

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

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F. Lehner^{1,2,*}, F. Joos^{1,2}, C. C. Raible^{1,2}, J. Mignot^{1,2,3}, A. Born^{1,2}, K. M. Keller^{1,2}, and T. F. Stocker^{1,2}

¹Climate and Environmental Physics, University of Bern, Switzerland

²Oeschger Centre for Climate Change Research, University of Bern, Switzerland

³LOCEAN Laboratory, Sorbonne Universités, France

* now at: National Center for Atmospheric Research, Boulder, USA

Received: 28 January 2015 – Accepted: 10 February 2015 – Published: 26 February 2015

Correspondence to: F. Lehner (lehner@climate.unibe.ch)

Published by Copernicus Publications on behalf of the European Geosciences Union.

This discussion paper is/has been under review for the journal Earth System Dynamics (ESD). Please refer to the corresponding final paper in ESD if available.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Abstract

Under the protocols of the Paleoclimate and Coupled Modelling Intercomparison Projects a number of simulations were produced that provide a range of potential climate evolutions from the last millennium to the end of the current century. Here, we present the first simulation with the Community Earth System Model (CESM), which includes an interactive carbon cycle, that continuously covers the last millennium, the historical period, and the twenty-first century. Besides state-of-the-art forcing reconstructions, we apply a modified reconstruction of total solar irradiance to shed light on the issue of forcing uncertainty in the context of the last millennium. Nevertheless, we find that structural uncertainties between different models can still dominate over forcing uncertainty for quantities such as hemispheric temperatures or the land and ocean carbon cycle response. Comparing with other model simulations we find forced decadal-scale variability to occur mainly after volcanic eruptions, while during other periods internal variability masks potentially forced signals and calls for larger ensembles in paleoclimate modeling studies. At the same time, we fail to attribute millennial temperature trends to orbital forcing, as has been suggested recently. The climate-carbon cycle sensitivity in CESM during the last millennium is estimated to be about $1.3 \text{ ppm } ^\circ\text{C}^{-1}$. However, the dependence of this sensitivity on the exact time period and scale illustrates the prevailing challenge of deriving robust constraints on this quantity from paleoclimate proxies. In particular, the response of the land carbon cycle to volcanic forcing shows fundamental differences between different models. In CESM the tropical land dictates the response to volcanoes with a distinct behavior for large and moderate eruptions. Under anthropogenic emissions, global land and ocean carbon uptake rates emerge from the envelope of interannual natural variability as simulated for the last millennium by about year 1947 and 1877, respectively.

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



1 Introduction

The last about 1000 years constitute the best opportunity prior to the instrumental period to study the transient interaction of external forcing and internal variability in climate, atmospheric CO₂, and the carbon cycle on interannual to multi-decadal time scales. In fact, the instrumental record is often too short to draw strong conclusions on multi-decadal variability. The relatively stable climate together with the abundance of high-resolution climate proxy and ice core data makes the last millennium an interesting target and testbed for modeling studies. Yet, the large and sometimes controversial body of literature on the magnitude and impact of solar and volcanic forcing on interannual to multi-decadal climate variability illustrates the challenges inherent in extracting a robust understanding from a period that is characterized by a small signal-to-noise ratio in many quantities and for which uncertainties in the external forcing remain (e.g., Wanner et al., 2008; PAGES 2k network, 2013; Schurer et al., 2013). In addition, a process-based quantitative explanation of the reconstructed preindustrial variability in atmospheric CO₂ and carbon fluxes is largely missing.

Compared to glacial–interglacial climate change, the last millennium experienced little climate variability, yet there is evidence for distinct climate states during that period (e.g., Lehner et al., 2012b; Keller et al., 2015). Within the last millennium the Medieval Climate Anomaly (MCA, ~ 950–1250 AD) and the Little Ice Age (LIA, ~ 1400–1700 AD) are two key periods of documented regional or global temperature excursions suggested to be driven by a combination of stronger solar irradiance and reduced volcanic activity and vice versa, respectively (e.g., Crowley, 2000; Mann et al., 2009; PAGES 2k network, 2013). Despite large efforts in reconstructing (e.g., PAGES 2k network, 2013) and simulating (e.g., Fernandez-Donado et al., 2013; Masson-Delmotte et al., 2013) the transition from the MCA to the LIA, substantial uncertainties remain with respect to the mechanisms at play. Recent studies point towards solar insolation playing a minor role for climate over the last millennium (Ammann et al., 2007; Schurer et al., 2014), while in turn regional feedback processes in response to volcanic eruptions and so-

ESDD

6, 351–406, 2015

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



lar variability need to be considered to explain decadal-scale climate variability (e.g., Lehner et al., 2013; Moffa-Sanchez et al., 2014). At high northern latitudes, the importance of millennial-scale orbital forcing is another debated issue (e.g., Kaufman et al., 2009).

5 Further, the last millennium offers the possibility to study the natural variability of the carbon cycle and its response to external forcing. Models with a carbon cycle module are extensively tested against present-day observations and widely used for emission-driven future projections (e.g., Hoffman et al., 2014). Yet, there are only few
10 last millennium simulations including a carbon cycle (e.g., Gerber et al., 2003; Jungclaus et al., 2010; Brovkin et al., 2010; Friedrich et al., 2012). The sensitivity of the carbon cycle to climate has been shown to be mostly positive, i.e., with warming additional CO₂ is released to the atmosphere (Ciais et al., 2013). However, the magnitude of this feedback remains poorly constrained by observations (Frank et al., 2010) and models (e.g. Friedlingstein et al., 2006). In particular, determining the role of the land
15 in past and future carbon cycle variability and trends is still challenging. In both idealized (Doney et al., 2006; Joos et al., 2013) and scenario-guided multi-model studies (Friedlingstein et al., 2014) the land constitutes the largest relative uncertainty in terms of intermediate- to long-term carbon uptake.

As for physical climate quantities, explosive volcanic eruptions constitute an important forcing for the carbon cycle. The sensitivity of the carbon cycle to such eruptions has been investigated by Jones and Cox (2001); Frölicher et al. (2011), or Brovkin et al. (2010) using different Earth System Models. For this short-lived forcing, the land response appears to be the driver of most post-eruption carbon cycle changes, with a range of magnitudes and time horizons associated with the different models. Further,
20 Frölicher et al. (2013) pointed out that the magnitude of the carbon cycle response to volcanoes depends critically on the climate state during the eruption.

25 The third Paleoclimate Modelling Intercomparison Project (PMIP3; Schmidt et al., 2011) and fifth Coupled Model Intercomparison Project (CMIP5; Taylor et al., 2012) represent joint efforts, in which different modeling groups perform identical experiments,

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



allowing for a systematic comparison of the models (e.g., Schmidt et al., 2014). Here we contribute to the existing set of simulations an integration from 850–2100 CE with the Community Earth System Model, including a carbon cycle module. The aims of this study are (i) to detect coherent large-scale features of forced variability in temperature and carbon cycle quantities, in particular in response to volcanic eruptions, (ii) to investigate the relative role of forcing uncertainty and model structural uncertainty, and (iii) to provide a preindustrial context to the future projections of climate change. The setup chosen here is unique in a number of ways and tailored to address the aims mentioned before: first, the carbon cycle is fully interactive with the other model components with the exception of the radiation code, which is fed by reconstructed CO₂. This allows us to study the isolated effect of climate on the carbon cycle, while guaranteeing an external forcing consistent with existing reconstructions. Second, the orbital parameters are held constant to study their importance relative to simulations with transient orbital parameters. Third, the solar forcing incorporated in the simulation has a larger amplitude than the majority of PMIP3 simulations and hence enables us to investigate whether the results are sensitive to this amplitude.

This paper is structured as follows: a description of the model and experimental setup is presented in Sect. 2. In Sects. 3 and 4 we address the general simulated climate and carbon cycle evolution and investigate forced and unforced variability of the simulated climate by comparing models to reconstructions and models to models. Sections 5 focuses on the response of the climate and carbon cycle to volcanic forcing. Section 6 deals with estimating the climate-carbon cycle sensitivity in CESM. A discussion and conclusions follow in Sect. 7.

2 Data and methods

2.1 Model description

The Community Earth System Model (CESM; Hurrell et al., 2013) is a fully-coupled state-of-the-art Earth System Model developed by the National Center for Atmospheric Research (NCAR) and was released in 2010. In terms of physics, CESM relies on the fourth version of the Community Climate System Model (CCSM4; Gent et al., 2011). Additionally, a carbon cycle module is included in CESM's atmosphere, land, and ocean components. The CESM version used here is release 1.0.1 in the so-called 1° version and includes components for the atmosphere, land, ocean, and sea ice, all coupled through a flux coupler.

The atmospheric component of CESM 1.0.1 is the Community Atmosphere Model version 4 (CAM4; Neale et al., 2010), which has a finite volume core with a uniform horizontal resolution of $1.25^\circ \times 0.9^\circ$ at 26 vertical levels. The land component is the Community Land Model version 4 (CLM4; Lawrence et al., 2011), which operates on the same horizontal grid as CAM4 and includes a prognostic carbon-nitrogen cycle that calculates vegetation, litter, soil carbon, vegetation phenology, and nitrogen states.

