure 2b](#) illustrates the  
264 first principal component of LSD (PC1-LSD) for GC2 and HiGEM, showing in both cases clear important  
265 multidecadal variability. Furthermore, [Figure 2c](#) shows the Fourier spectrum analysis of the annual PC1-LSD values,  
266 and most models show enhanced PC1-LSD variability for periodicities between 5 and 30 years.

267 In addition to the PC1-LSD index we consider a ~~t~~ is worth noting that the PC1-LSD index used here is different to  
268 the deep LSD index as introduced in [Robson et al \(2016\)](#). The deep LSD index is defined as the 1000-2500 m  
269 vertical mean of the spatially averaged density over the same region as PC1-LSD. We now compare how both  
270 indices represent the low-frequency changes in LSD, described in this paper as decadal running trends. A lead-lag  
271 correlation between the decadal trends in both PC1-LSD and deep LSD indices shows that they are strongly  
272 correlated in all models. However, some differences emerge when considering the lag of maximum correlation

273 (Figure 2d). This comparison might indicate, once again, that decadal variability of subsurface density is  
 274 concentrated at different depths in different models. It is also possible that both indices are sensitive to changes in  
 275 deep water formation in different locations (e.g. Irminger or GIN Seas), which could, hence, affect the depth and  
 276 maximum lag of the correlations. Nevertheless, we adopt PC1-LSD for the rest of the analyses, as it has the  
 277 advantage of adjusting in each model to the depths in which density variability is more prominent.

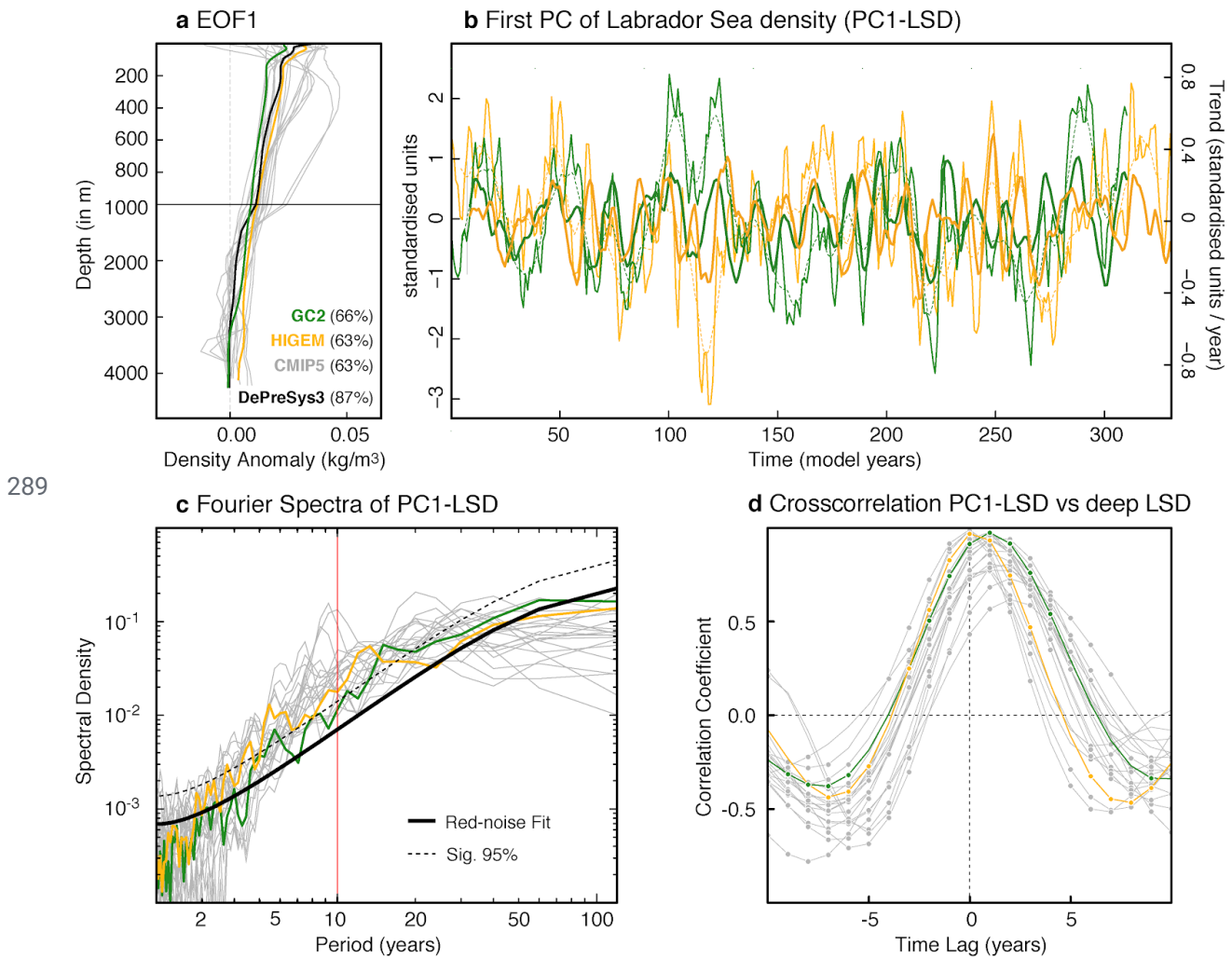


278





279 **Figure 1: a** Climatological density (computed as  $\sigma_2$  at all depth levels) difference between the subsurface (1000m)  
 280 and surface density ( $\sigma_2$ ) in the North Atlantic in the observational dataset EN4.2.1 (Good et al., 2013). The  
 281 reference period to compute the climatology is 1960-2013. The grey box (32°W-10°W and 47°N-63°N) encloses the  
 282 region where the ESPNA-T700 index in Figure 4d is computed. **b-d** Climatological mean of the spatially averaged  
 283 Labrador Sea (60–35°W, 50–65°N, blue box in panel a) temperature, salinity and density as a function of depth in  
 284 the simulation ensemble, the DePreSys3 assimilation run and EN4.2.1. The magenta (cyan) bars in the vertical axis  
 285 correspond to the depths that have been used to define the vertical stratification Labrador Sea indices. The horizontal  
 286 orange lines by the North American coast represent the location of the latitudinal cross-sections in Figures 10 and  
 287 11. For each model and dataset the climatology is computed for its whole length except for EN4.2.1, that is  
 288 computed for the overlap period with the DePreSys3 assimilation run.

