



1	N	orth Pacific subtropical sea surface temperature frontogenesis and
2		its connection with the atmosphere above
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ABSTRACT

The frontogenesis of the North Pacific subtropical sea surface temperature front 14 (NPSTF) occurring from October to the following February is examined 15 quantitatively based on the mixed-layer energy budget equation, with a focus on its 16 17 connection with the atmosphere above. Diagnosis results show that the net heat flux dominates the frontogenesis from October to December, while the meridional 18 19 temperature advection in the ocean contributes equally as or even more than the net 20 heat flux in January and February. The atmosphere is critical to the frontogenesis of 21 the NPSTF, including the direct effect of the net heat flux and the indirect effect through the Aleutian low. Further analyses demonstrate that the latent heat flux (the 22 shortwave radiation) dominates the net heat flux in October (from November to 23 24 February). The meridional temperature advection in the ocean is mostly owing to the 25 meridional Ekman convergence, which is related to the Aleutian low. Climatologically, the strengthening and southward migration of the Aleutian low from October to the 26 following February are characterized by the acceleration and southward shift of the 27 28 westerly wind to the south, respectively, which can drive southward ocean currents. Correspondingly, the southward ocean currents give the colder meridional advection 29 to the north of the NPSTF in January and February, favoring the frontogenesis. In 30 addition, the Aleutian low plays a role in transforming the dominant effect of the net 31 32 heat flux to the joint effect of the meridional temperature advection and the net heat 33 flux in January. CESM1.0.3 model with a slab ocean model further confirms the important influence of the atmosphere on the frontogenesis and on the meridional 34





- 35 temperature advection.
- 36 Key words: North Pacific subtropical sea surface temperature front; frontogenesis;
- 37 net heat flux; meridional temperature advection; Aleutian low





38 **1. Introduction**

39 The North Pacific Ocean is featured by two zonal sea surface temperature (SST) fronts at mid-latitude and subtropics, respectively. The mid-latitude front, with greater 40 magnitude, is referred to as the North Pacific subarctic SST front (NPSAF), and the 41 42 subtropical one is the North Pacific subtropical SST front (NPSTF). Due to the smaller magnitude, the NPSTF has been rarely studied. However, it also exerts 43 44 significant influences on the overlying atmosphere (Xie, 2004; Kobashi et al., 2008; 45 Wang et al., 2016; Zhang et al., 2017, 2018). On the synoptic scale, Kobashi et al. 46 (2008) found that the subsynoptic lows along the NPSTF are enhanced by the condensational heating and baroclinicity associated with the NPSTF during April to 47 May. On the interannual scale, the intensified NPSTF in spring can not only 48 49 accelerate the East Asian westerly jet (Zhang et al., 2017), but also serve as a precursor to the following La Niña event (Zhang et al., 2018). 50

From the respective of the seasonal variation, the NPSAF can exist throughout the 51 year, but the NPSTF is robust in winter and spring and is absent in summer and 52 53 autumn (Fig. 1; Kobashi and Xie, 2012). Thus, several studies have focused on the frontogenesis and frontolysis of the NPSTF (Qiu and Kawamura, 2012; Qiu et al., 54 2014; Roden, 1975; Kazmin and Rienecker, 1996). It is pointed out that the net heat 55 flux is responsible for the frontolysis of the NPSTF (Qiu and Kawamura, 2012; Qiu et 56 57 al., 2014). In terms of the frontogenesis, Roden (1975) found the meridional Ekman convergence is the primary reason for the frontogenesis of the NPSTF. However, 58 Kazmin and Rienecker (1996) diagnosed the mixed-layer energy budget equation 59





using the observation data from 1982 to 1990, and pointed out that both the net heat 60 61 flux and the Ekman convergence are frontogenetic and equally important to provide the observed frontogenesis in winter, rather than the Ekman convergence alone. This 62 finding is further confirmed by Dinniman and Rienecker (1999), based on the 10 63 64 years' (1985-1995) simulation of a primitive equation model (Geophysical Fluid Dynamics Laboratory's MOM2). However, they argued that these two factors are not 65 66 equally important: the net heat flux (the Ekman convergence) dominates the 67 frontogenesis in the western subtropical Pacific (the central and eastern subtropical 68 Pacific). Thus, the relative role of the net heat flux and the Ekman convergence in the frontogenesis of the NPSTF remains unclear, due to limited data used in previous 69 70 studies. Meanwhile, the net heat flux is associated with the air-sea interaction, and the 71 Ekman convergence is driven by the surface wind stress, implying that both frontogenesis factors are closely related to the atmospheric circulation. Kazmin (2017) 72 demonstrated the long-term (quasi-decadal) variability of the subtropical SST front is 73 determined by the variability of the meridional shear of the zonal wind. Thus, the role 74 75 of the atmosphere in the frontogenesis of the NPSTF deserves to further study.