The ocean component is the Parallel Ocean Program version 2 (POP2; Smith et al., 2010; Danabasoglu et al., 2012) with an nominal 1° horizontal resolution and 60 depth levels. The horizontal resolution varies and is higher around Greenland, to where the North Pole is displaced, as well as around the equator. Embedded in POP2 is the Biogeochemical Elemental Cycle model (BEC; Moore et al., 2004) that builds on a nutrient-phytoplankton-zooplankton-detritus food web model and distinguishes three phytoplankton functional types (Long et al., 2013). Carbon export and remineralization are parameterized according to Armstrong et al. (2002). Alkalinity, pH, partial pressure of CO_2 , and concentrations of bicarbonate, and carbonate ions are diagnosed from prognostic dissolved inorganic carbon, alkalinity, and temperature- and salinity-dependent equilibrium coefficients. Organic material reaching the ocean floor is remineralized instantaneously, i.e., no sediment module is included. River discharge from

ESDD

6, 351–406, 2015

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



CLM4 does not carry dissolved tracers but nitrogen deposition to the ocean surface has been prescribed. The sea ice component is the Community Ice Code version 4 (CICE4) from the Los Alamos National Laboratories (Hunke and Lipscomb, 2010), including elastic-viscous-plastic dynamics, energy-conserving thermodynamics, and a subgrid-scale ice thickness distribution. It operates on the same horizontal resolution as POP2.

2.2 Experimental setup

Table 1 provides an overview of the simulations conducted for this study. First, a 500 year control simulation with perpetual 850 CE forcing (hereafter CTRL) was branched from a 1850 CE control simulation with CCSM4 (Gent et al., 2011). However, restart files for the land component were taken from a 850 CE control simulation, kindly provided by the NCAR, in which the land use maps by Pongratz et al. (2008) were applied. This procedure has the advantage that the slow-reacting soil and ecosystem carbon stocks are closer to 850 CE conditions than in the 1850 CE control simulation. A transient simulation covering the period 850–2099 CE was then branched from year 258 of CTRL. Despite the shortness of CTRL leading up to the start of the transient simulation, most quantities of the surface climate, such as air temperature, sea ice, or upper ocean temperature, can be considered reasonably equilibrated at the start of the transient simulation, as the forcing levels due to TSI and most greenhouse gases are similar between 1850 and 850 CE (Landrum et al., 2013). However, weak trends in CTRL are still detectable in slow-reacting quantities such as deep ocean temperature (below 2000 m; $\sim -0.04\text{ }^{\circ}\text{C } 100\text{ yr}^{-1}$) or soil carbon storage ($\sim 4\text{ Pg C } 100\text{ yr}^{-1}$).

The applied transient forcing largely follows the PMIP3 protocols (Schmidt et al., 2011) and the Coupled Model Intercomparison Project 5 (CMIP5; Taylor et al., 2012), consisting of total solar irradiance (TSI), greenhouse gases (GHGs), volcanic and anthropogenic aerosols, and land use changes (Fig. 1). Here, the TSI reconstruction by Vieira and Solanki (2010, TSI_{VS09}) is used, to which a synthetic 11 year solar cycle is added (Schmidt et al., 2011). In light of the recently enlarged envelope of reconstructed TSI amplitude (Schmidt et al., 2012), we scale TSI by a factor of 2.2635 to

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



have an amplitude of 0.2 % between present-day (1961–1990 CE) and the late Maunder Minimum (1675–1704 CE), which is about twice as large as the 0.1 % used in most PMIP3 simulations:

$$TSI = 2.2635 \cdot (TSI_{VS09} - \overline{TSI_{VS09}}) + \overline{TSI_{VS09}}. \quad (1)$$

Figure 1a shows that the TSI used here lies in between the large-amplitude reconstruction by Shapiro et al. (2011) and the bulk of small-amplitude reconstructions of the original PMIP3 protocol (Schmidt et al., 2011). Note, that a recent detection and attribution study indicates small amplitude TSI reconstructions to agree better with temperature reconstructions over the last millennium than large amplitude reconstructions (Schurer et al., 2014), in agreement with Ammann et al. (2007). For the twenty-first century the last three solar cycles of the data set are repeated continuously. The insolation due to Earth's orbital configuration is calculated according to Berger (1978) with the orbital parameters held constant at 1990 CE values.

The volcanic forcing follows Gao et al. (2008) from 850–2001 CE. It provides estimates of the stratospheric sulfate aerosol loadings from volcanic eruptions as a function of latitude, altitude, and month and is implemented in CESM as a fixed single-size distribution in the three layers in the lower stratosphere (Neale et al., 2010). Post-2001 CE volcanic forcing remains zero.

Land use and land use changes (LULUC) are based on Pongratz et al. (2008) from 850 to 1500 CE, when this dataset is splined into Hurtt et al. (2011), a synthesis dataset that extends into the future. The two datasets do not join smoothly but exhibit a small step-wise change in the distribution of crop land and pasture at the year 1500 CE. Up until about 1850 CE global anthropogenic LULUC are small, however, can be significant regionally (Hurtt et al., 2011). Towards the industrial era LULUC accelerate, dominated by the expansion of crop land and pasture. Here, only net changes in land use area are considered. The impact of shifting cultivation and wood harvest on carbon emissions from land use is neglected; these processes are estimated to have contributed in the

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



order of 30% to the total carbon emissions from land use (Shevliakova et al., 2009; Houghton, 2010; Stocker et al., 2014).

The temporal evolution of long-lived greenhouse gases (GHGs: CO₂, CH₄, and N₂O) is prescribed based on estimates from high-resolution Antarctic ice cores that are joined with measurements at mid-twentieth century (Schmidt et al., 2011, and references therein). While the carbon cycle module of CESM interactively calculates the CO₂ concentration originating from land use changes, fossil fuel emissions (post-1750 CE, following Andres et al., 2012), and carbon cycle-climate feedbacks, it is radiatively inactive. Instead, ice core and measured data are prescribed in the radiative code, keeping the physical model as close to reality as possible. As a result, the impact of the interactive coupling of the carbon cycle module is minor for simulated climate, and limited to changes in surface conditions due to changing vegetation. For the extension of the simulation post-2005 CE the Representative Concentration Pathway 8.5 (RCP 8.5) is used, representing the unmitigated “business-as-usual” emission scenario, corresponding to a forcing of approximately 8.5 W m⁻² at the year 2100 (Moss et al., 2010).

Aerosols such as sulfate, black and organic carbon, dust, and sea salt are implemented as non-time-varying up to 1850 CE, perpetually inducing the spatial distributions of the 1850 CE control simulation during this time (Landrum et al., 2013). Post-1850 CE, the time-varying aerosol datasets provided by Lamarque et al. (2010, 2011) are used, whereby CESM only includes a representation of direct aerosol effects. Similarly, nitrogen (NH_x and NO_y) input to the ocean is held constant until it starts to be time-varying from 1850 CE onwards, also following Lamarque et al. (2010, 2011). Iron fluxes from sediments are held fixed (Moore and Braucher, 2008).

2.3 Other model simulations

Besides to output from current Model Intercomparison Projects, we compare CESM results to those from similar simulation with CCSM4 (Landrum et al., 2013) and IPSL-CM5A-LR (Sicre et al., 2013), two simulations without interactive carbon cycle. Further,

ESDD

6, 351–406, 2015

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



we compare to MPI-ESM (ECHAM5/MPIOM; Jungclaus et al., 2010; Friedrich et al., 2012) to assess the robustness of the simulated climate and carbon cycle variations in response to external forcing. The solar and volcanic forcing of CCSM4 and IPSL-CM5A-LR are identical to ours with the exception of the scaling of TSI that we applied to CESM. MPI-ESM uses Krivova et al. (2007) as TSI forcing and Crowley et al. (2008) as volcanic forcing. These differ from the CESM forcing in amplitude much more than in timing and therefore allow for a comparison of the forced response. All simulations, except ours, apply transient orbital forcing. If not using the full ensemble of MPI-ESM, we focus on the member “mil0021”. Another difference in terms of experimental setup between CESM and MPI-ESM is that MPI-ESM was run with a fully interactive carbon cycle, i.e., the prognostic CO₂ interacts with the radiation and through that again influences climate, while in our setup this is a one-directional interaction only. Further, MPI-ESM is coarser resolved than CESM in both ocean and atmosphere and applies the A1B scenario for the twenty-first century (IPCC, 2000), which corresponds roughly to the current intermediate scenario as compared to the high scenario RCP8.5 used in CESM.

3 General climate and carbon cycle evolution

3.1 Temperature

The simulated annual mean Northern Hemisphere (NH) surface air temperature (SAT) follows the general evolution of proxy reconstructions: a warm Medieval Climate Anomaly (MCA, ~ 950–1250 CE), a transition into the colder Little Ice Age (LIA, ~ 1400–1700 CE), followed by the anthropogenically driven warming of the nineteenth and twentieth century (Fig. 2a). The NH MCA-to-LIA cooling amounts to $0.26 \pm 0.18^\circ\text{C}$ (taking the time periods defined above, which are as in Mann et al., 2009), placing it in the lower half of reconstructed amplitudes that range from about 0.1 to 0.7°C (IPCC, 2013). The subsequent warming from 1851–1880 to 1981–2010 CE amounts to

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



1.23 ± 0.15 °C, while observations report only 0.71 ± 0.13 °C (Cowtan and Way, 2014). This overestimation by CESM takes place almost entirely after 1960 and arises largely from missing negative forcing from the indirect aerosol effect, which is not implemented in CAM4 (Meehl et al., 2012). The late twentieth century being the warmest period in the NH in the past millennium is consistent with reconstructions (e.g., PAGES 2k network, 2013).

In CESM, the inception of the NH LIA occurs in concert with decreasing TSI and a sequence of strong volcanic eruptions during the thirteenth century. Reconstructions differ substantially in this matter and start to cool as early as 1100 CE or as late as 1400 CE. Further, new regional multi-proxy reconstructions of temperature provide no support for a hemispherical or globally synchronous MCA or LIA but show a clear tendency towards colder temperatures and exceptionally cold decades over most continents in the second half of the millennium (PAGES 2k network, 2013; Neukom et al., 2014).