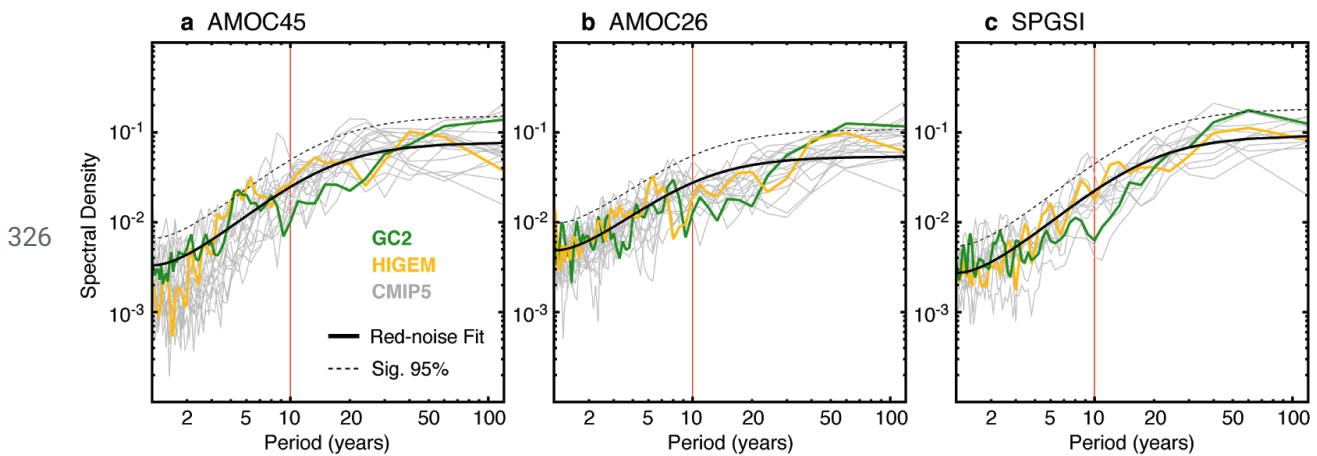


290 **Figure 2: a** First empirical orthogonal function (EOF) as a function of depth of the spatially averaged LSD in all the  
 291 preindustrial experiments and in the DePreSys3 assimilation run. The percentage of variance explained by this mode  
 292 in each model is included in brackets in the legend (for the CMIP5 runs, this represents the mean value across the  
 293 ensemble). **b** Associated principal component of the spatially averaged LSD (PC1-LSD) in the two high-resolution  
 294 experiments. The thin solid lines represent the raw yearly-resolved PC1-LSD timeseries, the thin dashed lines their  
 295 respective 10 year running means, and the thick (and slightly darker) lines represent their associated 10-year  
 296 running trends (centered around the last year of the decade over which the trend is computed). **c** Normalized Fourier

297 spectra of the PC1-LSD index in each of the preindustrial simulations. The black thick line represents a red noise  
 298 process with the same first autoregressive (AR1) coefficient as PC1-LSD in GC2, and the dashed line sets the 95%  
 299 confidence interval of this red-noise process. No major differences are found when using HiGEM's AR1 coefficient  
 300 instead. The red vertical line highlights the 10 year periodicity to separate the interannual from the decadal to  
 301 multi-decadal timescales. **d** Lead-lag correlations between the decadal trends in PC1-LSD, and those in the deep  
 302 LSD index from Robson et al. (2016), defined as the 1,000–2,500 m average density in the box 60–35°W, 50–65°N.  
 303 Positive lags indicate that PC1-LSD leads the changes in deep LSD. Full dots denote correlation values exceeding a  
 304 95% confidence level based on a student's t-test that takes into account the series autocorrelation.

### 305 3.2. Labrador Sea density linkages with the ocean circulation

306 The link between PC1-LSD and other ocean circulation indices in the North Atlantic is now examined. Three  
 307 indices are considered: the AMOC at two different latitudes, 26°N (i.e. the same latitude as the RAPID array) and  
 308 45°N to capture the typical variability of the subpolar AMOC, and an index of the subpolar gyre strength. The  
 309 AMOC indices are computed as the maximum of the North Atlantic overturning circulation at any depth.  
 310 Furthermore, the Ekman component is removed to focus on the slow wind-forced and the thermohaline-driven (i.e.  
 311 the only one that can be influenced by the PC1-LSD directly) AMOC changes. To compute the Ekman component,  
 312 we vertically integrate the Ekman velocities (after introducing a depth-uniform return flow to ensure no net  
 313 meridional mass transport), following Eq. 6 in Baehr et al. (2004) with a fixed Ekman layer depth of 50 meters. This  
 314 Ekman component is then removed at each depth level, prior to the calculation of the AMOC indices. The subpolar  
 315 gyre strength is computed as an average of the North Atlantic barotropic streamfunction in the Labrador Sea region  
 316 (60–35°W, 50–65°N), where the gyre strength is usually maximum. Since the SPG circulation is cyclonic and,  
 317 therefore, associated with negative barotropic streamfunction values, the subpolar gyre strength index (SPGSI) is  
 318 multiplied by -1 so that an intensification of the gyre corresponds to a positive value of the index. The Fourier  
 319 spectra of the raw ocean circulation indices (Figure 3) show that, similar to the PC1-LSD, all three indices have  
 320 strong multidecadal variability, with the largest differences with respect to PC1-LSD emerging for the timescales  
 321 between 10 and 30 years, in which the spectral power is comparatively weaker, in particular for the AMOC26 index,  
 322 and at 50 and longer timescales, in which the ocean circulation indices appear to have enhanced variability with  
 323 respect to PC1-LSD. Similar spectra, but with enhanced variance at short timescales and reduced variance at the  
 324 longest timescales are obtained for the AMOC indices when the Ekman component is kept (Supplementary Figure  
 325 1), which suggests that the low-frequency processes dominate the total AMOC variability.

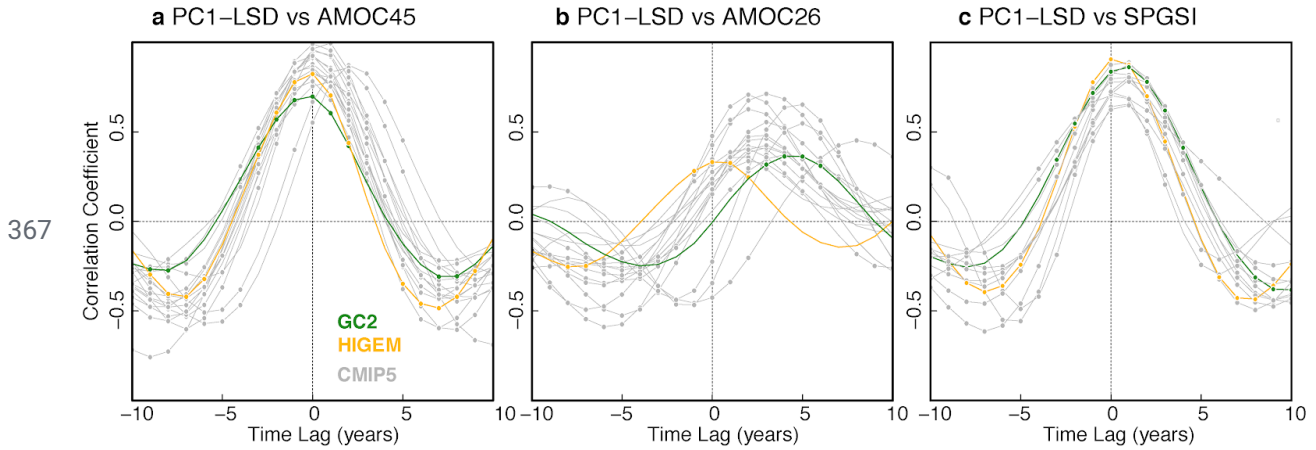


327 **Figure 3:** a-c Fourier spectra in the picontrol ensemble for the indices AMOC45N, AMOC26N and SPGSI.  
 328 Red-noise spectra corresponding to a 1<sup>st</sup> order autoregressive process fit to GC2 indices are provided as reference.