Therefore, this paper aims to figure out the relative importance of the net heat flux and Ekman convergence in the frontogenesis of the NPSTF, especially the role of the atmosphere in this process. The rest of the paper is organized as follows. We introduce the data and methods in Section 2. We analyze the frontogenesis of the NPSTF using the mixed-layer energy budget equation in Section 3 to explore the relative importance of the net heat flux and the oceanic meridional temperature advection





- 82 (including the Ekman convergence). Section 4 further investigates the roles of
- 83 atmosphere in the frontogenesis. Summary is given in Section 5.
- 84
- 85 2 Data and Methods
- 86 2.1 Data

We use the monthly ocean temperature, current velocities and wind stress from the 87 88 Simple Ocean Data Assimilation (SODA; Carton and Giese, 2008) version 2.2.4 at 89 $0.5^{\circ} \times 0.5^{\circ}$ grid with 40 levels from the depth of 5 to 2000 m. We also use surface heat 90 fluxes from the Objectively Analyzed Air-sea Fluxes Project (OAFlux; Yu and Weller, 2007) at 2.5°×2.5° grid to examine the mixed-layer energy budget. All heat fluxes are 91 defined to be positive downward. For consistency, all variables are interpolated onto 92 93 $0.5^{\circ} \times 0.5^{\circ}$ grid, and they cover the period from January 1984 to December 2009. The ocean temperature at $1.0^{\circ} \times 1.0^{\circ}$ grid with 27 levels from the International Pacific 94 Research Center (IPRC) Argo Product, together with ocean currents (on 40 levels) 95 and surface heat fluxes at 0.3°×1.0° grid from the NCEP Global Ocean Data 96 Assimilation System (GODAS; Saha et al., 2011) are used to confirm our results 97 based on the SODA data. These data are interpolated onto 1.0°×1.0° grid at 27 depths, 98 and only cover the period from January 2005 to December 2013. 99

The atmospheric data used in this study are monthly ERA-interim reanalysis from the European Center for Medium-range Weather Forecasts (ECMWF; Dee et al., 2011), including geopotential height and winds. They are on $1.5^{\circ} \times 1.5^{\circ}$ grid, and cover the period from January 1984 to December 2009.





104 2.2 The mixed-layer energy budget equation

The temporal variation of SST is governed by mixed-layer dynamics, which can
be represented by the mixed-layer energy budget equation (Dinniman and Rienecker,

107 1999; Zhang et al., 2013):

108
$$\frac{\partial SST}{\partial t} = -u \frac{\partial SST}{\partial x} - v \frac{\partial SST}{\partial y} - w \frac{\Delta T}{H} + \frac{Q_{net}}{\rho_0 c_p H} + R, \qquad (1)$$

109 where SST denotes sea surface temperature (here, we assume that SST equals 110 mixed-layer mean temperature), and ΔT represents the temperature difference 111 between the mixed layer and the interior ocean immediately below the mixed layer. u and v are mixed-layer zonal and meridional oceanic current velocities, respectively; w 112 113 is the vertical velocity at the bottom of the mixed layer. H is mixed-layer depth. Q_{net} is the net surface heat flux, including sensible and latent heat fluxes, as well as 114 115 longwave and shortwave radiation. A positive value of Q_{net} means that the ocean gains heat from the atmosphere. ρ_0 and c_p are the density and heat capacity of sea 116 water, respectively. R is the residual term, including sub-grid scale processes and 117 118 dissipation. The zonal temperature advection $(-u\partial SST/\partial y)$, meridional temperature 119 advection $(-v\partial SST/\partial y)$ and vertical temperature advection $(-w\Delta T/H)$ are intrinsic 120 processes in the ocean (Yu and Boer, 2004; Chen et al., 2014), while the net heat flux term $(Q_{net}/\rho_0 c_n H)$ represents air-sea interaction. The SST tendency $(\partial SST/\partial t)$ in a 121 particular month is obtained through the central finite difference. 122

123 Since the meridional gradient of SST overwhelmingly dominates over its zonal 124 counterpart in the frontal region, the gradient magnitude (GM) of the NPSTF is





