

# 1 On Studying Relations between Time Series in Climatology

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## 6 Abstract

7 Relationships between time series are often studied on the basis of crosscorrelation coefficients  
8 and regression equations. This approach is generally incorrect for time series irrespective of the  
9 crosscorrelation coefficient value because relations between time series are frequency-  
10 dependent. Multivariate time series should be analyzed in both time and frequency domains,  
11 including fitting a parametric (preferably, autoregressive) stochastic difference equation to the  
12 time series and then calculating functions of frequency such as spectra and coherent spectra,  
13 coherences, and frequency response functions. The example with a bivariate time series  
14 'Atlantic Multidecadal Oscillation (AMO) – sea surface temperature in Niño area 3.4 (SST3.4)'  
15 proves that even when the crosscorrelation is low, the time series' components can be closely  
16 related to each other. A full time and frequency domain description of this bivariate time series  
17 is given. The AMO – SST3.4 time series is shown to form a closed feedback loop system with  
18 a two-year memory. The coherence between AMO and SST3.4 is statistically significant at  
19 intermediate frequencies where the coherent spectra amount up to 55% of the total spectral  
20 densities. The gain factors are also described. Some recommendations are offered regarding  
21 time series analysis in climatology.

## 22 1 Introduction

23 Studying relations between time series on the basis of observations is a common task in all  
24 branches of Earth sciences. Normally, it requires getting quantitative answers to the following  
25 questions:

- 26 • what is the optimal time-domain stochastic model for a given multivariate time series?
- 27 • which components of the time series could be regarded as inputs and outputs of  
28 respective climatic system?
- 29 • is there any interaction between the inputs and the outputs (are there any closed-  
30 feedback loops within the system)?

- what are the statistical properties of the multivariate time series in the time and frequency domains?

In this article, we will apply the methods first developed in theory of information (Gelfand and Yaglom, 1956), time series analysis (Bendat and Piersol, 1966; Box and Jenkins, 1970), econometrics (Granger and Hatanaka, 1964; Granger, 1969), and in geophysics (Robinson, 1967; Robinson and Treitel, 1980) to study relations between multivariate time series of climatic data; the goal is to describe time series in the time and frequency domains, including climatic teleconnections that can hardly be found within the correlation/regression approach. Although methods of multivariate time series analysis are also described, mostly in the frequency domain, in climatology (von Storch and Ziers, 1999) and oceanography (Emery and Wilson, 2004), they are rarely applied in both time and frequency domains (e.g. Privalsky, 1988, 1995) or just in the frequency domain (e.g., Oladipo, 1987, Schneider and Schönwiese, 1989; Ghanbari et al., 2009; Park and Dusek, 2013).

## 2 Elements of multivariate time series analysis

Note first that the linear correlation/regression approach as a means of studying relations between scalar time series, including teleconnections within the climatic system, is generally inapplicable to time series analysis. The simplest example given in Privalsky and Jensen (1995) and repeated in Emery and Wilson (2004) is a zero crosscorrelation coefficient between two strictly linearly connected white noise sequences, one of which is obtained by applying a shift operator to the other. A low correlation coefficient may occur between any time series related to each other through more complicated but still strictly linear transformations. In particular, it can be a time series and its first difference, or any autoregressive-moving average (ARMA) time series and its innovation sequence, or the time series at the input and output of a linear filter. The general statement is that if a time series is obtained from another time series through a strictly linear inertial transformation the correlation coefficient between them will not be equal to 1 in spite of the strictly linear dependence between them.

Relations between two time series (say,  $A$  and  $B$ ), which are not mutually independent, always correspond to one of the following three situations:  $A$  affects  $B$  but not vice versa,  $B$  affects  $A$  but not vice versa,  $A$  and  $B$  affect each other (interaction). Again, the correlation/regression approach does not allow one to determine what the actual situation is. It can be done within the framework of the time-and-frequency domain analysis of multivariate time series.

1 The linear regression equation  $B = \alpha A + \varepsilon$ , where  $\alpha$  is a constant and  $\varepsilon$  is a white noise sequence  
 2 (regression error) means that the spectra  $s_A(f)$ ,  $s_B(f)$  of the time series  $A$  and  $B$  are identically  
 3 shaped because  $s_A(f) = \alpha^2 s_B(f) + \sigma_\varepsilon^2$  (where  $f$  is frequency and  $\sigma_\varepsilon^2$  the regression error variance).  
 4 This result is irrelevant when  $A$  and  $B$  are time-invariant random variables but if  $A$  and  $B$  are  
 5 time series, it puts an unnecessary limitation upon their properties. In the general case, the  
 6 shapes of the spectra are not identical, which would mean that  $\varepsilon$  is not white noise thus making  
 7 the regression equation inadequate. This is another reason why both the crosscorrelation  
 8 coefficient and respective linear regression equation cannot describe relations between time  
 9 series.