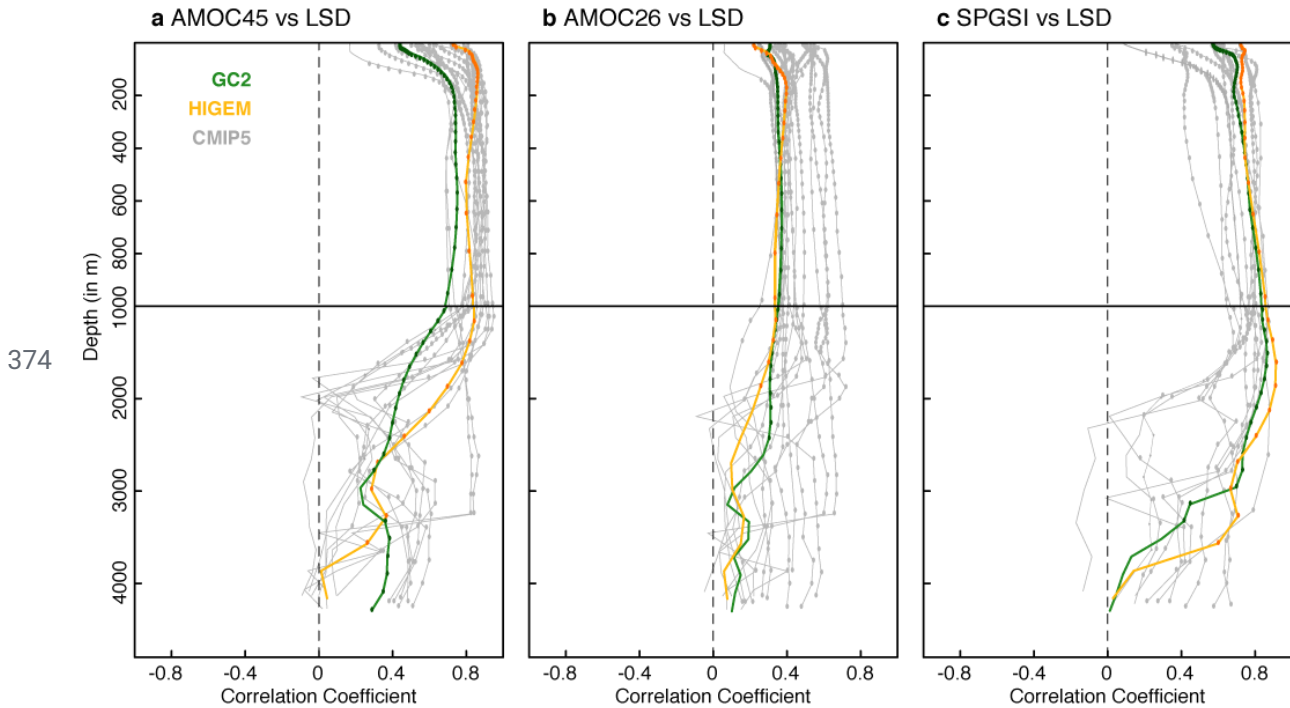
329 [Figure 4a](#) shows that decadal trends in PC1-LSD are associated with trends in the AMOC at 45°N (AMOC45).  
330 Nevertheless, there is some inter-model spread regarding the lag of maximum correlation, which ranges between 0  
331 and 2 years (with PC1-LSD leading), although both variables are in phase for the majority of models. The AMOC at  
332 26°N (AMOC26) is also positively related to PC1-LSD, with PC1-LSD leading AMOC26 by three years on average  
333 ([Figure 4b](#)). However, the average correlation between PC1-LSD and AMOC26 is weaker, and the spread in the  
334 magnitude and lag of the maximum correlation is larger than for AMOC45. Therefore, it appears that the link with  
335 the subtropics is weaker than for 45°N and that AMOC coherence between subpolar latitudes and the subtropics in  
336 coupled models is model dependent. This weaker link of PC1-LSD with the subtropical AMOC is not surprising, as  
337 the LSD anomalies need to propagate a longer distance along the western boundary, allowing for model differences  
338 in the representation of ocean currents and gyres to impact the timing and magnitude of the maximum correlations.  
339 The reasons for the spread in the relationship between PC1-LSD and AMOC26 are explored in Section 4. A strong  
340 relationship is also found between PC1-LSD trends and those in SPGSI ([Figure 4c](#)), of similar order than for  
341 AMOC45. Thus, overall, PC1-LSD is a good proxy for the large-scale ocean circulation in the Subpolar North  
342 Atlantic, and can also be a precursor for a fraction of the AMOC variability in the subtropical Atlantic.

343 PC1-LSD is also a good precursor of the full AMOC variability (i.e. including the Ekman transport), although the  
344 wind-induced fluctuations associated with the Ekman component can introduce differences in the lags of the  
345 maximum AMOC vs PC1-LSD correlations ([Supplementary Figure 2](#)). This different lag can be explained by the  
346 fact that when the Ekman component is included, the AMOC contains a signal that is instantaneously driven by  
347 basin-scale surface wind anomalies (such as the NAO) that are, ultimately, also linked to the heat loss in the  
348 subpolar North Atlantic, which induces a delayed influence on the PC1-LSD (Ortega et al. 2017). Hence, including  
349 Ekman can lead to counterintuitive relationships in some models in which the AMOC leads the PC1-LSD changes.  
350 For that reason, and to ease the interpretation of the lagged-relationships, the rest of the analysis is exclusively  
351 focused on the AMOC indices without ekman.

352 The role of PC1-LSD as a precursor of the AMOC result is further supported by a parallel analysis in [Figure 5](#),  
353 looking at the maximum correlation between the decadal AMOC trends and those in Labrador Sea density as a  
354 function of depth, when the latter leads the AMOC by up to 10 years. [Figure 5](#) reveals that the strongest link  
355 between the Labrador Sea densities and the AMOC, both at 45 and 26°N, occurs in its first 1000 m, the same levels  
356 where the first EOF of LSD show the maximum loadings ([Figure 2a](#)), which confirms the appropriateness of using  
357 PC1-LSD to represent the ocean circulation. The same analysis also supports a strong link between SPGSI and LSD,  
358 although in that case the largest correlations usually happen at deeper levels (between 1000 and 2000 m). Note also  
359 that the main conclusions drawn from PC1-LSD are also valid for the deep LSD index: however, the inter-model  
360 differences are larger in the cross-correlations with the AMOC indices ([Supplementary Figure 3](#)). This difference  
361 could reflect that the deep LSD index is more sensitive to other influences, like the Arctic overflows (Ortega et al.,  
362 2017), which can be very differently represented across models. Overall, the PC1-LSD index seems to be a better  
363 choice to describe multi-decadal North Atlantic variability in multi-model comparisons, as it selects the key depths  
364 for each model. However, PC1-LSD is mostly focused on near surface levels and, therefore, likely represents mostly  
365 Labrador Sea forced variability. Other indices describing densities at deeper levels might be preferable to compare  
366 Labrador Sea Waters of different origins across models, and to evaluate their realism against observations.



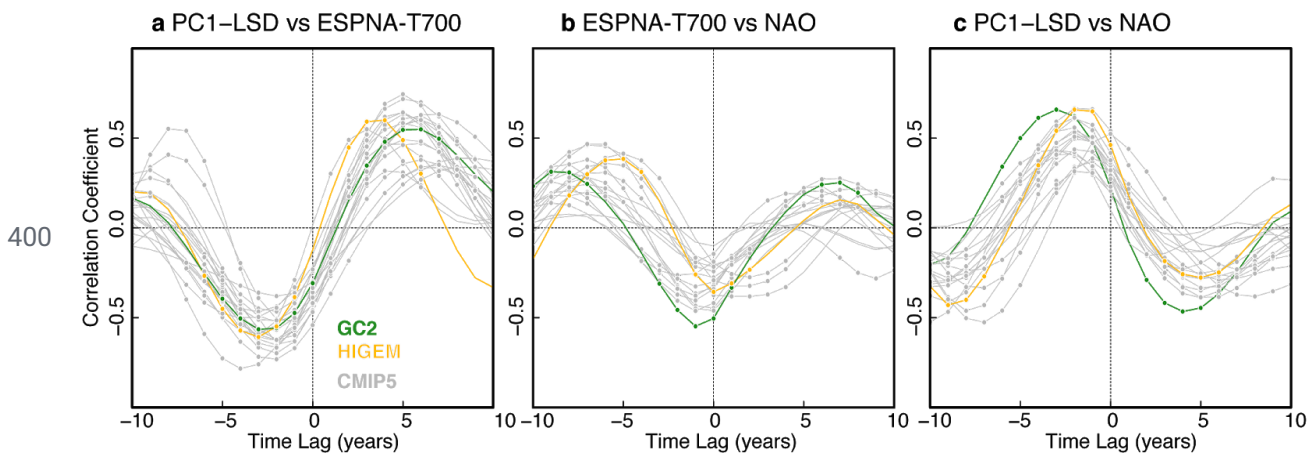
368 **Figure 4:** a Lead-lag correlations across the picontrol ensemble between the PC1-LSD index and the maximum  
 369 AMOC streamfunction at 45°N after the Ekman transport is removed (AMOC45). Correlations are based on 10-year  
 370 running trends. Significance is assessed as in Figure 2d and indicated with a circle. For positive lags, PC1-LSD  
 371 leads. **b-cd** The same as in *a* but between PC1-LSD and the maximum AMOC streamfunction at 26°N after the  
 372 Ekman transport is removed (AMOC26) and the subpolar gyre strength index (SPGSI) and the vertically averaged  
 373 top 700 m temperatures in the eastern subpolar gyre (ESPNA-T700, grey box in Figure 1a).