125 defined as $GM = -\partial SST/\partial y$ to measure the intensity of the NPSTF in a particular 126 month (Qiu and Kawamura, 2012; Qiu et al., 2014). Accordingly, GM is always 127 positive because the climatological mean SST is higher in the south. Its tendency can 128 be derived from Eq. (1) as follows,

129
$$\frac{\partial GM}{\partial t} = \frac{\partial}{\partial y} \left(u \frac{\partial SST}{\partial x} \right) + \frac{\partial}{\partial y} \left(v \frac{\partial SST}{\partial y} \right) + \frac{\partial}{\partial y} \left(w \frac{\Delta T}{H} \right) - \frac{\partial}{\partial y} \left(\frac{Q_{net}}{\rho_0 c_p H} \right) - \frac{\partial R}{\partial y}, (2)$$

130 A bigger (smaller) GM indicates a stronger (weaker) NPSTF. A positive GM tendency 131 $(\partial GM/\partial t)$ suggests a process that GM gradually increases, corresponding to the 132 frontogenesis of the NPSTF. A negative GM tendency indicates the decreasing of GM, 133 corresponding to the frontolysis of the NPSTF.

134 2.3 Definition of the mixed-layer depth

Three definitions of mixed-layer depth H are used in this study: (a) 135 136 $SST - T_H = 0.5^{\circ}C$ (Qiu et al., 2014), where T_H is the temperature at the base of the mixed layer, and the depth of 0.5° C lower than the SST is defined as H. (b) 137 $SST - T_H = 1.0^{\circ}C$ (Suga and Hanawa, 1990), so the depth of 1.0°C lower than 138 the SST is defined as H. (c) mixed-layer depth from the GODAS. Figure 2a shows the 139 latitude-time section of the climatological mean mixed-layer depth calculated by 140 141 method (a) averaged from 140°E to 170°W (longitudinal region of the NPSTF in Fig. 1; Zhang et al., 2017). The mixed-layer depth exhibits significant seasonal variation, 142 namely, deep in winter and spring with a maximum of 60-80 m and shallow in 143 summer with a minimum of 20 m. Figure 2b shows the latitude-depth section of the 144 145 climatological mean zonal current velocities and ocean temperature gradients





averaged in winter and spring when the NPSTF exists. The NPSTF is mainly located 146 147 between 24°N and 30°N, with the maximum center expanding from the surface to the depth of 60 m. The vertical scale of the maximum center is consistent with the deeper 148 mixed layer in winter and spring calculated by method (a), suggesting the variation of 149 150 mixed-layer-averaged temperature gradient can well represent the variation of the NPSTF. The mixed-layer depth is also computed by methods (b) and (c). Expect for 151 152 the deeper depth in winter and spring (~80 m), their temporal evolutions of the 153 mixed-layer depth agree well with that in Fig. 2a, and the diagnosis results of Eqs. (1) 154 and (2) do not change qualitatively (not shown). Therefore, method (a) is used to define the mixed-layer depth in this study. In addition, two subsurface subtropical 155 temperature fronts are located between 80 and 180 m in Fig. 2b, associated with the 156 157 two branches of the North Pacific subtropical countercurrent, consistent with the 158 findings of Kobashi et al. (2006).

159

160 **3 Frontogenesis of the NPSTF**

Figure 3 shows latitude-time sections of the climatological mean GM and its tendency averaged over $(140^{\circ}\text{E}-170^{\circ}\text{W})$. The GM tendency is positive and moves southward from September to the following February. The NPSTF that forms in December is characterized by the SST gradient of 0.6 °C (100 km)⁻¹, which is the threshold for the emergence and disappearance of the NPSTF according to Qiu et al. (2014). Then, it strengthens and slightly migrates southward until March, with a maximum of 0.9 °C (100 km)⁻¹ at 27°N. Although the NPSTF is still robust in spring,





- it exhibits an evident northward shift with a strengthening in the northern part and a
 weakening in the southern and central parts. It finally disappears in July, consistent
 with the previous studies (Dinniman and Rienecker, 1999; Qiu et al., 2014). In this
 study, we mainly focus on the frontogenesis period of the NPSTF, which is from
 October to the following February when the GM tendency is significantly positive. As
 the NPSTF is located between 24°N and 30°N during this period, the frontogenesis
 region of the NPSTF is defined as (140°E–170°W, 24°N–30°N).
- 175 3.1 SST variation