10 The problem is solved if one uses methods of time series analysis including simultaneous  
 11 description of multivariate time series in the time and frequency domains. It means fitting a  
 12 stochastic difference equation to the time series, analyzing its properties in the time domain and  
 13 then calculating and analyzing functions that describe the time series in the frequency domain.  
 14 For a number of considerations (see below), the approach used here will be limited to the  
 15 autoregressive (AR) case. Also, we will be regarding only the bivariate case. The extension to  
 16 higher dimensions is rather simple (*e.g.*, Bendat and Piersol, 1966; Robinson, 1967) and will  
 17 be briefly described at the end of this section.

18 Let the bivariate time series  $\mathbf{x}_n = [x_{1,n}, x_{2,n}]'$  be a (zero mean) sample record of an ergodic  
 19 discrete-time random process; here  $n = 1, \dots, N$  is the dimensionless argument,  $N\Delta t$  time series  
 20 length in time units  $\Delta t$ , and the strike means matrix transposition. The sampling interval  $\Delta t$  is  
 21 supposed here to be equal to 1.

22 In the time domain, the time series is described with a stochastic difference equation

$$23 \quad \mathbf{x}_n = \sum_{j=1}^p \Phi_j \mathbf{x}_{n-j} + \mathbf{a}_n, \quad (1)$$

24 where  $p < \infty$  is the order of autoregression,

$$25 \quad \Phi_j = \begin{bmatrix} \phi_{11}^{(j)} & \phi_{12}^{(j)} \\ \phi_{21}^{(j)} & \phi_{22}^{(j)} \end{bmatrix} \quad (2)$$

26 are matrix AR coefficients, and  $\mathbf{a}_n = [a_{1,n}, a_{2,n}]'$  is a bivariate innovation sequence (white noise)  
 27 with a covariance matrix

$$\mathbf{R}_a = \begin{bmatrix} R_{11} & R_{12} \\ R_{21} & R_{22} \end{bmatrix}. \quad (3)$$

The order  $p$  of an optimal AR model that agrees with the observed time series  $\mathbf{x}_n$  must be chosen on the basis of quantitative statistical considerations. Probably, the most efficient way to select an optimal order is to use order-selection criteria such as Akaike's AIC, Schwarz-Rissanen's BIC, Parzen's CAT, and Hannan-Quinn's  $\phi$  (e.g., Box et al., 2008; Parzen, 1977; Hannan and Quinn, 1979).

Properties of the time series  $\mathbf{x}_n$  in the frequency domain are defined with the spectral matrix

$$\mathbf{s}(f) = \begin{bmatrix} s_{11}(f) & s_{12}(f) \\ s_{21}(f) & s_{22}(f) \end{bmatrix}, \quad (4)$$

which is obtained through a Fourier transform of stochastic difference equation (1) fitted to the time series  $\mathbf{x}_n$ . Here  $f$  is frequency in cycles per sampling interval (in our case, year<sup>-1</sup>),  $s_{11}(f)$ ,  $s_{22}(f)$  are spectra and  $s_{12}(f)$ ,  $s_{21}(f)$  complex-conjugated cross-spectra of the time series components  $x_{1,n}$  and  $x_{2,n}$ . In particular, the coherence function

$$Co_{12}(f) = \frac{|s_{12}(f)|}{[s_{11}(f)s_{22}(f)]^{1/2}} \quad (5)$$

describes the linear dependence between the time series components  $x_{1,n}$  and  $x_{2,n}$  in the frequency domain. It can be thought of as a frequency-dependent set of correlation coefficients between components of a bivariate time series. It is the coherence function  $Co_{12}(f)$  (and not the crosscorrelation coefficient) that describes the degree of linear dependence between two scalar time series. Values of  $Co_{12}(f)$  satisfy the condition  $0 \leq Co_{12}(f) \leq 1$ .

The importance of the coherence function in time series analysis and modelling is illustrated with the following property. If the components of an ergodic bivariate time series present processes at the input and output of any linear time-invariant system, the coherence between them will be equal to 1 at all frequencies where the spectral density is not too close to zero.