375 **Figure 5:** a Maximum correlation (for any lag between 0 and 10 years) between the AMOC45 index (after the  
 376 Ekman transport is removed) and Labrador Sea Densities as a function of depth for all the simulations. Colored dots  
 377 indicate correlations that are significant at the 95% confidence level. **b-c** The same as in *a* but between the  
 378 AMOC26 index and LSD, and between the SPGSI and LSD, respectively.



380 Previous studies based on the GC2 picontrol simulation have suggested LSD to be also a potential predictor of  
 381 wide-spread cooling events in the eastern SPNA, like the observed cooling over 2005 to 2014<sup>5</sup> (Robson et al. 2016;  
 382 Ortega et al. 2017). We thereby continue our exploration of the PC1-LSD index by investigating its link with the  
 383 eastern SPNA in the multi-model ensemble. To explore this link we introduce a new index that represents the mean  
 384 potential temperature in the eastern SPNA region (32°W-10°W,47°N-63°N) averaged over the top 700 m of the  
 385 ocean (ESPNA-T700). Lead-lag correlations between the decadal trends in PC1-LSD and this index (Figure 6a)  
 386 show that there is a coherent relationship between both variables across models, with PC1-LSD increases  
 387 (decreases) being consistently followed by ESPNA-T700 warmings (coolings). Nevertheless, there are inter-model  
 388 differences concerning the magnitude and lag of the strongest positive correlations, revealing important uncertainty  
 389 in the relationship. The spread in the PC1-LSD vs ESPNA-T700 relationship is thus reminiscent of the spread found  
 390 between PC1-LSD and AMOC26, which suggests that they might be related. We also note significant negative  
 391 correlations when ESPNA-T700 leads PC1-LSD by 2-4 years that might be explained by the opposed (and nearly  
 392 concomitant) impacts that the NAO exerts on both variables (Figure 6b,c). Positive NAO phases, and associated  
 393 surface buoyancy forcing (Lozier et al., 2008) lead in first instance to negative SSTs (Barrier et al., 2014; Lohmann  
 394 et al., 2009) and an almost simultaneous cooling in ESPNA-T700 (Figure 6b). In comparison, on the western side  
 395 of the SPNA, positive NAO phases contribute to reduce vertical density stratification, favoring convection and a  
 396 more positive LSD index (Robson et al., 2016), which in the models lags the NAO by 2-3 years (Figure 6c). The  
 397 fact that correlations between NAO and ESPNA-T700 are weaker than between PC1-LSD and ESPNA-T700  
 398 suggests that the ocean might also be playing an additional role (besides the NAO) in controlling the ESPNA  
 399 temperatures.

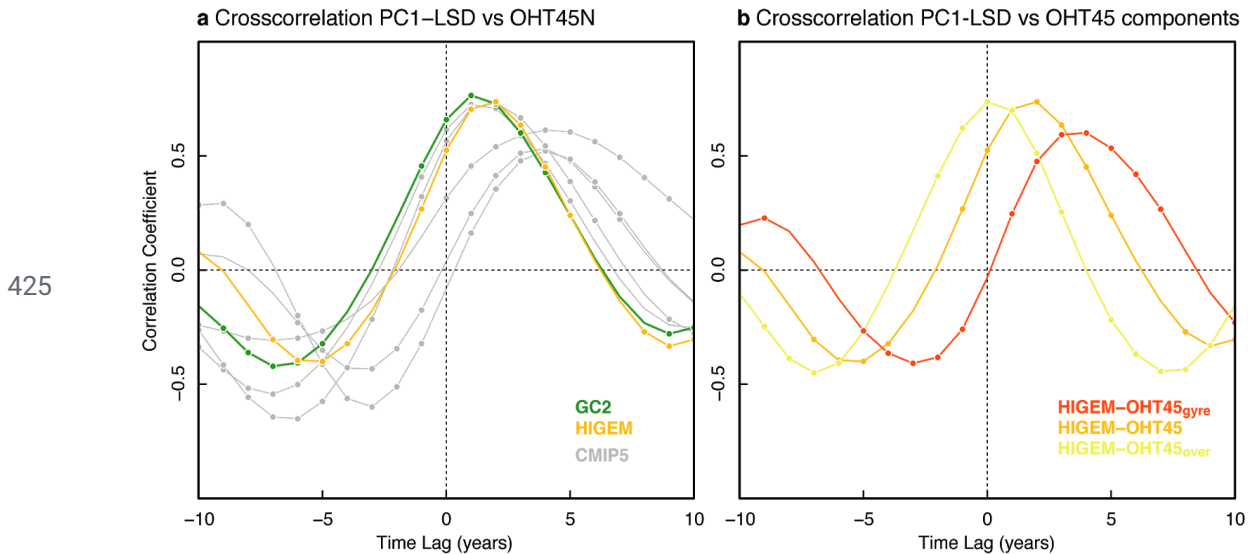


401 **Figure 6:** a Lead-lag correlations across the picontrol ensemble between the PC1-LSD index and the vertically  
 402 averaged top 700 m temperatures in the eastern subpolar gyre (ESPNA-T700; grey box in Figure 1a). Correlations  
 403 are based on 10-year running trends. Significance is assessed as in Figure 2d and indicated with a circle. For positive  
 404 lags, PC1-LSD leads. b-c The same as in a but between the North Atlantic Oscillation (NAO; defined as the  
 405 standardised difference in sea level pressure between the closest grid-points to Azores and Reykjavik) and the  
 406 ESPNA-T700, and between the NAO and the PC1-LSD, respectively. In these two cases, for negative lags the NAO  
 407 leads.

408 The link between PC1-LSD and the ESPNA could be explained through an influence of the PC1-LSD<sub>first</sub> on the  
 409 meridional ocean heat transport. This link is now investigated in the two eddy-permitting simulations (Figure 7) and  
 410 in the five CMIP5 models for which the ocean heat transport fields are publicly available. In the two high resolution



411 experiments and two of the CMIP5 ones the decadal trends in the meridional ocean heat transport at 45°N (OHT45)  
 412 are strongly linked with those in PC1-LSD. This is a similar relationship to the one previously found in [Figure 4](#)  
 413 between PC1-LSD and both the AMOC45 and SPGSI, but in this case with PC1-LSD leading with slightly longer  
 414 lead time. The other CMIP5 experiments support a weaker, yet significant, link, as well as a longer lag between  
 415 OHT45 and PC1-LSD. Altogether, [Figure 7a](#) confirms that PC1-LSD is a good precursor of the changes in the  
 416 meridional ocean heat transport, although with some differences across models which might reflect a different  
 417 representation of certain processes. The contributions of two different processes to this delay are further investigated  
 418 in HiGEM, for which OHT had been decomposed online at each time-step into vertical and horizontal heat  
 419 transports (as in Bryan, 1969), which can be respectively interpreted as the “overturning” (i.e. characterised by the  
 420 zonal mean transport) and “gyre” (i.e. characterised by variations from the zonal mean transport) components  
 421 (Robson et al., 2018a). While the overturning contribution (OHT45<sub>over</sub>) increases in phase with the AMOC45,  
 422 SPGSI and PC1-LSD changes ([Figure 7b](#)), the increase in the gyre component (OHC<sub>gyre</sub>) starts four years later. That  
 423 lag could be the time required in HiGEM for the propagation of mean and/or anomalous temperatures from the  
 424 southern to the northern branch of the SPG.