The last millennium simulation with CCSM4 shows a largely coherent behavior with CESM in terms of amplitude and decadal variability of NH SAT (850–1850 CE correlation of 5 year filtered annual means $r = 0.88$, $p < 0.001$). The difference in NH SAT due to the different TSI amplitudes in CESM and CCSM4 scales roughly with the regression slope of NH SAT vs. TSI of both CESM and CCSM4 (~ 0.13 °C per $W m^{-2}$), although internal variability can easily mask this relationship at times. For example, the Maunder Minimum (1675–1704 CE), the 30 year period with the lowest TSI values and – by construction of the TSI scaling – with the largest difference between CESM and CCSM4 ($1.5 W m^{-2}$), is only 0.14 °C cooler than in CCSM4 and not 0.20 °C as expected from the regression.

The NH temperature evolution of additional PMIP3 and CMIP5 simulations shows that the multi-model range is within the one of the reconstructions and encompasses the instrumental-based observations (Fig. 2a). Disagreement between models and reconstructions exists in particular on the magnitude of response to the eruptions at 1258 and around 1350 CE. The 1258 CE eruption is the largest volcanic event recorded for

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CEF. Lehner et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

the last millennium and its climatic impact was likely enhanced through the cumulative effect of three smaller eruptions following shortly after (Gao et al., 2008; Crowley et al., 2008; Lehner et al., 2013). However, the pronounced cooling that is simulated by the models for this cluster of eruptions is largely absent in temperature reconstructions. Conversely, around 1350 CE temperature reconstructions show a decadal-scale cooling presumably due to volcanoes that is absent in the models, as the reconstructed volcanic forcing shows only two relatively small eruptions around that time. Part of this incoherent picture may arise from the unknown aerosol size distribution (Timmreck et al., 2010) and geographic location of past volcanic eruptions (Schneider et al., 2009), and differences in reconstruction methods. As many proxy reconstructions of temperature rely heavily on tree ring data it is worth noting that the dendrochronology community currently debates whether the trees' response to volcanic eruptions resembles the true magnitude of the eruption (Mann et al., 2012; Anchukaitis et al., 2012; Tingley et al., 2014).

Disagreement among the models exists on the relative amplitude of the MCA, where most models show colder conditions than CESM and CCSM4. Remarkably, the simulation by IPSL-CM5A-LR applied the same TSI and volcanic forcing as CCSM4, yet it comes to lie at the lower end of the PMIP3 model range during the MCA. In other words, the way how models respond to variations in TSI and other forcings can still make a larger difference in the simulated amplitude than the scaling of TSI by a factor of 2, which in turn complicates a proper detection and attribution of solar forcing during the last millennium (Servonnat et al., 2010; Schurer et al., 2014). Further disagreement among the models exists on the response to volcanic eruptions, where CESM and CCSM4 are among the more sensitive models (an oversensitivity of CCSM4 to volcanoes based on twentieth century simulations was reported by Meehl et al., 2012). Turning to the century-scale change over the industrial era, CESM and CCSM4 are on the upper end of the CMIP5 range and show an overestimation of the observed warming.

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

[Title Page](#)

[Abstract](#)

[Introduction](#)

[Conclusions](#)

[References](#)

[Tables](#)

[Figures](#)

[◀](#)

[▶](#)

[◀](#)

[▶](#)

[Back](#)

[Close](#)

[Full Screen / Esc](#)

[Printer-friendly Version](#)

[Interactive Discussion](#)



The simulated mean SAT of the Southern Hemisphere (SH) generally shows a similar evolution as for the NH with the signature of the MCA and LIA superimposed on a weak millennial cooling trend (Fig. 2b). Models and reconstructions disagree to a larger extent in the SH than in the NH, in particular regarding cold excursions due to large volcanic eruptions, which are largely absent in the reconstructions. Similar results have been reported in a recent study on interhemispheric temperature variations that finds much less phasing of the two hemispheres in reconstructions than in models, potentially related to underestimated internal variability on the SH in models (Neukom et al., 2014). However, the uncertainties in the early period of the reconstructions prohibits to robustly answer the question whether the models are too global in their response to external forcing. Similar to the NH, the industrial warming in the SH from 1851–1880 to 1981–2010 CE ($0.53 \pm 0.07^\circ\text{C}$) is overestimated by CESM ($0.71 \pm 0.13^\circ\text{C}$).

The differential warming between the hemispheres in CESM is among the smallest among CMIP5 models (not shown). This is mainly due to the underestimated deep water formation in the Southern Ocean, leading to a comparably strong warming of the SH and likely an underestimation of the oceanic uptake of anthropogenic carbon (Long et al., 2013). With a transient climate response of 1.73°C and an equilibrium climate sensitivity of 3.20°C (Meehl et al., 2012), CESM lies in the middle of recent estimates of 1.0 to 2.5°C and 1.5 to 4.5°C , respectively (IPCC, 2013).

3.2 Orbital forcing

To detect and attribute the influence of orbital forcing on SAT trends during the last millennium, we compare our simulation with fixed orbital parameters to the CCSM4 simulation with time-varying orbital parameters (Fig. 3). While both models experience a negative long term trend in global TSI until about 1850 CE (Fig. 1), the difference arising from the different orbital setup can be seen best in Arctic summer land insolation (Fig. 3). Hence, Arctic summer land SAT has been proposed as a quantity to be affected by orbital forcing already on time scales of centuries to millennia (Kaufman et al., 2009). However, we find no detectable difference between the two simulations in

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



the trend of Arctic summer land SAT (Fig. 3b). In fact, the Arctic multi-decadal to centennial summer land SAT anomalies in CESM span a very similar range as in CCSM4, despite CESM not accounting for time-varying orbital parameters: Fig. 3c shows non-overlapping 100 and 200 year mean SAT anomalies plotted against the corresponding mean solar insolation. The results from CCSM4 suggest a clear relation of the two quantities, however, the results of CESM show that nearly identical SAT anomalies are possible without orbital forcing. In other words, while we detect a long-term cooling trend in Arctic summer SAT in both CESM and CCSM4, we fail to attribute this trend to orbital forcing alone, as suggested by Kaufman et al. (2009). This is confirmed in new simulations with decomposed forcing, again comparing simulations with fixed and time-varying orbital parameters (B. Otto-Bliesner, personal communication, 2014).

3.3 Carbon cycle

The prognostic carbon cycle module in CESM allows us to study the response of the carbon cycle to transient external forcing. The land biosphere is a carbon sink during most of the first half of the last millennium, but becomes a source as anthropogenic land cover changes start to have a large-scale impact on the carbon cycle. The ocean is a carbon source at the beginning and becomes a sink in the second half of the last millennium (not shown). The residual of these fluxes represents changes in the atmospheric reservoir of carbon, illustrated in Fig. 2c by the prognostic CO₂ concentration. The amplitude of the simulated concentration does not resemble the one reconstructed from ice cores (i.e., imposed on the radiative code of CESM), in particular the prominent CO₂ drop in the seventeenth century is not captured by CESM. This raises the question whether the sensitivity of the carbon cycle to external forcing is too weak in CESM, whether the imposed land use changes are too modest (Kaplan et al., 2011; Pongratz et al., 2011), whether major changes in ocean circulation are not captured by models (Neukom et al., 2014), or whether the ice core records are affected by uncertainties due to in-situ production of CO₂ (Tschumi and Stauffer, 2000). Ensemble simulations with MPI-ESM also do not reproduce the reconstructed amplitudes or the

drop (Jungclauss et al., 2010). Further, Earth System Models of Intermediate Complexity or vegetation models driven by GCM output do not reproduce the uptake of carbon by either ocean or land needed to explain the reconstructed amplitudes (Stocker et al., 2011; Gerber et al., 2003).

The rise in atmospheric CO₂ due to fossil fuel combustion is in good agreement with ice cores until about the 1940s. After that, a growing offset exists, leading to an overestimation of about 20 ppm by 2005 in CESM, qualitatively similar to the CMIP5 multi-model mean (Hoffman et al., 2014). From the observational estimates one can diagnose that the discrepancy arises primarily from overestimated carbon release from land (Fig. 4a; see also Hoffman et al., 2014; Lindsay et al., 2014). From 1750 to 2011 CE the cumulative total land release (including LULUC) is 80 Pg C (compared to 30 ± 45 Pg C from observational estimates, Ciais et al., 2013), while the cumulative net land uptake is 101 Pg C (160 ± 90 Pg C, Ciais et al., 2013). The ocean cumulative uptake of 154 Pg C compares more favorably to current estimates of 155 ± 30 Pg C (Ciais et al., 2013). Note, however, that given the overestimation of atmospheric CO₂, one would expect a higher ocean uptake. This bias originates largely from an underestimation of the uptake in the Southern Ocean (Long et al., 2013). Along with this goes an underestimated seasonal cycle in CESM, originating from a too weak growing season net flux in CLM4 (Keppel-Aleks et al., 2013). MPI-ESM, on the other hand, underestimates atmospheric CO₂ due to weak emissions from LULUC (Pongratz et al., 2008).