The spectral matrix (4) describes a linear stochastic system with the time series  $x_{2,n}$  and  $x_{1,n}$  as the system's input and output, respectively. The coherent spectral density, or coherent spectrum  $Cs_{12}(f) = Co_{12}^2(f)s_{11}(f)$  defines the part of the output spectrum  $s_{11}(f)$  that can be explained by the linear dependence between  $x_{1,n}$  and  $x_{2,n}$ . The coherent spectrum  $Cs_{21}(f) = Co_{12}^2(f)s_{22}(f)$  describes the share of the spectrum  $s_{22}(f)$  defined by the contribution of  $x_{1,n}$  to  $x_{2,n}$ . Finally, the complex-valued frequency response function  $G_{12}(f) = s_{12}(f)/s_{22}(f)$  shows in what manner

1 the spectral energy  $s_{22}(f)$  of the input  $x_{2,t}$  is transformed into the spectral energy  $s_{11}(f)$  of the  
 2 output  $x_{1,t}$  as well as the phase difference between them (the gain factor  $|G_{12}(f)|$  and phase factor  
 3  $\phi_{12}(f)$ , respectively). The spectral characteristics calculated on the basis of Gaussian AR models  
 4 *with properly selected autoregressive orders* satisfy the requirements of the maximum entropy  
 5 method in spectral analysis. This is one of the reasons for selecting the AR modelling.

6 In the general case of an  $M$ -variate time series  $\mathbf{x}_n = [x_{1,n}, \dots, x_{M,n}]'$ , the time domain model is still  
 7 given by Eq. (1) with the matrix AR coefficients

$$8 \quad \Phi_j = \begin{bmatrix} \phi_{11}^{(j)} & \phi_{12}^{(j)} & \dots & \phi_{1M}^{(j)} \\ \phi_{21}^{(j)} & \phi_{22}^{(j)} & \dots & \phi_{2M}^{(j)} \\ \vdots & \vdots & \ddots & \vdots \\ \phi_{M1}^{(j)} & \phi_{M2}^{(j)} & \dots & \phi_{MM}^{(j)} \end{bmatrix} \quad (6)$$

9 and with the innovation sequence  $\mathbf{a}_n = [a_{1,n}, \dots, a_{M,n}]'$ . Its covariance matrix

$$10 \quad \mathbf{R}_a = \begin{bmatrix} R_{11} & R_{12} & \dots & R_{1M} \\ R_{21} & R_{22} & \dots & R_{2M} \\ \vdots & \vdots & \ddots & \vdots \\ R_{M1} & R_{M2} & \dots & R_{MM} \end{bmatrix}. \quad (7)$$

11 Consequently, the spectral matrix (4) changes to

$$12 \quad \mathbf{s}(f) = \begin{bmatrix} s_{11}(f) & s_{12}(f) & \dots & s_{1M}(f) \\ s_{21}(f) & s_{22}(f) & \dots & s_{2M}(f) \\ \vdots & \vdots & \ddots & \vdots \\ s_{M1}(f) & s_{M2}(f) & \dots & s_{MM}(f) \end{bmatrix}, \quad (8)$$

13 where  $s_{ii}(f)$  and  $s_{ij}(f)$  are spectral (if  $i = j$ ) and cross-spectral (if  $i \neq j$ ) densities, respectively, of  
 14 the time series components  $x_{i,n}$ ,  $i = 1, \dots, M$ . The elements of the spectral matrix (8) are used to  
 15 calculate spectra, multiple and partial coherence functions, multiple and partial coherent  
 16 spectra, and  $M-1$  frequency response functions (see Bendat and Piersol, 2010). The spectral  
 17 matrix (8) is Hermitian.

18 It should be noted that if the multivariate time series is long (by orders of magnitude longer  
 19 than the largest time scale of interest) and if the spectra of its components are intricate, the  
 20 above-described approach may not be the best, especially in the time domain – because of the  
 21 high order of the optimal stochastic difference equation. In such cases the analysis may have to

1 be limited to a frequency domain description of the time series by using the nonparametric (e.g.,  
 2 Percival and Walden, 1993, Bendat and Piersol, 2010) approach. However, this can hardly  
 3 happen if one is interested in properties of the time series at climatic time scales.

### 4 **3 An example of bivariate time series analysis**

5 The example with actual climatic data given below proves that the components of a bivariate  
 6 time series can be connected to each other in spite of the fact that the crosscorrelation coefficient  
 7 between them is low. It also provides a simultaneous description of a climatic system in both  
 8 time and frequency domains. The case with a high crosscorrelation coefficient between the  
 9 components of the ENSO system (Southern Oscillation Index and SST variations in the Niño  
 10 area 3.4) has been treated in detail in Privalsky and Muzylev, 2013 where it was shown, in  
 11 particular, that both time series are close to white noise, interact with each other mostly through  
 12 the innovation sequence, and that the coherence function, coherent spectra, and the frequency  
 13 response functions between SOI and SST are frequency dependent.