176 Since the NPSTF is characterized by the meridional gradient of SST in the subtropics, the SST variation during the frontogenesis of the NPSTF is the first thing 177 we are interested in. Figure 4 portrays the temporal evolution of each term in Eq. (1) 178 179 over the NPSTF from October to the following February. As shown in Fig. 4a, the SST tendency is coherently negative during the frontogenesis, indicating that the SST 180 across the NPSTF gradually decreases. Note that the SST decreases more quickly in 181 the north than in the south, corresponding to the strengthening of the NPSTF. This 182 183 indicates that the largely decreasing SST in the north should be the key for the frontogenesis of the NPSTF. A diagnosis of each contributor on the right-hand side of 184 Eq. (1) is given in Figs. 4b-f. The SST tendency due to the net heat flux term (Fig. 4e) 185 bears similarities with the SST tendency in Fig. 4a in terms of spatial pattern and 186 187 magnitude, while the residual term (R) is mainly positive and facilitates an increasing SST. As for the oceanic intrinsic processes, the meridional temperature advection 188 serves as a much more important factor in determining the SST tendency compared to 189





190 the zonal and vertical temperature advections, especially in January and February. In 191 addition, the meridional temperature advection experiences a significant southward 192 displacement, which slightly increases the SST across the NPSTF in October and November and strongly decreases the SST in January and February. This is similar to 193 194 the southward migration of the GM tendency during the frontogenesis (Fig. 3). Overall, the SST across the NPSTF gradually decreases during the frontogenesis, 195 196 which is mainly attributed to the net heat flux term with some contributions from the 197 cold meridional advection in January and February. The residual term acts to suppress 198 this decreasing tendency.

199 3.2 GM variation

200 Figure 5a shows the temporal evolution of the climatological mean GM tendency 201 across the NPSTF from October to the following February. It is positive and moves 202 southward during the frontogenesis period, corresponding to the gradual enhancement of the NPSTF. Similar to the SST tendency from October to December (Fig. 4), the 203 GM tendency is mainly caused by the net heat flux term (Fig. 5e), while the residual 204 205 term acts to suppress the frontogenesis process (Fig. 5f). In January and February, the net heat flux term, together with the meridional temperature advection, favors the 206 frontogenesis of the southern and central NPSTF and suppresses the frontogenesis of 207 the northern NPSTF. The effect of R is nearly the opposite. Note that the magnitude of 208 209 the meridional temperature advection is quantitatively comparable to that of the net 210 heat flux term in January and February. Besides, the zonal and vertical temperature 211 advections (Figs. 5b and 5d) are negligible due to their smaller magnitudes. Figure 6a





further shows the regionally averaged GM tendency across the NPSTF during the 212 213 frontogenesis. The net heat flux term dominates the GM tendency from October to December and decreases after January. The meridional temperature advection 214 increases gradually from October to December, and plays an important role in January 215 216 and February. The residual term (R) mainly exerts an opposing influence on the frontogenesis except in January. These findings can be quantitatively illustrated in Fig. 217 218 6b. The net heat flux term controls the NPSTF frontogenesis from October to 219 December, while the meridional advection increases gradually and contributes equally 220 as the net heat flux in January and February. The results in January and February are consistent with those in Kazmin and Rienecker (1996), namely, the net heat flux and 221 the meridional Ekman convergence are equally important for the frontogenesis in 222 223 winter. In addition, the net heat flux also contributes to the disappearance of the 224 NPSTF in summer (not shown), which is consistent with the finding of Qiu et al. (2014). 225

Figure 7 shows the area mean GM tendency across the NPSTF calculated using 226 227 the Argo data from 2005 to 2013. Similar to Fig. 6a, the net heat flux term dominates from October to December and the meridional temperature advection works in 228 January and February. However, the effect of the meridional temperature advection is 229 overwhelmingly large in January and February, with much smaller net heat flux term 230 231 and R. This further confirms the dominant effect of the net heat flux term from 232 October to December and the important role of the meridional temperature advection in January and February for the frontogenesis of the NPSTF. Therefore, similar to the 233





previous studies (Kazmin and Rienecker, 1996; Dinniman and Rienecker, 1999), both
the net heat flux and oceanic meridional temperature advection contribute to the
frontogenesis of the NPSTF. As for the relative importance, the net heat flux
dominates the frontogenesis from October to December and then the meridional
temperature advection contributes equally as or even more than the net heat flux in
January and February.