The twenty-first century sees substantial emissions from fossil fuel burning under RCP 8.5 (Fig. 2c). In addition, LULUC is associated with a positive flux into the atmosphere, particularly until around 2050 CE (not shown). After accounting for LULUC (which constitutes a carbon loss for the land) the net land sink increases to about 7 Pg C yr⁻¹ at the end of the twenty-first century (Fig. 4a). The rate of ocean uptake, on the other hand, peaks around 2070 at about 5 Pg C yr⁻¹, despite that atmospheric CO₂ continues to rise (Fig. 2c). This decoupling of the trends in atmospheric CO₂ growth and ocean uptake flux is linked to non-linearities in the carbon chemistry. The change in dissolved inorganic carbon per unit change in the partial pressure of CO₂ decreases

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

[Title Page](#)

[Abstract](#)

[Introduction](#)

[Conclusions](#)

[References](#)

[Tables](#)

[Figures](#)



[Back](#)

[Close](#)

[Full Screen / Esc](#)

[Printer-friendly Version](#)

[Interactive Discussion](#)



with increasing CO₂, and thus the uptake capacity of the ocean. Additionally, differences in the ventilation time scales of the upper and the deep ocean likely play a role. While the surface ocean and the thermocline exchanges carbon on annual-to-multi-decadal time scales with the atmosphere, it takes century to ventilate the deep ocean as evidenced by chlorofluorocarbon and radiocarbon data (Key et al., 2004).

The prognostic atmospheric CO₂ increases to 1156 ppm by 2100 CE. This would imply a forcing of 7.6 W m⁻² from CO₂ relative to 850 CE, significantly more than the approximately 6.5 W m⁻² that are imposed by the radiative code (see Fig. 1c). This propagation of the twentieth century bias is consistent with the CMIP5 multi-model mean (Friedlingstein et al., 2014) and has motivated attempts to reduce such biases by using observational constraints for ocean ventilation (Matsumoto et al., 2004), the tropical land carbon storage sensitivity to temperature variations (Cox et al., 2013; Wenzel et al., 2014; Wang et al., 2014) and for the oceanic and terrestrial carbon fluxes (Steinacher et al., 2013). CESM with CLM4, however, shows very little sensitivity in tropical land carbon, in part due to the inclusion of an interactive nitrogen cycle, which – through enhanced photosynthetic uptake due to nitrogen fertilization – tends to counteract accelerated soil decomposition from warming (Lawrence et al., 2012; Wenzel et al., 2014). Together with the underestimated oceanic uptake this leads to the roughly 20 % larger airborne fraction in CESM as compared to the RCP 8.5.

Figure 4 puts the current and projected changes into perspective of preindustrial variability. Estimated interannual variability prior to 1750 CE is ± 0.94 Pg C yr⁻¹ (1 SD) for the net atmosphere–land and ± 0.42 Pg C yr⁻¹ for the net atmosphere–ocean flux. The much larger interannual variability in land than ocean flux is consistent with independent estimates and results from other models (e.g., Ciais et al., 2013). Large volcanic eruptions, as they have occurred in the last millennium, cause anomalously high uptake rates that for a short period of time are on par with current uptake rates (Fig. 4a and b, full range). We estimate when the anthropogenically forced, global-mean land and ocean uptake fluxes leave the bound of preindustrial natural variability (Hawkins and Sutton, 2012; Keller et al., 2014). As a threshold criteria, it is required that the decadal-

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



the preindustrial long-term negative trend are volcanic cooling events. In CESM many of these are global and are able to considerably cool the SH extra-tropics around 60° S, while in MPI-ESM the SH extra-tropics are only weakly affected. These differences are likely related to the Southern Ocean heat uptake rates in the two models (arising from under- and overestimation of Southern Ocean mixed layer depths in CESM and MPI-ESM, respectively; Danabasoglu et al., 2012; Marsland et al., 2003). This is evident also in the delayed warming at these latitudes in the twenty-first century in MPI-ESM as compared to CESM. The consistent SH high latitude positive anomalies before the thirteenth century, on the other hand, appear to be related to a positive phase of the Southern Annular Mode (SAM) in both models (not shown), a behavior common to most PMIP3 models. Note, however, that a recent reconstruction of the SAM finds the models to lack amplitude in their simulated variability, challenging the models' capabilities to represent SAM (Abram et al., 2014).

The phasing on interannual to decadal scales between the two models is largely restricted to periods of volcanic activity and within those mainly to land-dominated latitudes, except Antarctica, which shows no forced variability on these time scales (Fig. 5c). Despite the largest absolute temperature anomalies occurring in the Arctic, the correlations are highest in the subtropics, due to the smaller interannual variability there. Periods of centennial trends, such as the MCA or the Arctic cooling during the Maunder Minimum around 1700 CE, do not show up in the correlation analysis that focuses on 100 year windows, suggesting multi-decadal low-frequency forcing, such as centennial TSI trends, or internal feedback mechanisms to be responsible for the missing correlation. A regression analysis between the 5 year filtered annual TSI and SAT at each gridpoint (different filter lengths of up to 50 years have been tested as well without changing the results) reveals a clear link of the two quantities at high latitudes. In CESM this seems to be driven primarily by a displacement of the sea ice edge (Arctic) and Southern Ocean heat uptake (Fig. 6a). As the sea ice response has not been detected in an earlier model version (Ammann et al., 2007, their Fig. 4), it warrants the questions whether the regression of SAT on TSI might be biased by imprints of volca-

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



noes (Lehner et al., 2013), even when the timeseries are filtered, especially in a model like CESM that has a very strong volcanic imprint. Forthcoming simulations with solar-only forcing will be able to answer that question. MPI-ESM, on the other hand, shows a similar polar amplification signal from solar forcing, but not as clearly linked to sea ice (Fig. 6b). MPI-ESM also displays a stronger land-ocean contrast than CESM.

In addition to the comparison with MPI-ESM, Fig. 5d shows results from the correlation analysis between CESM and CCSM4, two simulations that in terms of physics differ only in their applied TSI amplitude and orbital parameters. Not unexpected, there are generally more robust signals of forced variability as compared to CESM vs. MPI-ESM (Fig. 5c), very likely due to the identical physical model components in CESM and CCSM4. Similarly, global mean SAT shows generally stronger phasing between CESM and CCSM4 (Fig. 5e). However, the latitudinal and temporal pattern of the CESM vs. CCSM4 analysis agrees well with the one arising from CESM vs. MPI-ESM (Fig. 5c; with exception of the much stronger phasing in CESM and CCSM4 during the volcanic eruptions in the 1450s) and suggest the physical mechanism behind periods of phasing to be robust across the two models.

Applied to ocean temperature, the above approach enables us to investigate the penetration depth of a forced signal seen at the surface (Fig. 7). Indeed, most of the surface signals also show up as significant correlations down to depths of about 150–200 m, whereby their timing suggests again volcanic forcing as the origin. The Atlantic Meridional Overturning Circulations (AMOC) in the CESM and MPI-ESM shows no significant correlation, however, the highest correlation occurs during the thirteenth century and coincides with a phasing of the upper ocean temperatures due to strong volcanic forcing (Fig. 7d). The correlation between CESM and CCSM4 at that time is even higher and points to a significant imprint of the volcanic forcing on ocean circulation (Otterå et al., 2010; Swingedouw et al., 2013). However, during the remainder of the millennium, no phasing of the AMOC is found.

4.2 Carbon cycle

We apply the same correlation analysis to zonally integrated land and ocean carbon fluxes from the two models to detect forced variability in the carbon cycle. Compared to SAT hardly any phasing can be found between the models in atmosphere-to-land carbon fluxes (not shown), which is due to its large interannual variability and to distinctly different responses to external forcing in the two models, as will be illustrated in Sect. 5. Similarly, there is little model phasing in net atmosphere-to-ocean carbon fluxes (not shown). Results become somewhat clearer when considering globally integrated upper-ocean dissolved inorganic carbon (DIC; Fig. 8). While there appear to exist spurious trends at depth in both models, there are periods of coherent carbon draw-down coinciding with volcanic eruptions around 1450 and 1815 CE in response to temperature-driven solubility changes. Interestingly, MPI-ESM shows a distinct behavior for the strong eruption of 1258 CE, with a prolonged ocean carbon loss after a weak initial uptake. CESM shows a stronger and more sustained carbon uptake, leading to no correlation between the two models for this eruption. The reasons for this discrepancy are discussed in Sect. 5.

Generally, the largest changes in upper-ocean carbon storage occur in response to volcanoes and take place in the tropical Pacific, with other significant changes occurring in the North and South Pacific, the subtropical Atlantic and the Arctic (Sect. 5). Within the tropical oceans, the models show different characteristics: CESM shows a larger variability in DIC than MPI-ESM and, when influenced by anthropogenic emissions in the twentieth and twenty-first century, takes up a larger portion of the total ocean carbon uptake than in MPI-ESM (not shown). In MPI-ESM, the Southern Ocean shows stronger variability and larger carbon uptake in the twenty-first century, illustrating the different behavior of the two models in terms of ocean carbon cycle variability and trend magnitude, closely related to the different mixed layer depth in the Southern Ocean region.

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



5 Volcanic forcing

To further isolate the response of the climate system and carbon cycle to volcanic eruptions, a Superposed Epoch Analysis is applied to both simulations. Thereby, composite time series for the strongest three (top3) and following strongest seven eruptions (top10), by measure of optical depth anomaly, over the period 850–1850 CE are calculated for the CESM and MPI-ESM (Fig. 9). The time series are calculated as de-seasonalized monthly anomalies to the 5 years preceding an eruption.

The physical parameters global mean surface air temperature and global mean precipitation decrease in both models after volcanic eruptions, although the response of CESM is stronger by roughly a factor 2–2.5 (Fig. 9a, b, f, g). Consequently, CESM temperature and precipitation take longer (~ 15 years) to relax back to pre-eruption values than MPI-ESM (~ 9 years).

The atmospheric carbon inventory, on the other hand, shows a remarkably different response in the two models. In CESM the atmosphere initially loses about 2–3 Pg C, ir-respectively of the eruption strength, with the minimum occurring after about 1–2 years. In the top10 case values return to normal after about 16 years, while in the top3 case they tend to return already after about six years, and overshoot. This overshoot is not straightforward to understand and did not seem to occur in earlier versions of the model (Frölicher et al., 2011). In MPI-ESM the response is a priori more straightforward and slower: in the top10 case the atmosphere loses about 2.5 Pg C, reaches a minimum after 2–4 years, and returns to pre-eruption values after 10–16 years. The top3 case reaches its minimum (–6 Pg C) a bit faster, but then takes about 20 years to return to pre-eruption values (Brovkin et al., 2010).