14 The El Niño–Southern Oscillation (ENSO) system is believed to affect many phenomena in the  
 15 Earth climate (*e.g.*, Philander, 1990). We will construct an autoregressive model of the bivariate  
 16 time series  $\mathbf{x}_n = [x_{1,n}, x_{2,n}]'$ , which consists, respectively, of the annual Atlantic Multidecadal  
 17 Oscillation (AMO) time series  $x_{1,n}$  (Kaplan SST V2 data provided by the NOAA/ OAR/ESRL  
 18 PSD, Boulder, Colorado, USA, see Enfield et al., 2001) and sea surface temperature SST3.4 –  
 19 an oceanic component of the ENSO system ( $x_{2,n}$ ). The SST3.4 data is computed from the  
 20 HadSST3 file at the University of East Anglia site (see Morice et al., 2011), for the same time  
 21 interval from 1876 through 2014 (Fig. 1a).

22 As seen from the figure, the two components behave in a different manner: AMO contains  
 23 much stronger low-frequency variations than SST3.4. The correlation between AMO and  
 24 SST3.4 (Fig. 1b) is very low, with the crosscorrelation coefficient 0.06 and the maximum  
 25 absolute values of the crosscorrelation function below 0.26. With the crosscorrelation-based  
 26 approach that prevails in climatology, the conclusion would have to be that the two scalar time  
 27 series are either not related to each other at all or that the connection between them is very  
 28 weak. And it would not be correct.

29 Consider first the time domain properties. All four above-mentioned order selection criteria  
 30 selected the same order  $p = 2$ . The respective AR(2) model is

$$\begin{aligned}
 31 \quad x_{1,n} &\approx 0.59x_{1,n-1} + 0.09x_{2,n-1} + 0.21x_{1,n-2} - 0.07x_{2,n-2} + a_{1,n} \\
 x_{2,n} &\approx -1.28x_{1,n-1} + 0.29x_{2,n-1} + 1.00x_{1,n-2} - 0.22x_{2,n-2} + a_{2,n}
 \end{aligned} \tag{9}$$

1 with the innovation covariance matrix

$$2 \quad \mathbf{R}_a = \begin{bmatrix} 0.014 & 0.007 \\ 0.007 & 0.290 \end{bmatrix}. \quad (10)$$

3 All coefficients in Eq. (9) are statistically significant at a confidence level of 0.95.

4 Obviously, Eq. (9) describes a closed feedback loop system: AMO ( $x_{1,n}$ ) depends upon two of  
5 its previous values and upon two previous values of SST3.4 ( $x_{2,n}$ ) and SST3.4 depends upon  
6 two previous values of both SST3.4 and AMO. The eigen-frequency of this system was found  
7 to be  $0.24 \text{ year}^{-1}$  so that respective period of about 4 years is the time required to close the  
8 system's feedback loop.

9 The stochastic difference equations (9) and the innovation sequence covariance matrix (10)  
10 allow one to understand how much of the variances of AMO and SST3.4 can be explained by  
11 the “deterministic” components of the model (9) that describe the dependence of  $x_{1,n}$  (AMO)  
12 and  $x_{2,n}$  (SST3.4) upon their own past values and upon past values of the other time series. The  
13 variance of AMO  $\sigma_1^2 \approx 0.035(\text{°C})^2$  while, according to Eq. (10), the variance  $R_{11}$  of the  
14 innovation sequence  $a_{1,n}$  is  $0.014(\text{°C})^2$ . Therefore, the part of the AMO variance  $\sigma_1^2$  which  
15 cannot be explained by the dependence of the time series upon their past behavior is  $0.014/0.035$   
16  $\approx 0.4$ . Thus, Eq. (9) allows one to explain about 60% of the AMO variance by its dependence  
17 upon its own past values and upon the past values of SST3.4.

18 The SST3.4 variance  $\sigma_2^2 \approx 0.37(\text{°C})^2$  while the variance  $R_{22}$  of the innovation sequence  $a_{2,n}$  is  
19  $0.29 (\text{°C})^2$ . Thus, the unexplainable part of the SST3.4 variance amounts to almost 80% of the  
20 SST3.4 variance. The results for AMO and SST3.4 differ because AMO and SST3.4 have very  
21 dissimilar spectra (Fig. 2a).

22 Both AMO and SST3.4 time series are Gaussian so that their autoregressive spectral estimates  
23 satisfy the requirements of the maximum entropy spectral analysis. The AMO spectrum  $s_{11}(f)$   
24 quickly decreases with frequency, which is characteristic of spatially-averaged climatic  
25 processes. Such behavior of the spectrum means that AMO is strongly dependent on its past  
26 values. The SST3.4 spectrum  $s_{22}(f)$  is not monotonic, has a wide maximum at intermediate  
27 frequencies and, in general, does not differ much from a white noise spectrum. The dependence  
28 on its past values is weak.