240

241 **4 Roles of the Atmosphere**

242 4.1 Decomposition of the net heat flux

243 The net heat flux term is critical for the frontogenesis of the NPSTF from October

to December, which can be decomposed as follows:

245
$$\frac{Q_{net}}{\rho_0 c_p H} = \frac{Q_s}{\rho_0 c_p H} + \frac{Q_L}{\rho_0 c_p H} + \frac{Q_{LR}}{\rho_0 c_p H} + \frac{Q_{SR}}{\rho_0 c_p H},$$
(3)

where $Q_{\rm S}$, $Q_{\rm L}$, $Q_{\rm LR}$, and $Q_{\rm SR}$ represent sensible heat flux, latent heat flux, 246 longwave radiation, and shortwave radiation, respectively. Figure 8 shows the 247 248 temporal evolution of the GM tendency induced by individual heat flux terms in Eq. (3). The positive latent heat flux term primarily contributes to the positive GM 249 250 tendency in October, together with the sensible heat flux and the longwave radiation terms. The shortwave radiation term evidently strengthens in November and 251 December, and appears to be the dominant factor in January and February. Meanwhile, 252 the other three terms act to suppress the frontogenesis, especially the latent heat flux 253 254 term. Therefore, the four components of the net heat flux jointly contribute to the





- frontogenesis of the NPSTF, with a leading effect of the latent heat flux term in October and shortwave radiation term from November to February. Note that the temporal variation of the net heat flux term is consistent with that of the latent heat flux term. Moreover, the quick decrease of the net heat flux term in January is mainly attributed to the reduction of the latent heat flux term.
- 260 4.2 Cold meridional advection

261 As discussed above, the meridional temperature advection plays an important role 262 in the frontogenesis of the NPSTF in January and February (Fig. 6), which transports 263 the cold water from north to decrease the SST across the NPSTF. Figure 9 gives the 264 meridional Ekman convergence of $\partial (V_F \partial SST/\partial y)/\partial y$ calculated by the meridional Ekman velocity $V_E = -\tau_x / \rho_0 f H$, where τ_x is the zonal component of wind stress 265 and f is the Coriolis parameter. The meridional Ekman convergence moves southward 266 from October and strengthens in January and February, similar to the meridional 267 temperature advection (Fig. 5c). Moreover, in terms of contribution to the 268 frontogenesis, the Ekman convergence accounts for at least 75% of the meridional 269 270 temperature advection in January and February. Thus, the meridional temperature advection in January and February is mostly owing to the meridional Ekman 271 convergence. Note that $\tau_x = c_D \rho_a U^2$, where c_D is the drag coefficient, ρ_a is air 272 density and U is the surface zonal wind speed. Accordingly, the meridional Ekman 273 274 convergence must be associated with zonal wind speed. In the following, we focus on 275 possible atmospheric influence on the meridional temperature advection.

Figure 10 shows the climatological 1000-hPa geopotential height and wind fields





277	during the frontogenesis of the NPSTF. The weak Aleutian low primarily appears to
278	the east of the Bering Sea in October, with its center located over the Alaska Bay (Fig.
279	10a). The westerly wind to its south prevails zonally over 45° – 50° N. In November,
280	the Aleutian low develops and extends westward, longitudinally covering the whole
281	North Pacific (Fig. 10b). Accordingly, the prevailing westerly wind strengthens. The
282	Aleutian low keeps strengthening and heads southward from ~40°N in December to
283	~35°N in February (Figs. 10c-e). Meanwhile, the associated westerly wind is further
284	enhanced and shifted southward. Correspondingly, the westerly wind stress is
285	enhanced and moves southward from December to the following February. It can
286	force southward Ekman ocean currents in the Northern Hemisphere according to
287	$V_E = -\tau_x / \rho_0 f H$, leading to cold meridional advection. Moreover, the southward shift
288	of the westerly wind is consistent with the southward migration of the meridional
289	temperature advection in Fig. 4c. The above process can be seen clearly in Fig. 11.
290	Both the westerly wind and southward meridional ocean currents are obviously
291	increased, and move southward with the Aleutian low. Accordingly, the cold
292	meridional advection is enhanced and moves southward, cooling the SST across the
293	NPSTF in January and February. Li (2010) found that an Aleutian low-like anomalous
294	wind stress can decrease the SST in the mid-latitude North Pacific (north of 25° N) in
295	numerical models. Further analysis revealed that it is the cold meridional advection,
296	induced by the Aleutian low-like anomalous wind stress, acts to decrease the SST
297	north of 25°N. This previous study suggested that the strengthening and southward
298	migration of the Aleutian low can decrease the SST across the NPSTF via the cold





meridional advection. In addition, both the westerlies and the southward currents 299 300 reach the southern latitude of 28°N, resulting in colder SST in the northern NPSTF than in the southern NPSTF, corresponding to the frontogenesis of the NPSTF. The 301 cooler SST in the northern is also associated with the fact that the northern SST 302 303 cooling contributes greatly during the frontogenesis (Fig. 4a). Thus, the meridional Ekman convergence dominates the cold meridional advection, which may be related 304 305 to the strengthening and southward migration of the Aleutian low from October to the 306 following February. The associated westerly wind, together with the wind-driven 307 southward currents, is strengthened and shifts southward to induce cooler SST in the northern NPSTF, favoring its frontogenesis. 308