Partitioning these atmospheric carbon changes into land and ocean changes indicates that the land is primarily responsible for the differing response behavior of the two models, confirming the findings in the previous section. While in both models the land drives the atmospheric change by taking up carbon initially, it is released back to the atmosphere within about 3 years in CESM, but kept in the land for at least 15 years

ESDD

6, 351–406, 2015

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



in MPI-ESM (and up to 50 years for the 1258 CE eruption; Brovkin et al., 2010). In the top3 case of CESM the land starts to even loose carbon after about 5 years, causing the overshoot seen in the atmospheric carbon.

A closer look at CESM reveals a distinct response to the top3 and the top10 volcanoes. The response to top3 must be understood as an interplay of a number of processes: the initial global cooling triggers a La Niña-like response and a corresponding cloud and precipitation reduction that is particularly pronounced over tropical land, where also large changes in carbon storage occur (see Fig. 11a–c for the spatial pattern). Figure 10 and the following analysis therefore focuses on tropical land. Direct solar radiation decreases, indirect radiation increases, with a net decrease (Fig. 10d). These unfavorable conditions cause a reduction in net primary productivity and a strong decrease of vegetation (–8 Pg C; Fig. 10a and e). At the same time, decomposition of dead biomass becomes less efficient due to reduced temperature (similar to, e.g., Frölicher et al., 2011). Despite the simultaneous decrease in net primary production this results in a build-up of dead biomass of about 5 Pg C (Fig. 10b). Although carbon loss due to fire increases, it cannot get rid of the large amount of dead biomass immediately (Fig. 10f). While vegetation decrease and dead biomass buildup balance each other, the soil takes up about 2 Pg C (Fig. 10c), stores it for at least 16 years, and is therefore responsible for the initial net land uptake seen in Fig. 9e (see also Fig. 11c left). After about two years, tropical precipitation increases again and puts a halt to the decrease in vegetation (Figs. 10a and 11b right). The vegetation does not recover fully for another about 20 years. The dead biomass, on the other hand, gets decomposed entirely within about 15 years and therefore turns the land into a carbon source, causing the overshoot in CO₂. In the top10 case, the precipitation and radiation response is about half of the top3 case, and so is the vegetation decrease. Consequently, vegetation recovers faster. The decomposition of dead biomass, however, takes about the same amount of time as in the top3 case as the decomposition rates are similar for both cases. Hence, the land acts as a more sustainable carbon sink in the top10 case. In MPI-ESM it is the soil as well which acts as main land carbon storage pool, while the

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



vegetation decrease is significantly less than in CESM (Brovkin et al., 2010), leading to the different response behavior of the two land models, particularly striking in the top3 case. Note that there are subtle regional differences between CESM and the earlier version of the carbon cycle-enabled NCAR model CSM1.4-carbon (Frölicher et al., 2011): tropical Africa sees a reduction of land carbon in CESM, related to a persistent increase in cloud cover and precipitation after volcanoes, while CSM1.4-carbon saw a decrease in precipitation and an increase in land carbon.

The ocean, on the other hand, shows a qualitatively similar response in CESM and MPI-ESM with an uptake of carbon and a gradual relaxation back to pre-eruption values over 20 or more years. In CESM the radiative cooling leads to increased uptake in the Western Pacific, while in the Eastern Pacific, cooling is less as this region is more controlled by upwelling rather than direct radiative forcing, as suggested by Maher et al. (2014) (Fig. 11d). Two or more years after the volcano a La Niña-like pattern settles in both surface temperature as well as carbon uptake. Some model differences exist, e.g., in the top3 case of MPI-ESM the ocean starts to release carbon, compensating the persistent positive anomaly in the land inventory (imposed on the ocean via atmospheric CO₂ concentration Brovkin et al., 2010), a feature not present in CESM, in which the land does not store the anomalous carbon as long. In CESM the tropical oceans appear to be more sensitive to volcanic forcing than to TSI variations. The equatorial Pacific shows the strongest response in DIC to volcanoes (Fig. 11d), while the response to TSI variations of comparable radiative forcing is up to an order of magnitude weaker and confined to higher latitudes (not shown). Overall it seems therefore that the response of the land vegetation governs the overall different responses in the two models.

In an attempt to validate the two models, one is restrained to the well-observed eruption of Pinatubo in 1991 CE, as the CO₂ records from ice cores do not adequately resolve short-term variations induced by volcanoes over the last millennium. Figure 12 shows the global temperature and atmospheric carbon response to Pinatubo as extracted from observations, CESM, and the 3-member ensemble of MPI-ESM. Note that

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



the effects of El Niño-Southern Oscillation and anthropogenic emissions have been removed from the CO₂ observations to obtain a tentative estimate of the actual CO₂ response to the Pinatubo eruption (Frölicher et al., 2013). The initial cooling of about –0.5 °C and the relaxation back to initial temperatures around 1998 CE is captured well by both models. The MPI-ESM ensemble, however, shows a large and robust variation around 1995 CE, seemingly related to a phasing of ENSO variability in response to the eruption (see also Zanchettin et al., 2012). Further, the magnitude of atmospheric carbon response matches better in CESM, although the overshoot of the observation-based estimate is not captured. CESM’s response also falls within the range of the earlier model version (Frölicher et al., 2013). It remains unclear whether this mismatch reflects a model-deficiency or is due to uncertainties arising from removing the ENSO signal from the CO₂ observations. However, the mechanisms described above that lead to an atmospheric CO₂ overshoot for large eruptions in CESM offer an opportunity for reconciliation of this discrepancy. Further, the precipitation response (and therewith the cloud and surface short-wave response) to volcanic eruptions is not well constrained due to the small number of observed eruptions (Trenberth and Dai, 2007). Biases in the representation of these processes can influence a model’s carbon cycle response.

6 Climate-carbon cycle sensitivity

Due to the absence of large anthropogenic disturbances of the carbon cycle, the last millennium represents a testbed to estimate the climate-carbon cycle feedback sensitivity γ (ppm °C⁻¹) and can thus potentially help to constrain this quantity (e.g., Woodwell et al., 1998; Joos and Prentice, 2004; Scheffer et al., 2006; Cox and Jones, 2008; Frank et al., 2010). Here, we use the experimental setup of CESM to estimate γ , mimicking to some extent the methods by Frank et al. (2010) and Jungclaus et al. (2010).

We focus on the period before significant LULUC (850–1500 CE) and apply different low-pass filters of 20 to 120 years, taking 5 year increments, to the time series of NH

Applying the identical analysis to CTRL reveals other time scales of climate-carbon cycle feedback, suggesting maximum lags of less than 10 years and a γ of 2.3 (1.4–2.9) ppm °C⁻¹. A later peak in the lag correlation of CTRL clusters at 73.3 ± 1.1 years, i.e., close to where the forced simulation shows its highest lag correlation, but these lag correlations are much weaker ($r \sim 0.4$ compared to $r \sim 0.7$ in the forced simulation). This is generally consistent with the finding by Jungclaus et al. (2010) that a forced simulation exhibits increased power on lower frequencies compared to a control simulation.

7 Discussion and conclusions

This study presents a simulation from 850 to 2100 CE with the fully-coupled CESM, including carbon cycle, and investigates the imprint of external forcing on different climate and carbon cycle diagnostics. For comparison we draw on a number of PMIP3 simulations, particularly, comparable simulations with CCSM4 and MPI-ESM. The evolution of NH SAT during the preindustrial era in CESM is in reasonable agreement with both reconstructions and other models, albeit the uncertainties in reconstructions and forcing are still considerable. Comparing to more reliable data in the twentieth century, the anthropogenic warming in CESM is overestimated due to a lack of negative forcing from indirect aerosol effects. On the SH, CESM and most other models do not capture the evolution of the mean SAT as well. The discrepancies could be explained by (i) significant model biases in SH and also interhemispheric SAT variability (Neukom et al., 2014), (ii) spectral biases in proxies used in the reconstructions (Franke et al., 2013), (iii) uncertainties in the external forcing (Masson-Delmotte et al., 2013), or (iv) natural internal variability (Bothe et al., 2013). Unfortunately, these potential explanations are neither exclusive nor independent. Arguments for model bias come from the fact that reconstructed interhemispheric SAT variability lies outside the models' range over 40 % of the time (Neukom et al., 2014); but these arguments are weakened by the uncertainty in external forcing. We show here that implementing the same TSI forcing in two

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Swingedouw et al., 2013). Under anthropogenic emissions, land and ocean carbon uptake rates emerge from the envelope of natural variability as simulated for the last millennium by about 1947 and 1877 CE, respectively. Atmospheric CO₂ and global temperature emerge by 1755 and 1966 CE, suggesting that changes in carbon-cycle related variables would be easier to detect than temperature given sufficient observational data (Keller et al., 2015).

We find forced decadal-scale variability in CESM and MPI-ESM in response to major volcanic eruptions in both SAT and upper-ocean temperature, while the response in carbon cycle quantities is less coherent among models. Outside volcanically active periods large parts of the decadal-scale variations cannot be attributed to external forcing, suggesting that internal variability masks external forcing influence. Note, however, that recent work suggest that small volcanic eruptions, which are typically not well-resolved in reconstructions of volcanic activity, exert a significant cumulative effect on global temperature and climate (Ridley et al., 2014).