29 Though the crosscorrelation coefficient between AMO and SST3.4 is very low, the maximum  
30 entropy estimate of coherence  $C_{012}(f)$  shown in Fig. 2b is statistically significant everywhere  
31 except at the frequencies below  $0.13 \text{ year}^{-1}$  and above  $0.44 \text{ year}^{-1}$ . It is bell-shaped and exceeds

1 0.6 within the frequency band from  $0.20 \text{ year}^{-1}$  to  $0.36 \text{ year}^{-1}$ , that is, roughly, at time scales  
2 between 3 and 5 years. Its maximum value is 0.74 at  $f \approx 0.27 \text{ year}^{-1}$ . This behavior of the  
3 coherence function means that the components of this bivariate time series contribute to each  
4 other up to about 55% of the spectral density at intermediate frequencies.

5 The coherence between AMO and SST3.4 is weak at the low-frequency end, where AMO's  
6 spectral energy is much higher than elsewhere. The high coherence occurs at intermediate  
7 frequencies where the spectral density of AMO is much lower. The strong dependence of AMO  
8 on its past values and the relative closeness of the SST3.4 spectrum to a constant seem to be the  
9 reasons why the stochastic model (9) can explain so much of the total AMO variance and less  
10 of the total SST3.4 variance.

11 The contribution of SST3.4 to the AMO spectrum is  $Cs_{12}(f) = Co_{12}^2(f)s_{11}(f)$  where  $s_{11}(f)$  is  
12 the AMO spectrum (Fig. 3a). Respective contribution of AMO to SST3.4 shown in Fig. 3b is  
13  $Cs_{21}(f) = Co_{12}^2s_{22}(f)$ . These coherent spectrum estimates are statistically significant within the  
14 frequency band from  $0.18 \text{ year}^{-1}$  to  $0.38 \text{ year}^{-1}$  where they amount to 25% – 55% of the spectral  
15 densities  $s_{11}(f)$  and  $s_{22}(f)$ . This is a substantial contribution but it occurs within the frequency  
16 band where the spectral density of AMO is at least an order of magnitude smaller than at lower  
17 frequencies where the coherence between AMO and SST3.4 is not significant. The  
18 crosscorrelation coefficient that “integrates” the coherence function is small in spite of the  
19 relatively close connection between the two processes at moderate frequencies for at least two  
20 reasons: the complex structure of the interdependence between the time series components  
21 expressed by Eqs. (9) and (10) which cannot be described with a linear regression equation and  
22 the low absolute contribution of SST3.4 to the AMO variance.

23 This example also shows that using proper methods of analysis allows one to avoid filtering of  
24 time series in order to suppress ‘noise’. Indeed, though the low-frequency variations dominate  
25 the spectrum of AMO, the coherence function has revealed the ‘signal’ – a teleconnection  
26 between AMO on SST3.4 at intermediate frequencies where the AMO spectrum is low. This is  
27 another useful property of autoregressive time and frequency domain models.

28 Our frequency-domain results generally agree with the earlier results by Park and Dusek (2013)  
29 regarding the connection between AMO and the Multivariate ENSO Index (MEI) at  
30 intermediate frequencies. The authors used a nonparametric spectral estimation – singular  
31 spectrum analysis keeping 10 first empirical orthogonal functions that cover slightly over 75%  
32 of the time series total variances. At frequencies above  $0.15 \text{ year}^{-1}$ , our estimate of coherence



1 is quite similar to the coherence estimate in Park and Dusek (2013). However, the authors did  
2 not estimate the frequency response function (FRF) because, according to them, “a physically-  
3 based transfer function is likely [to be] nonlinear”. Actually, a nonlinear theoretical model of  
4 the FRF between AMO and SST3.4 time series would have been in disagreement with  
5 observations because both time series are Gaussian. At frequencies from  $0.1 \text{ year}^{-1}$  to  $0.4 \text{ year}^{-1}$ ,  
6 the gain factors of the empirical FRFs AMO – SST3.4 and SST3.4 – AMO (Fig. 4) behave  
7 in the same manner as the coherent spectra shown in Fig. 3 and their values at intermediate  
8 frequencies amount to approximately 0.1 and 4, respectively. The coherent spectra and gain  
9 factors are shaped similarly because the coherence function is bell-shaped and the AMO and  
10 SST3.4 spectra are rather flat at intermediate frequencies.