426 **Figure 7: a** Lead-lag correlations in a subset of the picontrol experiments between the PC1-LSD index and the  
 427 ocean heat transport across the 45°N transect (OHT45N). Note that the ocean heat content is only available for 5  
 428 models of the CMIP5 ensemble. Correlations are based on 10-year running trends. **b** The same as in *a* but only in  
 429 HiGEM for the different terms of the OHT45N. For positive lags, PC1-LSD leads.

430

#### 431 4. Characteristics of the inter-model spread in subpolar to subtropical AMOC

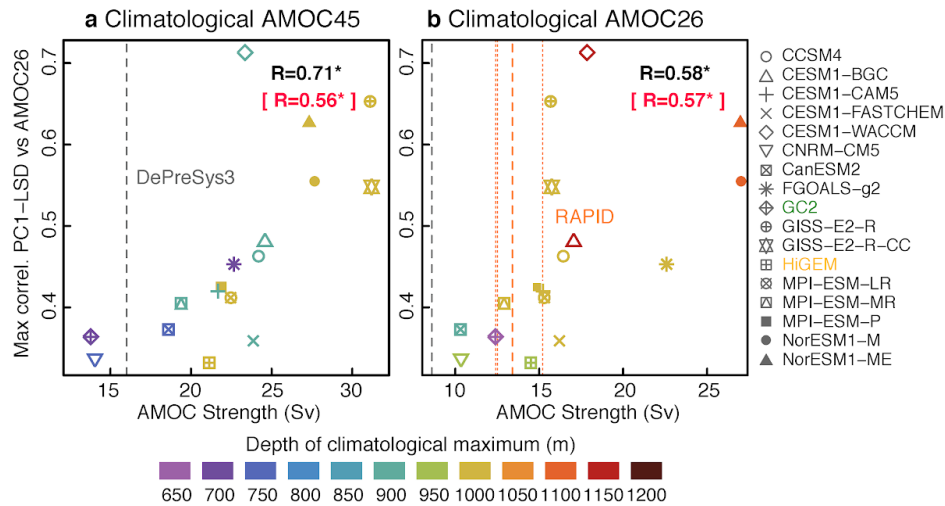
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433 This section investigates which particular climatological model features are linked to the large inter-model spread in  
 434 the PC1-LSD vs AMOC26 relationships. The most relevant model features thus identified will improve our process  
 435 understanding, and can eventually be used to identify which models are most realistic and, in turn, can deliver more  
 436 reliable projections of the future changes in the North Atlantic.

437 [Figure 8](#) shows that models that simulate a stronger and deeper climatological AMOC (both at 45°N and 26°N) tend  
 438 to have a stronger correlation between PC1-LSD and the subtropics. All these linear relationships between  
 439 climatological AMOC strength and depth and the PC1-LSD vs AMOC26 connectivity are significant at the 95%

440 confidence level. These climatological AMOC values (without Ekman) can be put in context with those from  
 441 RAPID observations and DePreSys3. RAPID observational uncertainties have been considered by including the  
 442 mean values over three different non-overlapping periods (i.e. 2004-2007, 2008-2012 and 2013-2016; dotted lines in  
 443 Figure 8). The scatterplots show that the majority of models whose climatological AMOC26 lies within the  
 444 RAPID/DePreSys3 climatological spread have a relatively weak link between PC1-LSD and AMOC26, although  
 445 some models supporting a strong link are also included or remain close to the RAPID/DePreSys3 values. However,  
 446 caution is recommended, e.g., before defining emerging constraints, because model and observations are not directly  
 447 comparable for numerous reasons. For example, both RAPID and DePreSys3 cover shorter periods than the  
 448 simulations and relate to different background forcing conditions (present day vs preindustrial) which might imply  
 449 different mean states (Thornalley et al., 2018). Also, climatological values of the AMOC26 strength are notably  
 450 weaker in DePreSys3 than in RAPID, a difference that is not explained by the different temporal periods covered by  
 451 each dataset (not shown) and that implies that DePreSys3 might be underestimating too the real AMOC45 strength.  
 452 This underestimation might be larger than shown in Figure 8, as Furthermore, evidence suggests RAPID  
 453 calculations from mooring arrays might be underestimating the AMOC strength by  $\sim 1.5$  Sv (Sinha et al., 2018).

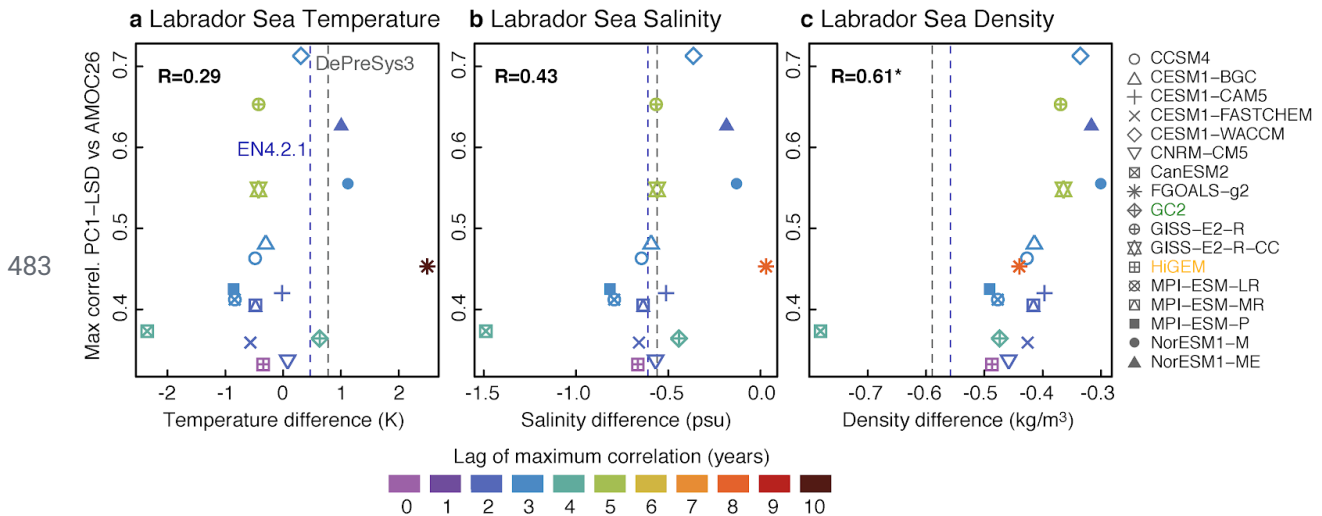
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455 **Figure 8: a-b** Scatterplot of the maximum cross-correlation value in Figure 4b ~~s~~ at any lag between PC1-LSD and  
 456 AMOC26 ~~N~~ against the climatological AMOC45 ~~N~~ and AMOC26 ~~N~~ means, respectively. All AMOC indices refer to  
 457 the values after the Ekman transport signal is removed. The maximum correlations are based on 10-year running  
 458 trends, and always happen when PC1-LSD leads the AMOC26 index. Colors indicate the depth at which the  
 459 climatological AMOC maximum occurs. The correlation coefficient between the maximum PC1-LSD correlation  
 460 and the climatological mean AMOC is shown in the top-left corner in black. In magenta, the analogous correlation  
 461 but against the depth of the mean climatological AMOC is shown. The presence of an asterisk indicates that the  
 462 correlation is significant at the 95% confidence level. The dashed grey vertical lines mark the climatological AMOC  
 463 strength value in the DePreSys3 assimilation run. The orange vertical lines indicate the climatological value from  
 464 RAPID observations (Smeed et al., 2018) from 2004 to 2016 (dashed), and in three non-overlapping sub-periods of  
 465 4 years (dotted).