309 Note that the rapid decrease of the net heat flux term in January is mainly due to 310 the reduction of the latent heat flux term. The latent heat flux term can be calculated by $Q_L = \rho_a L C_E U_{10m} (q_s - q_a)$, where L is the latent heat of vaporization, C_E is the 311 312 bulk coefficient, U_{10m} represents the 10-m wind speed (Qiu et al., 2014). According to Eq. (2), the GM tendency is proportional to the meridional gradient of the 10-m wind 313 speed (- $\partial U_{10m}/\partial y$). Figure 12 shows the temporal evolutions of - $\partial U_{10m}/\partial y$ across 314 the NPSTF and GM tendency associated with the latent heat flux term. The 315 meridional gradient of wind speed gradually decreases from October to the following 316 February, consistent with the GM tendency calculated by the latent heat flux term, 317 especially from December to February. Interestingly, the decreasing - $\partial U_{10m}/\partial y$ is 318 319 also consistent with the southward migration of the Aleutian low (blue line in Fig. 12). This southward shift leads to a gradual increase in the wind speed to the south of the 320





Aleutian low (to the north of the NPSTF), corresponding to the decrease of $-\partial U_{10m}/\partial y$ between the NPSTF and its northern region, further resulting in the decease of the net heat flux term during the frontogenesis. Therefore, the Aleutian low acts to decrease the effect of the net heat flux and to increase the effect of the meridional temperature advection during the frontogenesis, which may also play an important role in transforming the dominant effect of the net heat flux to the joint effect of the meridional temperature advection and net heat flux in January.

328 4.3 Model results

329 Next, we use the NCAR Community Earth System Model version 1.0.3 (CESM1.0.3) to further investigate the role of the atmosphere. In the configurations 330 used here, the Community Atmosphere Model version 5.1 is coupled to the 331 332 Community Land Model version 4, the Los Alamos Sea Ice Model version 4 and a slab ocean model (Kiehl et al., 2006). The slab ocean model is essentially a 333 0-dimensional model and an approximation of the well-mixed ocean mixed layer. The 334 thermodynamic calculation uses a specified mixed-layer depth, and the temperature of 335 336 the slab is calculated based on the mixed-layer depth and surface fluxes. It means that the ocean dynamic processes can be ignored and the SST variation responds to the 337 atmosphere. In addition, the atmosphere and land models use a horizontal resolution 338 of $1.9^{\circ} \times 2.5^{\circ}$ (latitude \times longitude). A nominal resolution of 1° (gx1v6) is used for the 339 340 ice and ocean models. The model is run for 50 years, and the first 35 years are used 341 for spin-up to ensure that all model components reach their equilibrium states (Deng at al., 2017). The SST and meridional oceanic current velocity from the last 15 model 342





343 years are used for analyses.

344 Figure 13a shows the regionally-averaged GM and GM tendency over the NPSTF in the CESM1.0.3 simulation. The simulated GM tendency is positive from October 345 to the following February and turns negative in March, consistent with the 346 347 observations (Fig. 3), especially the southward shift during the frontogenesis. Accordingly, the NPSTF appears in December and disappears in July. Since the SST 348 349 in the slab ocean model is mostly due to the surface heat fluxes, it implies that the net 350 heat flux could cause the appearance of the NPSTF, further confirming the important 351 effect of the atmosphere on the frontogenesis of the NPSTF. Recall that in spring when the NPSTF is the strongest, the observed GM tendency is positive north of the 352 NPSTF and negative south of the front, corresponding to the northward shift of the 353 354 front (Fig. 3). However, this is absent in the CESM1.0.3 model, suggesting ocean dynamics may play an important role in the northward migration process, which 355 needs further exploration. Figure 13b shows the meridional temperature advection 356 term in the slab ocean model. It moves southward during the frontogenesis and 357 358 enhances gradually to a comparable value of the whole GM tendency in January and February, consistent with the observations (Fig. 5c). This corresponds to the 359 southward migration of the Aleutian low (Fig. 13c), confirming the atmospheric 360 influence on the meridional temperature advection. 361

362

363 **5. Summary**

364 We investigated the frontogenesis of the NPSTF occurring from October to the





following February based on the mixed-layer budget equation, with a focus on the 365 366 role of the atmosphere. In terms of the relative importance of the net heat flux and the Ekman convergence term, we find that the different terms dominated in different 367 periods of the frontogenesis. The net heat flux dominates the frontogenesis of the 368 369 NPSTF from October to December, while the meridional temperature advection 370 contributes equally as or even more than the net heat flux in January and February. 371 The zonal and vertical temperature advections can be neglected due to their smaller 372 magnitudes, while *R* acts to suppress the frontogenesis except in January.