Volcanoes trigger a coherent global response in SAT and precipitation that is qualitatively in line with earlier studies on the volcanic influence on climate and carbon cycle (e.g., Jones and Cox, 2001; Brovkin et al., 2010; Frölicher et al., 2011, 2013). However, the carbon cycle response, in particular on land, shows fundamental model differences in terms of perturbation amplitude and persistence after volcanic eruptions. These differences arise from a differing land vegetation responses in the two models. The extent to which such structural uncertainties matter is illustrated by the large spread in the airborne fraction of CO₂ between these two (and other) models in the twenty-first century (see also Friedlingstein et al., 2014). In particular, known biases in CESM's carbon uptake in response to anthropogenic emissions in the twentieth and twenty-first century lead to a 20% overestimation of the atmospheric CO₂ concentration and the corresponding prognostic radiative forcing as compared to the prescribed RCP8.5 at year 2100 CE.

The climate-carbon cycle sensitivity of CESM as estimated from the anthropogenically unperturbed first part of the last millennium is about 1.3 ppm °C⁻¹, with a depen-

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



5 dency on the filtering and the exact time period considered. Generally, the sensitivity of the carbon cycle to temperature variations in CESM is comparably small (Frank et al., 2010) and reveals a strong component of unforced natural variability. In a transient last millennium simulation with small temperature variations, the proper detection of a lead–lag relation between temperature and the carbon cycle is complicated by the superposition of perturbations and responses. In addition to the classic climate-carbon cycle sensitivity experiments (e.g., Arora et al., 2013) it is therefore desirable to conduct step function-like sensitivity experiments in order to isolate the response of the carbon cycle to a particular external forcing (Gerber et al., 2003).

10 Despite the challenges that paleoclimate modelling faces, a number of lessons regarding forcing and structural uncertainties can be learned from these experiments. In order to better understand the role of internal vs. externally-forced variability – which remains particularly critical for a period of relatively weak external forcing, such as the last millennium – larger simulation ensembles and ensembles with decomposed forcing should become a standard procedure in paleoclimate modelling. At the same time, uncertainties in forcings and reconstruction need to be further reduced to be able to better validate models in the past with the goal of constraining their future response. Key targets for such constrains remain the sensitivity of temperature to solar and volcanic forcing and the climate-carbon cycle sensitivity.

20 *Acknowledgements.* We gratefully acknowledge Axel Timmermann, Bette Otto-Bliesner, Peter Lawrence, and Rosie Fisher for valuable discussions. We are grateful to the NCAR in Boulder, USA, for providing the code of the CESM, to the World Climate Research Programme's Working Group on Coupled Modelling, which is responsible for CMIP, to the climate modeling groups for producing and making available their model output. This study is supported by the Swiss National Science Foundation (grant no. 200020 147174), the European Commission through Seventh Framework Program (FP7) projects CARBOCHANGE (grant no. 264879) and Past4Future (grant no. 243908). J. Mignot has benefited from the support of the French Agence Nationale de la Recherche (HAMOC: ANR 13-BLAN-06-0003). The simulations for this study were performed on a CRAY XT5 and XE6 at the Swiss National Supercomputing Centre (CSCS) in Lugano.

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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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ESDD

6, 351–406, 2015

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Table 1. List of simulations conducted for this study. See text for details regarding the forcing. TSI = total solar irradiance, GHGs = greenhouse gases, E_{CO_2} = anthropogenic CO_2 emissions from fossil fuel burning and cement production. LULUC = land use and land use change.

Forcing	Control simulation (CTRL) 850 CE (500 years)	Transient simulation (CESM) 850–2099 CE
TSI	1360.228 W m^{-2}	adjusted Vieira and Solanki (2010) and Lean et al. (2005)
Volcanic	none	Gao et al. (2008)
GHGs	CO_2 (279.3 ppm) CH_4 (674.5 ppb) N_2O (266.9 ppb)	Schmidt et al. (2011)
E_{CO_2}	none	Andres et al. (2012) and Moss et al. (2010)
Aerosol	1850 CE from Lamarque et al. (2010)	Lamarque et al. (2010, 2011)
Orbital	1990 CE after Berger (1978)	1990 CE after Berger (1978)
LULUC	850 CE from Pongratz et al. (2008)	Pongratz et al. (2008) and Hurtt et al. (2011)

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

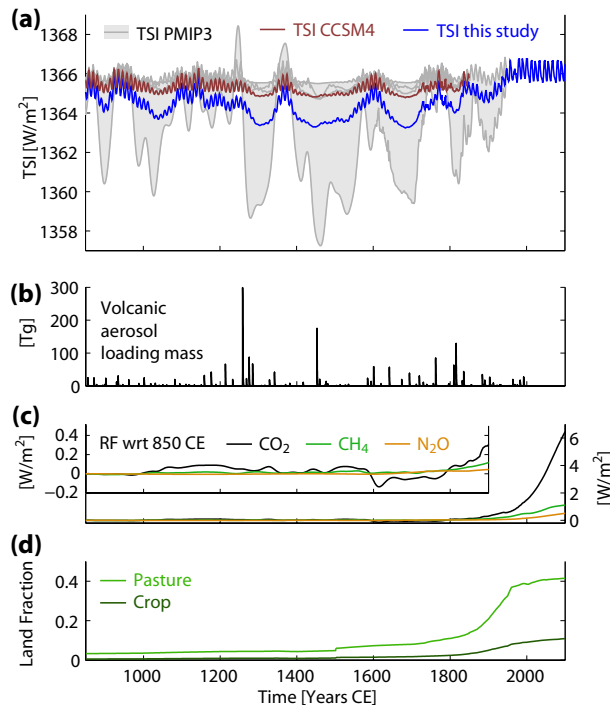


Figure 1. Forcings used in the last millennium simulation with CESM. **(a)** TSI in comparison with the different TSI reconstructions proposed by PMIP3. **(b)** Volcanic forcing as total volcanic aerosol mass. **(c)** Radiative forcing (RF, calculated according to IPCC, 2001) from the greenhouse gases CO_2 , CH_4 , and N_2O . **(d)** Major changes in land cover (as fraction of global land area). See text for details.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

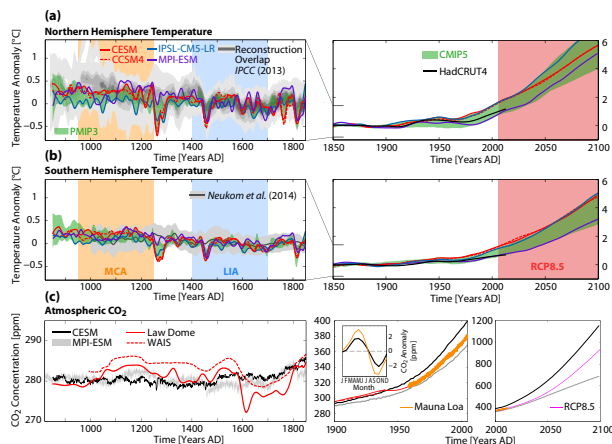


Figure 2. (a) Northern Hemisphere and (b) Southern Hemisphere temperature anomalies in model simulations and reconstructions. The anomalies are with reference to 1500–1899 CE (left panels) and 1850–1899 CE (right panels). Gray shading in (a) indicates the reconstruction overlap (IPCC, 2013), in (b) the reconstruction by Neukom et al. (2014). The 5–95 % range of the simulations from the third Paleoclimate Modelling Intercomparison Project (PMIP3) and the fifth Coupled Model Intercomparison Project (CMIP5; applying the RCP 8.5) are given in green and red shading, respectively. Note that MPI-ESM applies the A1B scenario (IPCC, 2000), which has a weaker forcing than RCP 8.5. Hemispheric means from observations are shown as thick black line (Cowtan and Way, 2014). All time series have been smoothed by a local regression filter which suppresses variability higher than 30 years. The Medieval Climate Anomaly (MCA) and the Little Ice Age (LIA) are indicated as defined in Mann et al. (2009). (c) Evolution of atmospheric CO₂ in CESM (black), MPI-ESM (grey; ensemble range), from ice cores (red), from measurements (orange), and from RCP8.5 used to force the radiative code in CESM (magenta). The small inset in the middle panel shows the observed annual cycle at Mauna Loa, Hawaii, and a 2° × 2° average over Hawaii from CESM, both derived from the period 1958–2012.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

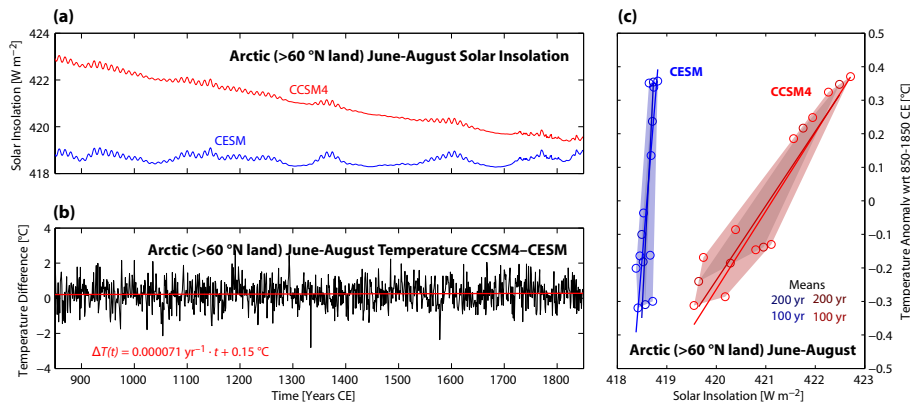


Figure 3. (a) Mean June–August (JJA) Arctic (> 60° N land) solar insolation in CCSM4 with time-varying orbital parameters and CESM with fixed orbital parameters. (b) Arctic JJA temperature difference between CCSM4 and CESM. The least-squares linear trend of this temperature difference is given in red. (c) Arctic JJA temperature anomalies (from their 850–1850 AD mean) vs. solar insolation as 100 and 200 year averages (10 and 5 circles, respectively) from CCSM4 and CESM (red and blue, respectively). The least-squares linear trend for each cloud of 100 and 200 year averages is given in the respective color. The shading envelops the range of temperature vs. solar insolation for each cloud of means.