466 A potentially important factor behind the inter-model spread in Figure 4b is the mean density stratification in the  
 467 Labrador Sea. Figure 9 suggests that, indeed, the PC1-LSD vs AMOC26 spread is partly influenced by the density  
 468 stratification in this region. ~~Models with a weaker density stratification have a stronger relationship between~~  
 469 ~~PC1-LSD and AMOC26. However, no link is found between the PC1-LSD vs AMOC26 relationship and~~  
 470 ~~temperature or salinity stratification in the Labrador Sea.~~ Models that have a weaker density stratification (here  
 471 defined as the difference between the top 100 m, and the average between 500-1000 m), and thus favor deeper

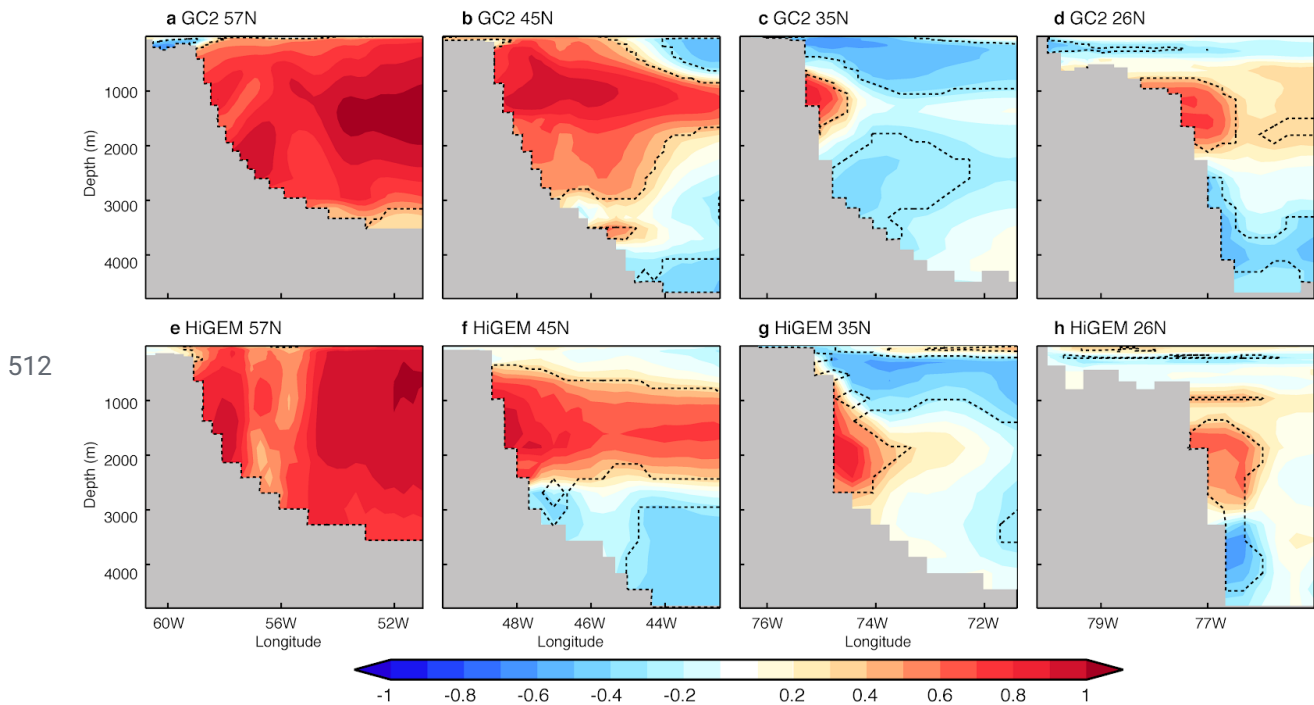
472 convection in the Labrador Sea, generally exhibit a stronger link between PC1-LSD and AMOC26. This result is  
 473 robust for other stratification indices based on different depth levels (See [Supplementary Figure 4](#)). Differences in  
 474 density stratification across-models can be due to a combination of different factors, from differences in the local  
 475 buoyancy fluxes (driven by differences in the atmospheric circulation), to differences in the representation of the  
 476 Arctic overflows, which are parameterised in some models (e.g. the CESM family; Danabasoglu et al., 2010), and  
 477 explicitly resolved in others. No robust link between the PC1-LSD vs AMOC relationship and both temperature and  
 478 salinity stratification in the Labrador Sea has been found. ~~is weaker than for density for all depths considered, to  
 479 compute the stratification index, and is generally not significant.~~ It is also worth mentioning that all models except  
 480 CanESM2 are more weakly stratified in the Labrador Sea than the observations (represented herein by the  
 481 DePreSys3 assimilation run and EN4.2.1). Hence, the real link of LSD with the AMOC26 may not be as strong as  
 482 some models suggest.



484 **Figure 9:** a Scatterplot of the maximum ~~cross-correlation value in Figure 4b~~ ~~at any lag~~ between PC1-LSD and  
 485 AMOC26 (without the Ekman component) against the climatological mean of the Labrador Sea temperature  
 486 stratification index (computed as the difference of the vertical means in the levels 0-100 m minus the vertical means  
 487 in the levels 500-1000 m; see Fig. 1). The maximum correlations are based on 10-year running trends. The  
 488 correlation coefficient between the two metrics is shown in the top-left corner. The presence of an asterisk indicates  
 489 that the correlation is significant at the 95% confidence level. Colors indicate the lag at which the maximum  
 490 correlation between PC1-LSD and AMOC26 is obtained. The grey (bluegreen) vertical lines depict the mean  
 491 stratification value in the DePreSys3 assimilation run (EN4.2.1). In both cases, their overlap period is used to  
 492 compute the climatology (i.e., 1960-2013). **b-c** The same as in **a** but for the Labrador Sea salinity and density  
 493 (defined as  $\sigma_2$ ), respectively.