373 The atmosphere is critical to the frontogenesis of the NPSTF, including the direct effect of the net heat flux and the indirect effect through the Aleutian low. A 374 decomposition of the net heat flux term reveals that its four components jointly 375 376 contribute to the frontogenesis, with a leading role by the latent heat flux in October and by shortwave radiation from November to the following February. Further 377 analyses of atmospheric effects on the oceanic process show that the meridional 378 Ekman convergence dominates the meridional temperature advection, and which is 379 380 associated with the Aleutian low variation. The strengthening and southward migration of the Aleutian low are characterized by the acceleration and southward 381 shift of the westerly wind to the south, which benefits southward ocean currents. 382 Accordingly, the cold meridional advection due to the southward currents induces 383 384 cooler SST in the northern NPSTF than in the southern NPSTF, and favors the 385 frontogenesis of the NPSTF in January and February. In addition, the reduction of the latent heat flux term (dominating the net heat flux term variation) during the 386





387	frontogenesis also results from the southward shift of the Aleutian low, suggesting
388	that the Aleutian low also plays a role in transforming the dominant effect of the net
389	heat flux to the joint contributions of meridional temperature advection and the net
390	heat flux in January. CESM1.0.3 model with the slab ocean model confirms the
391	important influence of atmosphere on the frontogenesis and on meridional
392	temperature advection.
393	
394	Data availability. The SODA, Argo and GODAS data can be downloaded fro
395	m: <u>http://apdrc.soest.hawaii.edu/data/data.php</u> . The OAFlux data are from: <u>ftp://ft</u>
396	p.whoi.edu/pub/science/oaflux/data v3/monthly/radiation 1983-2009/, and the ER
397	A-interim data are form: http://apps.ecmwf.int/datasets/data/interim-full-moda/levt
398	<u>ype=sfc/</u> .
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482 Captions

- 483 **Figure 1.** Climatological meridional SST gradients ($|\partial SST/\partial y|$, units: °C (100 km)⁻¹)
- 484 in (a) winter, (b) spring, (c) summer, and (d) autumn.

Figure 2. (a) Latitude-time section of the climatological monthly mean mixed-layer depth (units: m) calculated by $SST - T_H = 0.5 \,^{\circ}\text{C}$. (b) Latitude-depth section of the climatological zonal current velocity (black contour; units: m s⁻¹), superimposed with ocean temperature gradient (shading; units: $^{\circ}\text{C}$ (100 km)⁻¹); both are averaged over winter and spring. All three fields are averaged zonally over (140°E–170°W).

491 Figure 3. Latitude-time section of the climatological monthly mean gradient
492 magnitude (GM) of the NPSTF (black contour; units: °C (100 km)⁻¹) and its
493 tendency (shading; units: °C (100 km)⁻¹ month⁻¹), averaged zonally over (140°E–
494 170°W).

Figure 4. Latitude-time section of each term (shading; units: °C month⁻¹) in Eq. (1) 495 from October to the following February, averaged zonally over (140°E-170°W). 496 497 (a) shows the total SST tendency ($\partial SST/\partial t$), and (b-f) illustrate the components on the right-hand side of Eq. (1), namely, zonal temperature advection (Uadv), 498 meridional temperature advection (Vadv), vertical temperature advection (Wadv), 499 the net heat flux (Qnet), and the residual term (R). The black contours in each 500 panel are the same, indicating the climatological monthly mean GM (units: °C 501 $(100 \text{ km})^{-1}$), averaged zonally over $(140^{\circ}\text{E}-170^{\circ}\text{W})$. 502

503 **Figure 5.** Same as Fig. 4, except for the terms (units: °C (100 km)⁻¹ month⁻¹) in Eq.