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

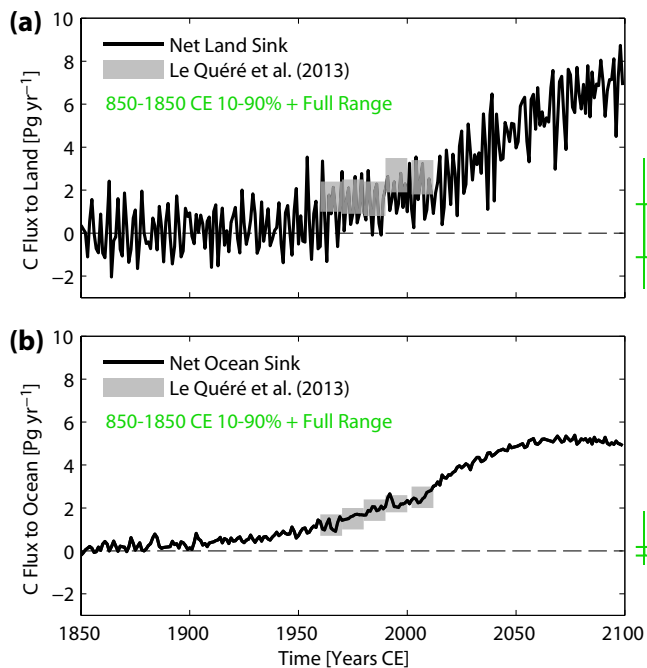


Figure 4. Annual mean net carbon flux from the atmosphere to (a) land and (b) ocean in CESM. Green bars give the full and 10–90 % range from the preindustrial part of the simulation. Observational estimates are from Le Quéré et al. (2013).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



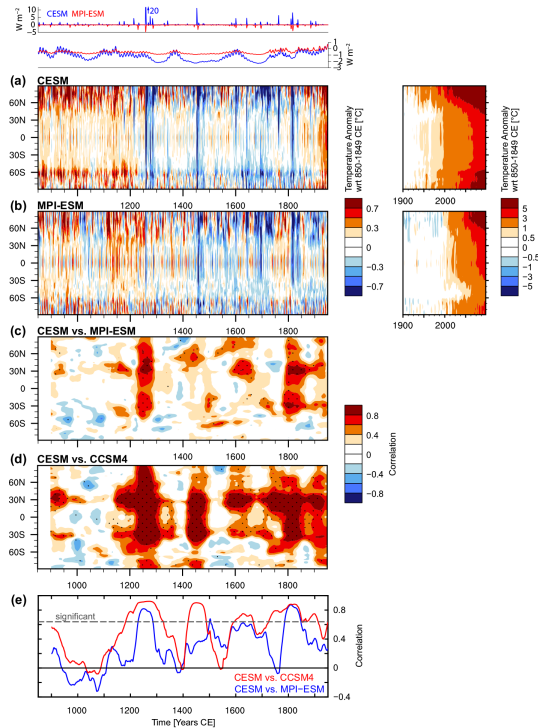


Figure 5. 5 year filtered zonal mean anomalies of surface air temperature (SAT), relative to 850–1849 CE from **(a)** CESM and **(b)** MPI-ESM. **(c)** 100 year running-window correlation of zonal mean SAT from CESM and MPI-ESM. 0.75 Tukey window has been applied to the data before correlation to weaken sharp transitions. Stippling indicates significance (5% level), taking into account autocorrelation estimated from the entire time period. **(d)** As **(c)** but for the correlation of CESM with CCSM4. **(e)** As **(d)** but for global mean SAT. Small inset on top shows volcanic and solar forcing of CESM and MPI-ESM. Volcanic forcing of CESM scaled to have the same radiative forcing as MPI-ESM for Pinatubo in 1991 CE. Solar forcing relative to 1850 CE.

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

Title Page	
Abstract	Introduction
Conclusions	References
Tables	Figures
◀	▶
◀	▶
Back	Close
Full Screen / Esc	
Printer-friendly Version	
Interactive Discussion	



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

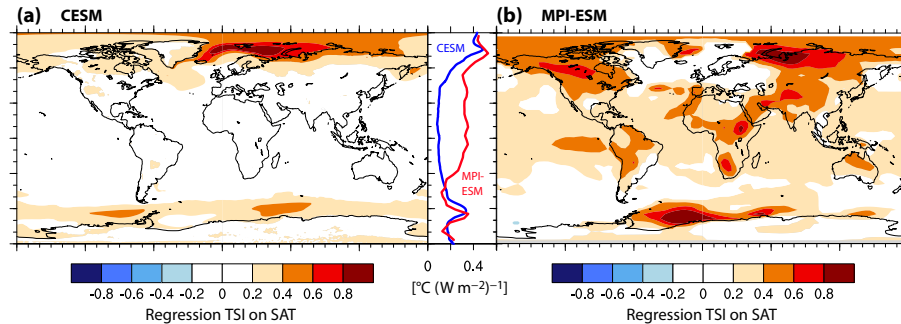


Figure 6. Regression of total solar irradiance (TSI) on surface air temperature (SAT) for the period 850–1850 CE in **(a)** CESM and **(b)** MPI-ESM. Time series at each gridpoint have been 5 year filtered. Only significant regression coefficients at the 5% level are shown. The small panel shows zonal means.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[◀](#)[▶](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

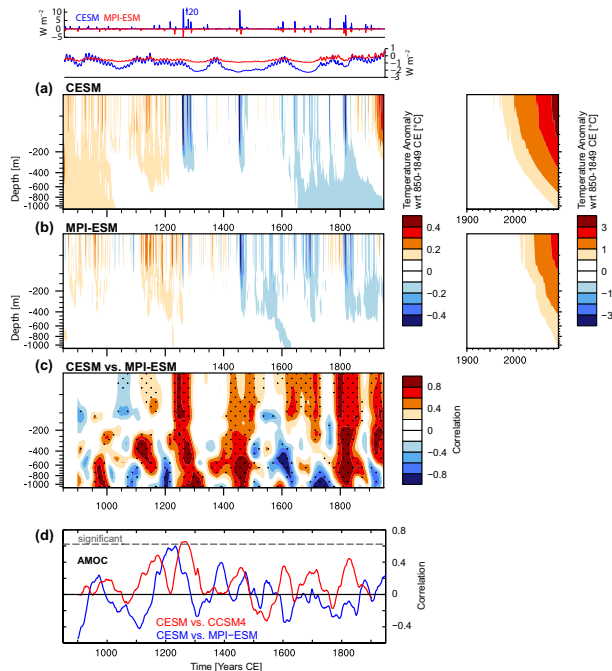


Figure 7. 5 year filtered zonal mean anomalies of horizontally averaged ocean temperature, relative to 850–1849 CE from (a) CESM and (b) MPI-ESM. (c) 100 year running-window correlation of zonal mean SAT from CESM and MPI-ESM. A 0.75 Tukey window has been applied to the data before correlation to weaken sharp transitions. Stippling indicates significance at the 5% level, taking into account autocorrelation estimated from the entire time period. (d) 100 year running-window correlation of the Atlantic Meridional Overturning Circulation (AMOC) in CESM and MPI-ESM.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

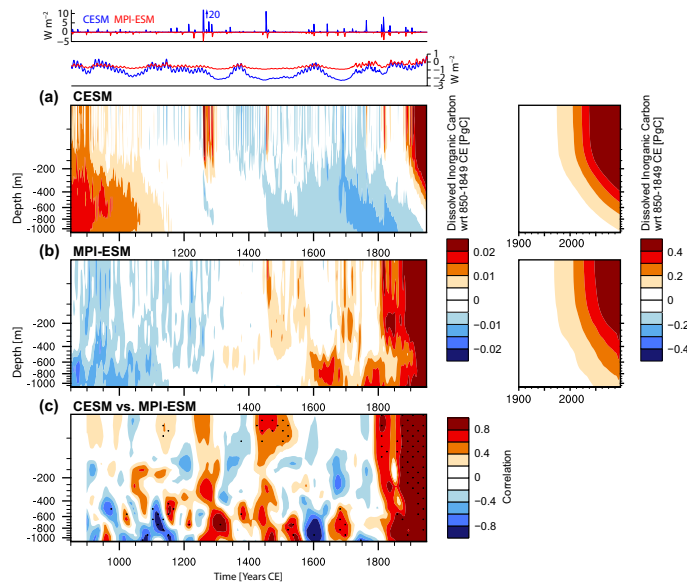


Figure 8. 5 year filtered zonal mean anomalies of horizontally integrated dissolved inorganic carbon (DIC), relative to 850–1849 CE from (a) CESM and (b) MPI-ESM. (c) 100 year running-window correlation of zonal mean SAT from CESM and MPI-ESM. A 0.75 Tukey window has been applied to the data before correlation to weaken sharp transitions. Stippling indicates significance at the 5 % level, taking into account autocorrelation estimated from the entire time period.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