494 Another key aspect of the PC1-LSD vs AMOC26 connectivity is the western boundary density (WBD). Indeed,  
 495 boundary density is critical to the mechanism through which LSD influences the AMOC at lower latitudes. Positive  
 496 (negative) LSD anomalies propagate equatorward following this boundary, and as they do so they strengthen  
 497 (weaken) the zonal density gradient, triggering a thermal wind response that accelerates (decelerates) the AMOC. In  
 498 the following we investigate differences in the propagation of boundary densities across models, and if these  
 499 differences can affect the inter-model PC1-LSD vs AMOC26 spread. [Figure 10](#) focuses on the two high-resolution  
 500 simulations, where important differences already manifest. It represents the in-phase correlations of PC1-LSD with  
 501 the density fields (defined as  $\sigma_2$ ) near the western boundary at four different longitudinal transects: 57°N (cutting  
 502 across the Labrador Sea), 45°N, 35°N and 26°N. In both models, the depth of the maximum correlation near the

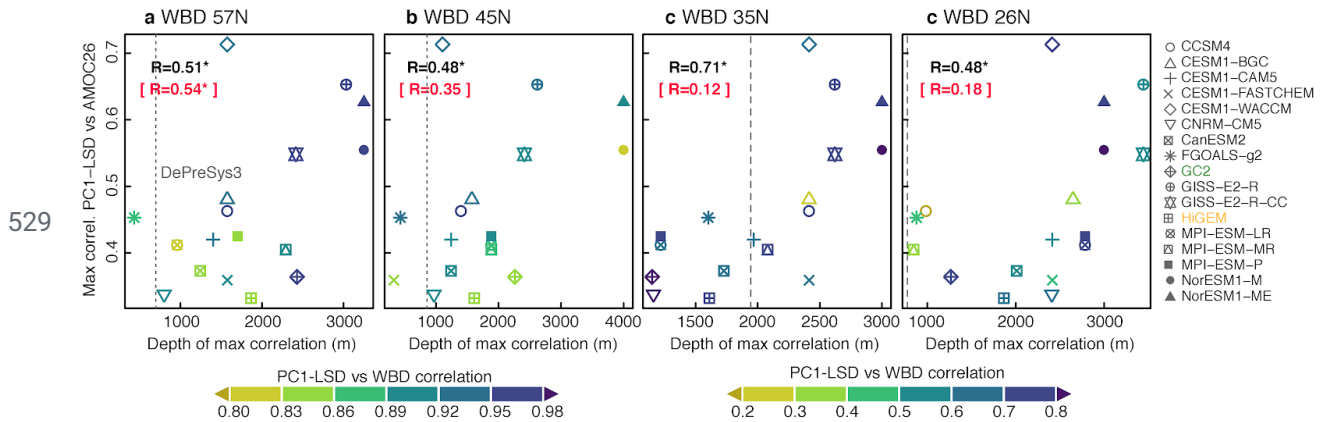
503 continental shelf is coherent across latitudes. However, in HiGEM these occur at deeper levels (1000 to 3000 m)  
 504 compared to GC2 (1000 to 2000 m), and the difference is especially clear at 35°N, where the highest correlations  
 505 occur at ~2000 m in HiGEM, while only at 1000 m in GC2. Similar depth differences are also found at 26°N, but  
 506 with slightly weaker correlations. In addition to the difference in the depth of the maximum correlation between  
 507 HiGEM and GC2, there are differences in the vertical structure between the two models. For example, at 35°N in  
 508 GC2, density anomalies on the western boundary form a tripole (low correlation above and below the maximum  
 509 correlation at ~1000 m), but in HiGEM the density anomalies form a dipole (Figure 10g). We note some differences  
 510 in bathymetry at this latitude (which is steeper in HiGEM), which might partly explain some of the differences in  
 511 terms of the density correlation structure.



513 **Figure 10:** a In-phase correlation in GC2 between the PC1-LSD index and the density fields across a zonal section  
 514 at 57°N located in the vicinity of the western Atlantic boundary. Thin dashed contours enclose areas where the  
 515 correlation significance exceeds the 95% confidence level. Correlations are based on 10-year running trends. **b-d**  
 516 The same as in *a* but for zonal sections at 45°N, 35°N and 26°N. **d-h** The same as in *a-d* but for HiGEM.

517 [Figure 11](#) shows that the diversity in the depth of these boundary densities is even more evident when including the  
 518 CMIP5 models. The depth of the maximum correlation between PC1-LSD and the western boundary density at the  
 519 four latitudinal sections relates linearly (and significantly at the 95% confidence level) across models with their  
 520 PC1-LSD vs AMOC26 correlation. In this case, models exhibiting maximum correlations with the deeper WBDs at  
 521 deeper levels generally show stronger links between PC1-LSD and the subtropical AMOC. In DePreSys, our  
 522 observationally-constrained reference (dashed grey lines in [Figure 11](#)), these maximum correlations tend to occur at  
 523 relatively shallow levels when compared with the multi-model ensemble. We have also checked if models with  
 524 stronger correlations without the WBDs (as represented by the PC1-LSD and WBD maximum correlations at every  
 525 latitudinal section) also support a stronger link between the PC1-LSD and the AMOC (correlations in magenta in  
 526 [Figure 11](#)), but this linearity assumption only holds true at 57°N (correlations in magenta in [Figure 11](#)). This

527 suggests that the depth along which WBDs propagate southward, and/or the vertical structure of anomalies, are the  
 528 key aspects to understand and potentially narrow down the spread.



530 **Figure 11:** a Scatterplot of the maximum cross-correlations value in Figure 4 at any lag between PC1-LSD and  
 531 AMOC26N (without the Ekman component) against the depth at which the maximum correlations at any lag  
 532 between PC1-LSD and the WBD at 57°N occur. The maximum correlations are based on 10-year running trends.  
 533 The correlation coefficient between the two metrics is shown in black the top-left corner. Likewise, another  
 534 correlation coefficient in magenta is shown, computed between the PC1-LSD and AMOC26N maximum correlation  
 535 and the PC1-LSD and WBD at 57°N maximum correlation. The presence of an asterisk indicates that the  
 536 correlation is significant at the 95% confidence level. Colors indicate the maximum correlation between PC1-LSD  
 537 and the WBD. The grey vertical lines depict the corresponding depth of maximum correlation for the DePreSys3  
 538 assimilation run. **b-d** The same as in **a** but for the WBD at 45, 35 and 26°N, respectively.

### 539 5. Conclusions and discussion

540

541 This article has explored, in a multi-model context, the linkages between subsurface density in the subpolar North  
 542 Atlantic (SPNA) and the ocean circulation further south. In particular, it has explored the role of Labrador Sea  
 543 density (LSD) in driving Western Boundary Density anomalies (WBD) and the ocean circulation, and the impact on  
 544 upper ocean temperature changes in the SPNA. The analysis was based on two control simulations with  
 545 eddy-permitting models, a preindustrial one with HadGEM3-GC2 and a present day one with HiGEM, and on 20  
 546 CMIP5 preindustrial experiments. Furthermore, where possible these characteristic model features have been  
 547 computed in observational datasets, as well as in a simulation assimilating observations. The major findings are  
 548 listed below:

549

- 550 - All the simulations show clear multidecadal variability in Labrador Sea density. There is also a close link  
 551 between LSD and the strength of the subpolar Atlantic Ocean circulation, with positive density anomalies  
 552 leading to a strengthening of the Atlantic Meridional Overturning Circulation (AMOC) at 45°N and the  
 553 Subpolar Gyre (SPG) circulation.
- 554 - The relationship between anomalous LSD and the strength of the AMOC at 26°N - the latitude of the  
 555 RAPID array measurements - is also positive in the simulations, but there are significant inter-model  
 556 differences, both in the strength of the relationship and the lag of maximum correlation. This uncertainty  
 557 implies that the connectivity of LSD with the subtropics and latitudinal AMOC coherence is  
 558 model-dependent.
- 559 - The connectivity between anomalies in LSD and AMOC at 26°N is sensitive to different model features,  
 560 including the strength and depth of the climatological AMOC maximum, the mean density stratification in



561 the Labrador Sea, and the depths at which the LSD propagates southward along the western boundary.  
562 Stronger LSD connectivity with the subtropics tends to occur in models with a stronger and deeper AMOC,  
563 weaker Labrador Sea stratification and western boundary density propagating at deeper levels.  
564 - Observationally derived constraints of the model based relationships tend to suggest that the link between  
565 LSD and the subtropical AMOC is weak. This suggests that observations of AMOC via RAPID may not be  
566 representative of the basin wide buoyancy forced AMOC variability. However, caution is advised because  
567 simulations and observations are not directly comparable, and so significant uncertainty remains in  
568 constraining the relationship between LSD and subtropical AMOC.  
569 - The multi-model ensemble does also support a significant lagged relationship between LSD and the upper  
570 ocean temperature in the eastern SPNA, in line with previous studies linking LSD to the recently observed  
571 changes in the North Atlantic. However, models disagree regarding the strength of the link (correlations  
572 between 0.3-0.7), and the maximum lag (3 to 10 years).