11 The phase factors in this case cannot give explicit information about the AMO – SST3.4 system  
12 because its feedback loop is closed (interaction between AMO and SST3.4). We cannot  
13 compare our spectra with those shown in Park and Dusek (2012) because their spectral  
14 estimates are given without confidence bounds but generally the shapes of the spectra at  
15 frequencies below  $0.5 \text{ year}^{-1}$  seem to be rather similar.

#### 16 **4 Conclusions and comments**

17 1. Relations between time series should not be studied on the basis of crosscorrelation  
18 coefficients and regression equations. An efficient approach within the framework of  
19 time series analysis includes two stages both involving parametric (preferably,  
20 autoregressive) modelling:

- 21 • fitting a stochastic difference equation to the time series (time domain), analyzing the  
22 selected model, and
- 23 • using the fitted equation to calculate and analyze frequency domain characteristics  
24 (spectra and coherent spectra, coherence function(s), gain and phase factor(s)).

25 This two-pronged approach is little known in climatology and related sciences.

26 2. Methods of multivariate time series analysis should be used in all cases, irrespective of  
27 the value of the crosscorrelation coefficient. The crosscorrelation and regression  
28 coefficients do not generally describe relations between time series. In particular, a low  
29 crosscorrelation coefficient does not necessarily mean the lack of even a strictly linear  
30 dependence between the time series.

- 1       3. The stochastic difference equation (9) with the innovation sequence covariance matrix  
2       (10) shows quantitatively that AMO and SST3.4 interact with each other so that AMO  
3       and SST3.4 can be regarded as either inputs or outputs to the AMO – SST3.4 system.  
4       It also reveals that the system's memory extends for two years. The dependence of AMO  
5       and SST3.4 upon their own past and upon the past of the other time series explains about  
6       60% and 25% of the AMO and SST3.4 variances, respectively.
- 7       4. The frequency domain analysis of the system shows that the spectra of AMO and  
8       SST3.4 behave in a different manner, with the AMO spectrum decreasing fast with  
9       frequency and with a relatively flat SST3.4 spectrum.
- 10      5. In spite of the very low crosscorrelation coefficient between the time series of AMO  
11      and SST3.4, a close linear dependence between them was shown to exist at intermediate  
12      frequencies corresponding to time scales from 3 to 5 years. The coherence between  
13      AMO and SST3.4 is statistically significant in a wide frequency band centered at 0.26  
14      year<sup>-1</sup> where the coherence peaks at 0.74. This result has been obtained earlier by Park  
15      and Dusek (2013) for a similar climatic system.
- 16      6. The coherent spectra AMO – SST3.4 and SST3.4 – AMO are statistically significant at  
17      frequencies from 0.18 to 0.38 year<sup>-1</sup> contributing 25%-55% of the spectral densities of  
18      SST3.4 and AMO to each other. The gain factors in the band behave in a manner similar  
19      to the behavior of the coherent spectra and peak at about 0.1 and 4.0, respectively.

20    These conclusions provide answers to the questions formulated at the introduction to this work.

21    Time series can often be treated as Gaussian. The ability to use a Gaussian approximation is  
22    important because for such time series the nonlinear approach cannot give better results than  
23    what is obtained within the linear approximation. This latter statement holds, in particular, for  
24    time series extrapolation within the framework of the Kolmogorov-Wiener theory. Also, as  
25    shown by Choi and Cover (1984), the random process that maximizes the entropy rate under  
26    constraints on  $p$  first values of the correlation function is a Gaussian autoregressive process of  
27    order  $p$ .

28    Many time series, especially at climatic time scales, are short, that is, their length does not  
29    exceed the time scales of interest by orders of magnitude. Therefore, the non-parametric  
30    analysis in the frequency domain may not be efficient because, with short time series, it would  
31    produce less reliable results. Besides, the nonparametric approach does not allow one to obtain

1 explicit stochastic models in the time domain. (These are two more reasons to prefer the  
2 parametric modeling.)