504	(2).
505	Figure 6. (a) The area mean GM tendency (units: $^{\circ}C$ (100 km) ⁻¹ month ⁻¹) over the
506	NPSTF from October to the following February. (b) The contribution percentages
507	(units: %) of the right-hand side terms in Eq. (2) to the left-hand side term. The
508	black dashed line in (a) is the GM tendency of the NPSTF. Green, red, purple,
509	blue, and brown indicate zonal temperature advection (Uadv), meridional
510	temperature advection (Vadv), vertical temperature advection (Wadv), the net heat
511	flux (Qnet), and the residual term(R), respectively, in both (a) and (b). Note that
512	the ratios of Qnet to $\partial GM/\partial t$ in October and November in (b) are 165% and 112%,
513	respectively; we cap them at 100%.
514	Figure 7. Same as Fig. 6a, except using the Argo data from 2005 to 2013.
515	Figure 8. The area mean GM tendency (units: °C (100 km) ⁻¹ month ⁻¹) induced by the
516	net heat flux term (Qnet, blue), sensible heat flux term (Q_s , green), latent heat flux
517	term (Q_L , red), longwave radiation term (Q_{LR} , purple), and shortwave radiation
518	term (Q_{SR} , brown) over the NPSTF from October to the following February.
519	Figure 9. Same as Fig. 5c, except for the meridional temperature advection term
520	calculated by the Ekman velocity.
521	Figure 10. Climatological monthly mean geopotential height (shading; units: m ² s ⁻²)
522	and wind velocities (vector; units: $m s^{-1}$) at 1000 hPa in (a) October, (b) November,
523	(c) December, (d) January, and (e) February.
524	Figure 11. Latitude-time sections of (a) the climatological monthly mean geopotential
525	height (shading; units: m ² s ⁻²) and zonal wind speed at 1000 hPa (black contour;





- 526 units: $m s^{-1}$), (b) the climatological monthly mean meridional ocean currents (units:
- 527 m s⁻¹). All variables are averaged zonally over $(140^{\circ}\text{E}-170^{\circ}\text{W})$.

528	Figure 12. Meridional gradient of 10-m wind speed $(-\partial U_{10m}/\partial y, \text{ black, units: } 10^{-5})$
529	s ⁻¹) and GM tendency calculated by the latent heat flux (Q_L , red, units: °C (100
530	km) ⁻¹ month ⁻¹) over the NPSTF. The blue curve (AL) is the latitude of
531	climatological geopotential height at 900 m ² s ⁻² averaged zonally over (140°E–
532	170°W), representing the southward migration of the Aleutian low.

533	Figure 13. Latitude-time sections of (a) the total GM tendency ($\partial GM/\partial t$; shading;
534	units: °C month ⁻¹) and climatological monthly mean GM (black contour; units: °C
535	(100 km) ⁻¹), (b) meridional temperature advection (Vadv; black contour; shading;
536	units: $^{\circ}$ C month ⁻¹) and climatological monthly mean GM (units: $^{\circ}$ C (100 km) ⁻¹) in
537	Eq. (2), and (c) climatological monthly mean geopotential height (shading; units:
538	$m^2\ s^{\text{-}2})$ and zonal wind speed at 1000 hPa (black contour; units: $m\ s^{\text{-}1})$ from the
539	CESM1.0.3 simulation outputs, averaged zonally over (140°E–170°W).
540	







543 Figure 1. Climatological meridional SST gradients ($|\partial SST/\partial y|$, units: °C (100 km)⁻¹) in

- 544 (a) winter, (b) spring, (c) summer, and (d) autumn.
- 545







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Figure 3. Latitude-time section of the climatological monthly mean gradient
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- 601 December, (d) January, and (e) February.
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- 607 All variables are averaged zonally over (140°E–170°W).
- 608







610 Figure 12. Meridional gradient of 10-m wind speed $(-\partial U_{10m}/\partial y, \text{ black, units: } 10^{-5} \text{ s}^{-1})$

611 and GM tendency calculated by the latent heat flux $(Q_L, \text{ red, units: }^{\circ}C (100 \text{ km})^{-1}$

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Figure 13. Latitude-time sections of (a) the total GM tendency ($\partial GM/\partial t$; shading; units: °C month⁻¹) and climatological monthly mean GM (black contour; units: °C (100 km)⁻¹), (b) meridional temperature advection (Vadv; black contour; shading; units: °C month⁻¹) and climatological monthly mean GM (units: °C (100 km)⁻¹) in Eq. (2), and (c) climatological monthly mean geopotential height (shading; units: m² s⁻²) and zonal wind speed at 1000 hPa (black contour; units: m s⁻¹) from the CESM1.0.3 simulation outputs, averaged zonally over (140°E–170°W).