573  
574 We have shown that, in coupled climate models at least, subsurface density anomalies in the western SPNA are an  
575 important predictor of the wider North Atlantic ocean circulation and upper ocean temperature in the SPNA. This  
576 importance on the ocean circulation is especially clear at the latitudes of the SPNA itself. Given the important role  
577 of the wind in driving lower latitude AMOC anomalies, and the range of processes by which wind can act on the  
578 AMOC (Duchez et al., 2014b, 2014a; Kanzow et al., 2010; Polo et al., 2014; Zhao and Johns, 2014) it is not  
579 surprising that the relationship between LSD and AMOC at 26°N is much weaker. Nevertheless, the reasons behind  
580 the large spread in these relationships across models is not so clear.

581 We have tried to constrain this uncertainty by looking at a range of observed metrics that may explain the spread in  
582 the correlation strength, including the density anomalies on the western boundary, the stratification of the Labrador  
583 Sea, and the mean-AMOC strength. Overall, these constraints point to a relatively weak relationship between LSD  
584 and AMOC at 26°N on decadal timescales (i.e.  $r \sim 0.4$ ) in the real world. However, there are many reasons why this  
585 number is still very uncertain, and further work is needed to assess its validity. A caveat of this study is that the  
586 simulations and observation-based datasets employed are not directly comparable, as they differ in the background  
587 radiative forcing levels, the length of the period used to compute the climatologies, and even the way some indices,  
588 like the AMOC, are computed. We also recognise that there is large uncertainty within the observationally derived  
589 metrics. For instance, the assimilation run in DePreSys3, which is used to constrain relationships, clearly  
590 underestimates the mean AMOC strength at 26°N with respect to RAPID (see Figure 8b) and, therefore, might be  
591 also underestimating the AMOC at higher latitudes. Our findings might also be limited by model deficiencies. There  
592 is also emerging evidence that current models underestimate AMOC variability and North Atlantic variability on  
593 decadal timescales (Roberts et al., 2013; Cheung et al., 2017; Yan et al., 2018), which can degrade decadal  
594 predictability in the region and even lead to overly weak linkages between the AMOC and the AMV (Yan et al.,  
595 2018). The AMV is indeed a mode of variability that also shows important differences across models, in different  
596 aspects like its periodicity, amplitude, spatial structure and climate footprints (Medhaug and Furevik, 2011; Zhang  
597 and Wang 2013; Kavvada et al., 2013), inter-model differences that could be partly connected with those herein  
598 reported for the PC1-LSD vs AMOC relationships. We also recognise that there is large uncertainty within the  
599 observationally derived metrics. For example, the assimilation run in DePreSys3, which is used to constrain  
600 relationships, clearly does not represent the mean AMOC strength at 26°N from RAPID (see Figure 8b). Models  
601 also generally tend to generally underestimate the depth of the return flow, and this may still affect how density  
602 anomalies project on the basin-wide AMOC. It has also been argued that ocean-only models produce too much deep  
603 water in the western basin and Labrador Sea (i.e., Li et al., 2019), and recent observations even challenge the  
604 prevailing view from models that Labrador Sea convection dominates the AMOC variability (Koenigk and Brodeau,  
605 2017), suggesting that the key deep water formation occurs in the Irminger Sea, a few hundred kilometers north

606 eEast of the Labrador Sea (Lozier et al., 2019). Therefore, further in-depth study is warrantedrequired to narrow  
607 down understand the uncertainty in the real AMOC and PC1-LSD relationships.

608 Most of the models considered in this study have relatively coarse resolution, including non-eddy oceans ( $\geq$   
609  $1^\circ \times 1^\circ$ ), which means that they might be missing some key dynamics for the AMOC (Johnson et al., 2019) that could  
610 be important to represent realistic linkages. The current analysis also, however, includes two models at  
611 eddy-permitting resolution (HadGEM3-GC2 and HiGEM), whose relationships lie within the spread of those in the  
612 coarser models, suggesting that climatological features (like Labrador Sea stratification) can be more important than  
613 the representation of mesoscale processes. However, it could be that higher resolution is needed (e.g. enabling  
614 meso-scale eddies in subpolar latitudes) to identify substantial differences (Hirschi et al., 2020; Johnson et al.,  
615 2019). A recent analysis based on HadGEM3-GC3.1 (a later version of HadGEM3-GC2) configured at different  
616 horizontal resolutions has shown that long-standing model biases, affecting including in the North Atlantic, are  
617 reduced at eddy-resolving resolution ( $1/12^\circ \times 1/12^\circ$  in the ocean), and that the strength of the AMOC, the boundary  
618 currents and the northward heat transport is higher than for the coarser resolutions (Hirschi et al., 2020; Roberts et  
619 al., 2019). High resolution coupled models also generally support the new view from OSNAP observations in which  
620 the largest fraction of AMOC variability (on sub annual to decadal timescales) originates at the eastern SPNA  
621 (Hirschi et al., 2020). Eddy-resolving resolutions have also been shown in a multi-model study (Roberts et al., 2020)  
622 to represent the AMOC response at  $26^\circ\text{N}$  differently in future projections, leading to stronger declines than in  
623 non-eddy simulations, declines mostly associated with a weakening in the Florida Current. Roberts et al. (2020)  
624 also compares the meridional coherence of the AMOC, which does not seem to be resolution-dependent, a result  
625 that is in line with another multi-model comparison between non-eddy and eddy-permitting simulations (Li et al.,  
626 2019).

627 Despite the current limitations in the models considered for this study, it is important to highlight that they provide a  
628 rather consistent picture of a chain of relationships in the North Atlantic that is able to explain some of the recent  
629 observed trends (Robson et al. 2016). This paper has broadly characterized this behaviour, and highlighted the  
630 uncertainty. These relationships are also consistent with the mechanisms proposed by Yeager and Robson (2017) to  
631 explain high levels of predictive skill in the SPNA on decadal timescales. Our analysis has also helped to identify  
632 specific metrics (such as LSD stratification and the depth of the boundary density) that could be used as emergent  
633 constraints for future projections, i.e. to subset the simulations expected to more realistically represent the future  
634 changes in the region. Having a more realistic subpolar gyre stratification at present day conditions has been shown  
635 in CMIP5 simulations to increase the probability of a future collapse in convection (Sgubin et al., 2017), that would  
636 lead to a widespread SPG cooling. It remains to be tested if similar conclusions can be drawn from eddy-resolving  
637 simulations.

638 **Code availability.** The main scripts used in the analysis and other supporting information that may be useful to  
639 reproduce the results of this article are archived at the Barcelona Supercomputing Center and will be shared upon  
640 request by the corresponding author.

641 **Data availability.** Outputs from the CMIP5 simulations can be downloaded from the corresponding ESGF node:  
642 <https://esgf-node.llnl.gov/projects/cmip5/>. EN4 observations used in this study correspond to version 2.1 of the  
643 dataset, available at <https://www.metoffice.gov.uk/hadobs/en4/download-en4-2-1.html>. Outputs from the GC2,  
644 HiGEM and DePreSys3 simulations are available upon request to the corresponding author.

645 **Author contributions.** P. O., J. R. and R. S. conceived the study, which was later discussed and refined with the  
646 other co-authors. M. M. downloaded and processed the CMIP5 data, computing the main climate indices. P. O. led  
647 the analysis, and together with J. R. prepared the manuscript with contributions from all co-authors.

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