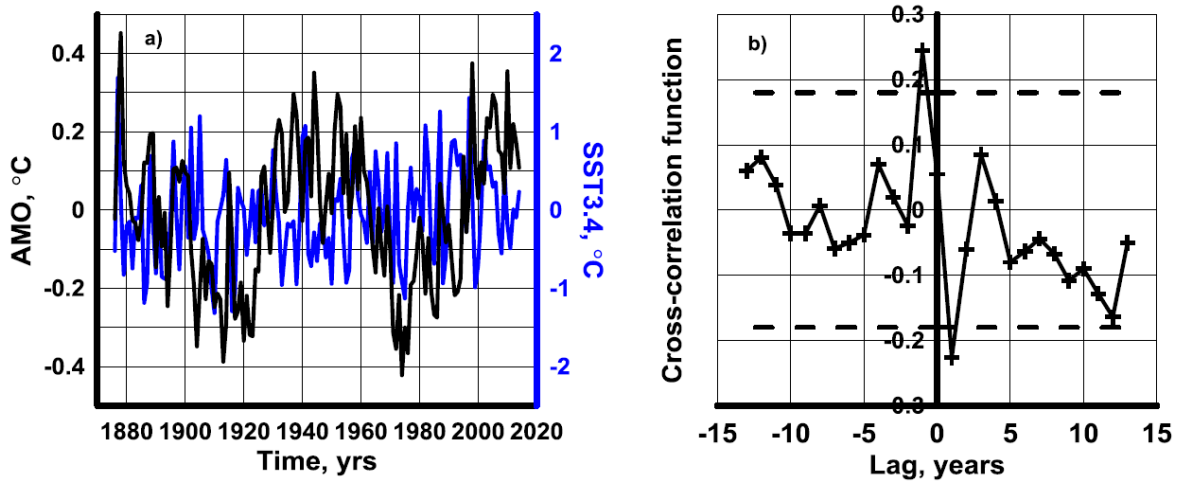
3 A parametric (first of all, autoregressive) analysis in time and frequency domains is effective  
4 because it results in relatively accurate estimates due to the postulation of a stochastic model  
5 for the time series. In particular, the frequency domain estimates obtained with the properly  
6 selected order satisfy, in the Gaussian case, the requirements of the maximum entropy spectral  
7 analysis. However, it is not correct to say that any autoregressive spectral estimate has this  
8 important property. The number of parameters to be estimated should always be much smaller  
9 than the time series length. If, for example, the length  $N\Delta t$  of a bivariate time series is 128 years,  
10 one should hardly expect statistically reliable results for models with AR orders higher than 5  
11 (20 AR coefficients plus 3 elements of the noise covariance matrix to be estimated). This is one  
12 of the reasons why it is not possible to make any realistic conclusions about the presence of 60-  
13 70 years long “periods”, “cycles”, “oscillations”, or about any other features at such large time  
14 scales unless the time series is at least 300 - 400 years long. For parametric models, it is strongly  
15 recommended to determine the model’s order and, consequently, the number of parameters to  
16 be estimated, on the basis of order selection criteria; an improper selection of the order  
17 invalidates the results of analysis.