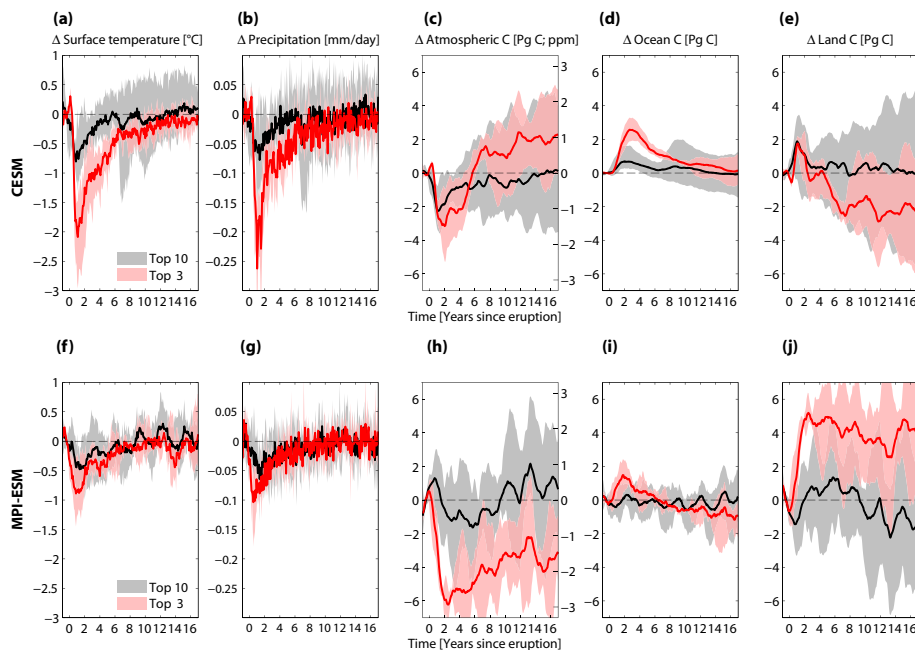


Figure 9. Superposed Epoch Analysis of the strongest three (top3) and following strongest seven eruptions (top10) of the period 850–1850 CE in **(a–e)** CESM and **(f–j)** MPI-ESM for **(a and f)** global mean surface air temperature, **(b and g)** global mean precipitation, **(c and h)** atmospheric carbon given in PgC on the left y axis and in ppm on the right y axis, **(d and i)** ocean carbon, and **(e and j)** land carbon. Time series are deseasonalized and calculated as anomalies to the mean of the preceding five years. The shading shows the 10–90 % range.

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Printer-friendly Version](#)
[Interactive Discussion](#)


Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

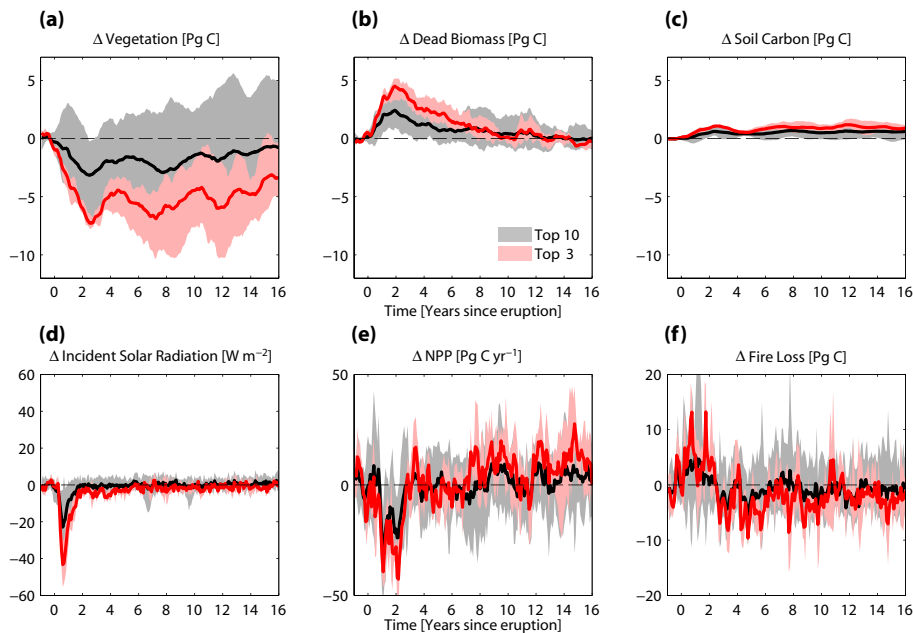


Figure 10. Superposed Epoch Analysis of the strongest three (top3) and following strongest seven eruptions (top10) for tropical land (25° S to 25° N) in CESM during the period 850–1850 CE. Land carbon inventory changes split up in (a) vegetation, (b) dead biomass (litter and wooden debris), and (c) soil. Further, changes in (d) solar radiation, (e) net primary production (NPP), and (f) loss of carbon through fire. Time series are deseasonalized and calculated as anomalies to the mean of the preceding five years. The shading shows the 10–90 % range.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

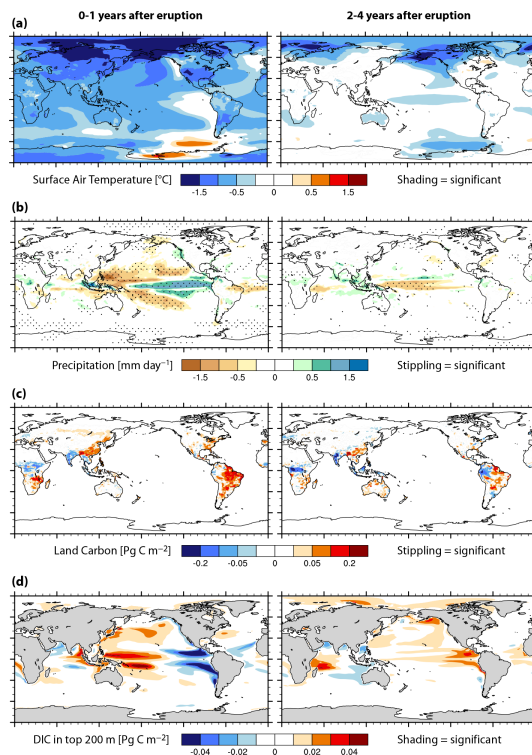


Figure 11. Composites of top10 post-volcanic eruption years as anomalies to the preceding 5 years, averaged over (left) the first 2 years starting with the year of the eruption, and (right) the following three years. **(a)** Surface air temperature, **(b)** precipitation, **(c)** total land carbon, **(d)** dissolved inorganic carbon (DIC) integrated over the top 200 m. Shading or stippling indicates significance at the 5 % level. Note, that for land carbon at an individual grid cell hardly any significant changes are detected due to the large inter-annual variability.

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Printer-friendly Version](#)
[Interactive Discussion](#)

Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

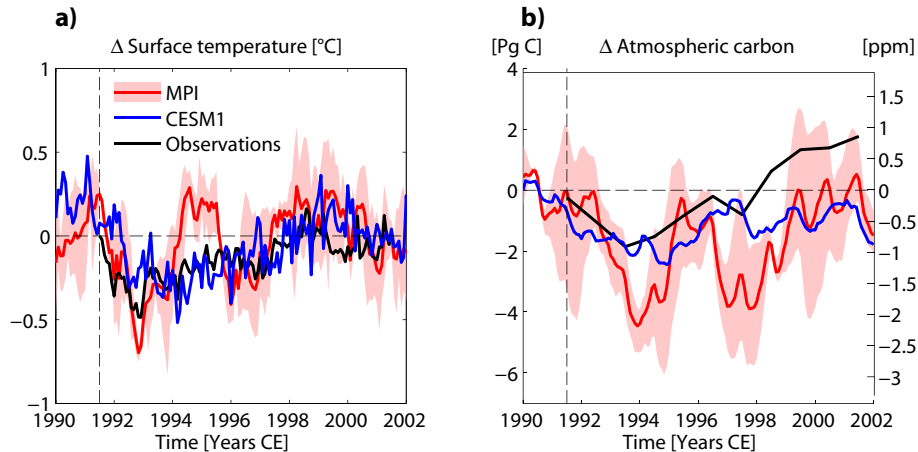


Figure 12. Global mean changes in response to Pinatubo. **(a)** Global mean surface air temperature and **(b)** atmospheric carbon, both deseasonalized and linearly detrended over 30 years centered on June 1991; temperature observations were corrected for El Niño-Southern Oscillation and other dynamical components (Thompson et al., 2009), CO₂ observations were corrected for El Niño-Southern Oscillation and anthropogenic emissions (Frölicher et al., 2013).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Climate and carbon cycle dynamics in a CESM simulation from 850–2100 CE

F. Lehner et al.

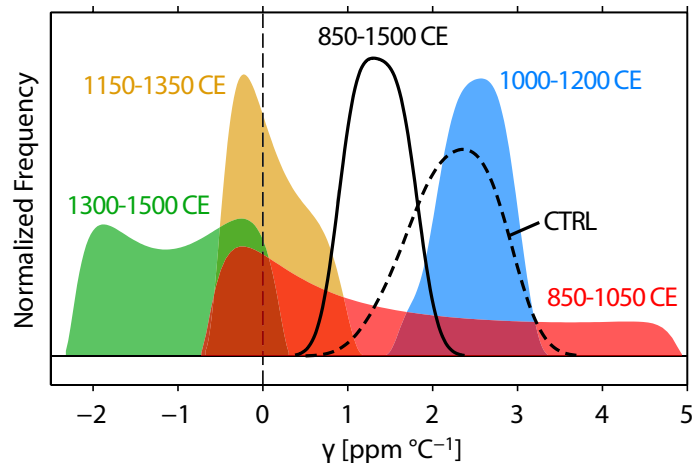


Figure 13. Temporal dependence of the climate carbon cycle sensitivity γ in CESM. Normalized probability density functions (PDF) of γ for 200 year windows overlapping by 50 years (color-filled), for the full period 850–1500 CE (black solid), and for the CTRL (black dashed). The spread of each PDF arises from the range of low-pass filters applied (20 to 120 years).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