## 18 **References**

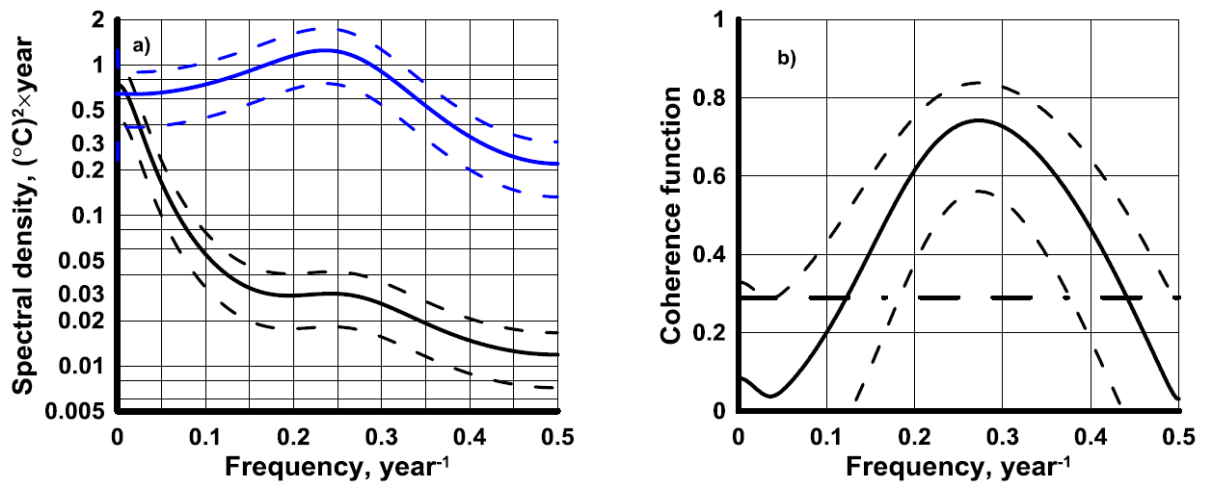
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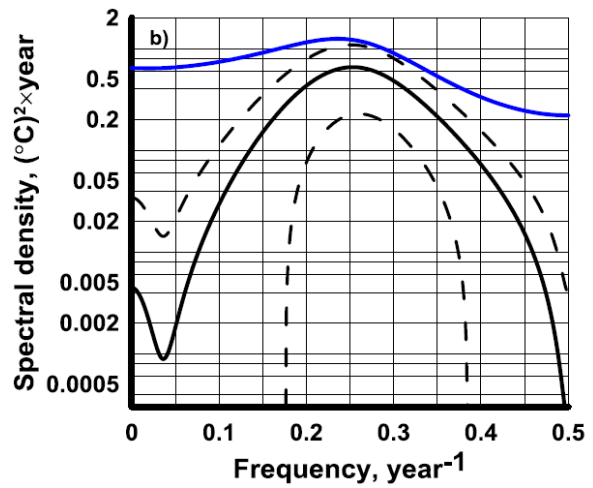
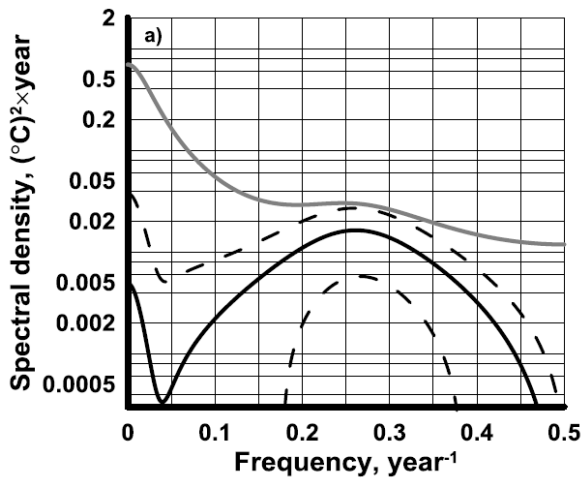
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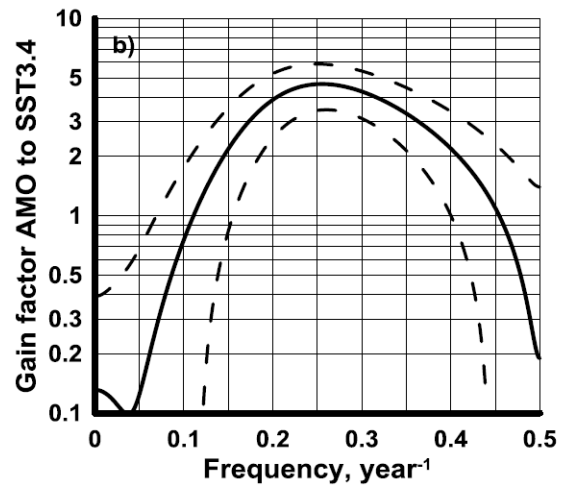
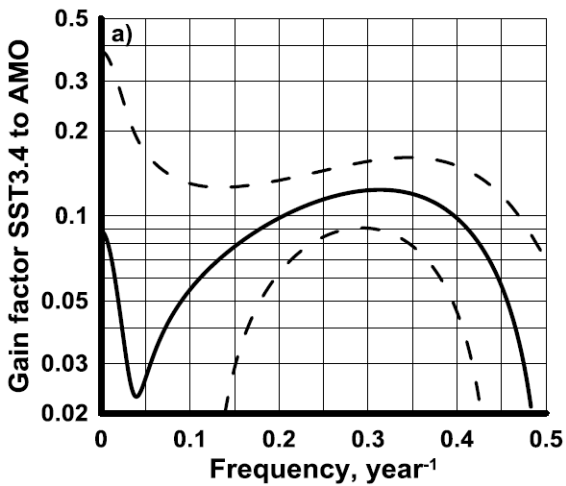
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 2 Figure 1. Observed AMO and SST3.4, 1876-2014 (a) and AMO–SST3.4 crosscorrelation  
 3 function estimate (b); the dashed lines show approximate 90% confidence bounds.



4  
 5  
 6 Figure 2. (a) spectra of AMO (black lines) and SST3.4 (blue lines) with approximate 90%  
 7 confidence bounds (dashed lines); (b) coherence function AMO – SST3.4 (solid line) with  
 8 approximate 90% confidence bounds (dashed lines). The dot-dashed line shows the  
 9 approximate upper 90% confidence bound for the true zero coherence. For confidence limits see  
 10 Privalsky et al. (1987).



1  
 2 Figure 3. Coherent spectra: contribution of SST3.4 to the AMO spectrum (a), and contribution  
 3 of AMO to the SST3.4 spectrum (b). Dashed, grey and blue lines show approximate 90%  
 4 confidence bounds and AMO and SST3.4 spectra, respectively.



5  
 6 Figure 4. Gain factors SSR3.4-AMO (a) and AMO-SST3.4 (b) with approximate 90%  
 7 confidence bounds (dashed lines).